

Fundamental Physics with Hydrogen in emission and in absorption

Matteo Viel - SISSA (Trieste, Italy)

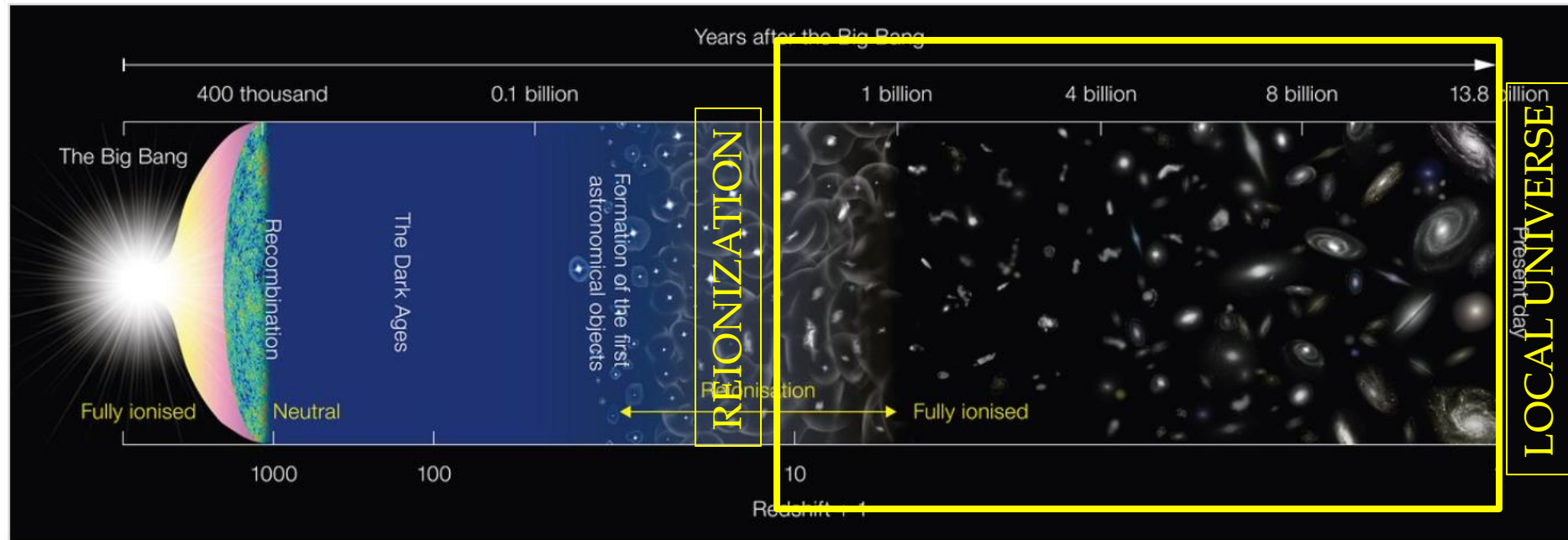
12/06/25

COSMOFONDUE



Promises of the post-reionization Universe

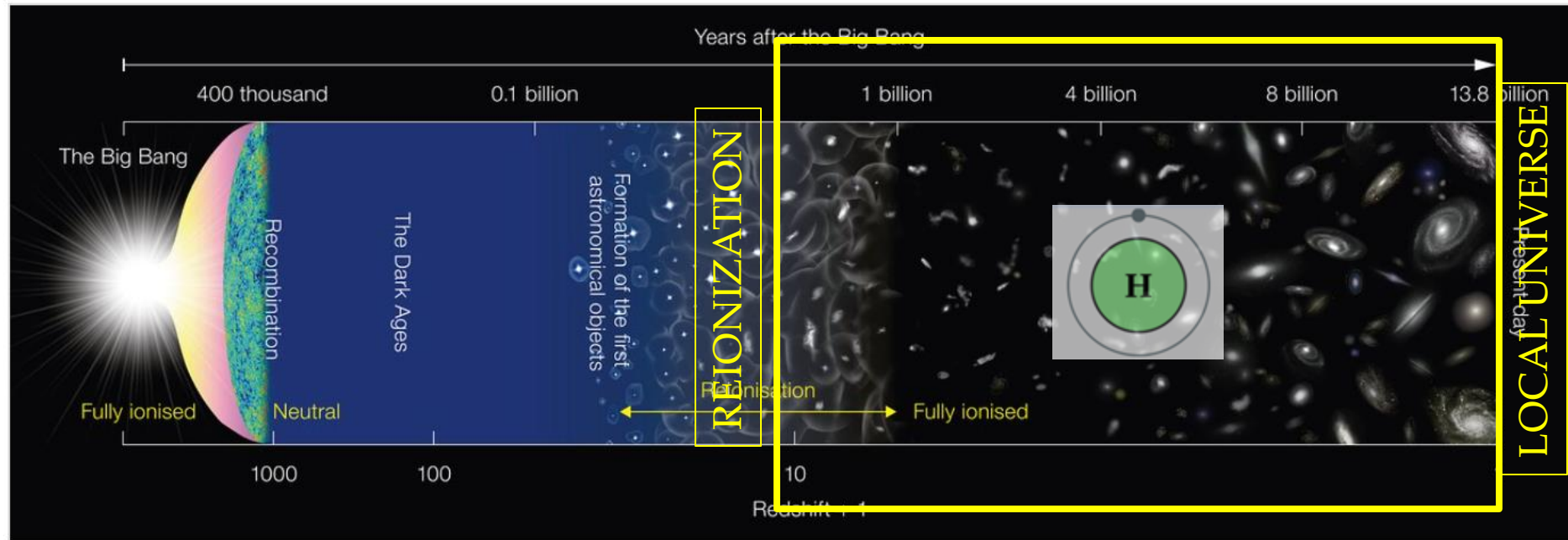
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- Ideal place to test **structure formation** processes
- and cosmological models in and beyond Λ CDM (Universe being more linear)

Promises of the post-reionization Universe

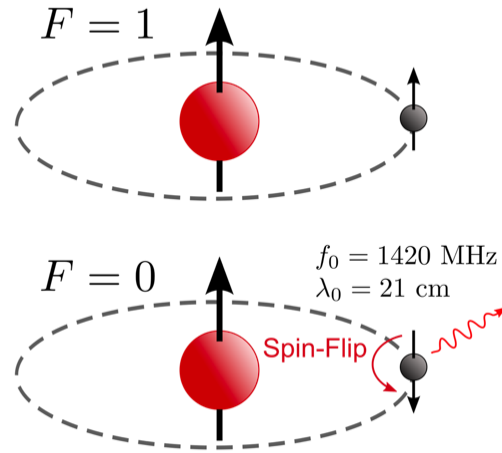
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- Ideal place to test **structure formation** processes
- and cosmological models in and beyond Λ CDM (Universe being more linear)
- **Large volume** to be probed
- ... but HI tracer is sensitive to **small scale** (astro) physics (intrinsically no threshold)
- And can probe **underdense regions** far from galaxies too

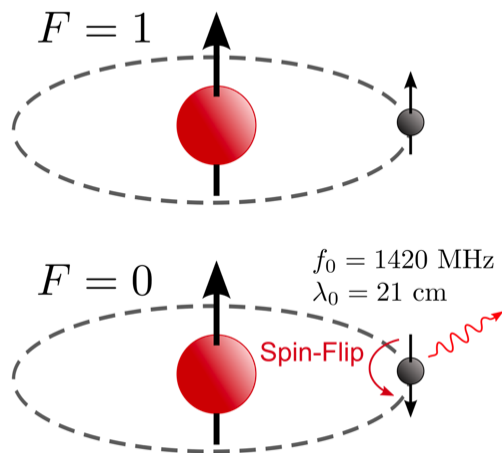
Clustering of LSS tracers ($> \text{Mpc}$) is coupled to the astrophysics at scales: $O(10 \text{ pc})$ - $O(100 \text{ kpc})$

EMISSION



- Spin-flip electron transition
- In the post-reio epoch HI is mainly in haloes
- Which dominate the emission signal
- **Mass weighted view** of the HI distribution
- Need to specify $M_{\text{HI}}(M_{\text{halo}})$ or Line luminosity as a function of M_{halo}

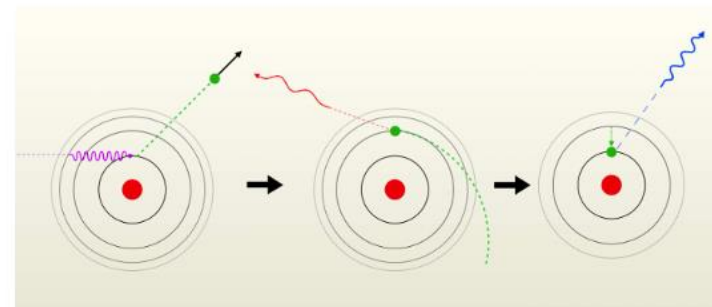
EMISSION



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ABSORPTION

HI

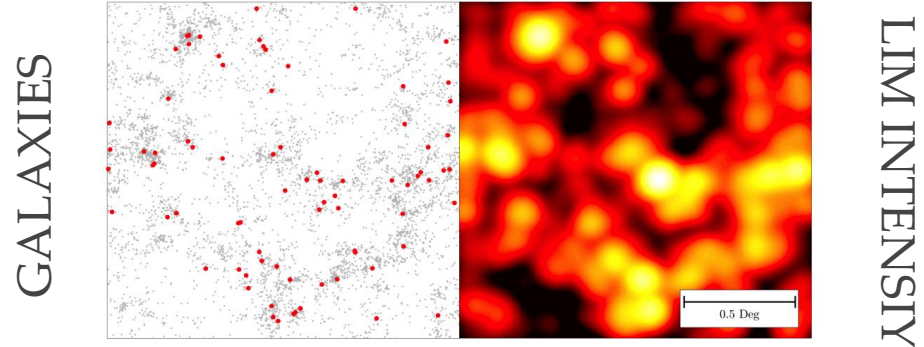
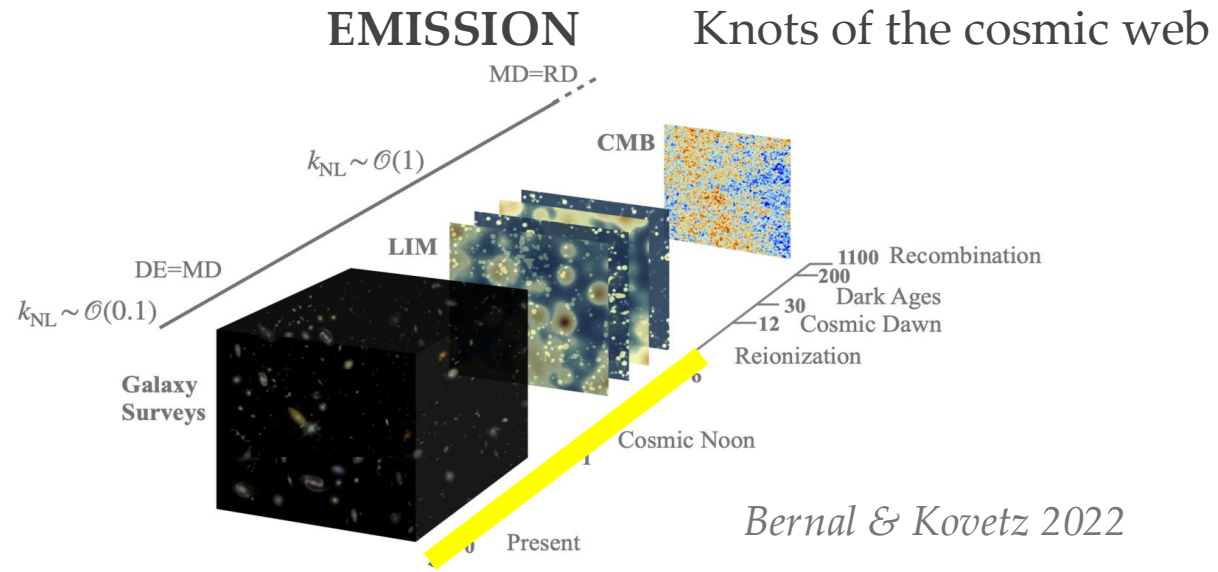


$$\lambda = \lambda_0(1 + z)$$
$$\lambda_0 = 1215.67 \text{ \AA}$$

- Lyman- α scattering
- Need a bright source behind
- Neutral fraction in most of the volume is $<1.e-5$
- **Volume weighted view** of the HI distribution
- Need to model Flux-DM density relation (very non-linear transform)

Holistic view on HI distribution

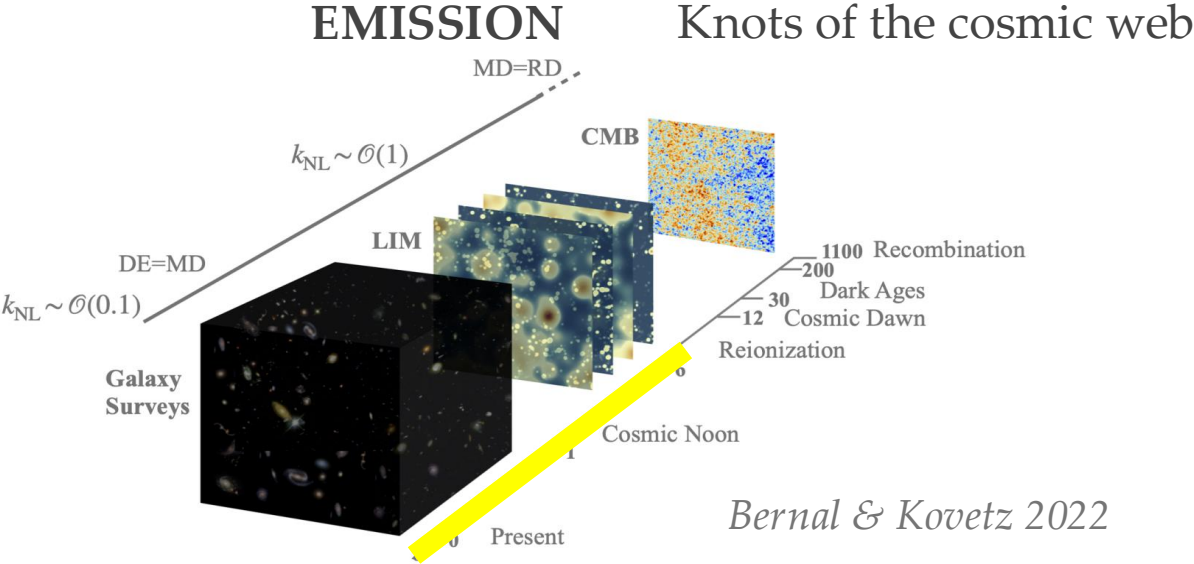
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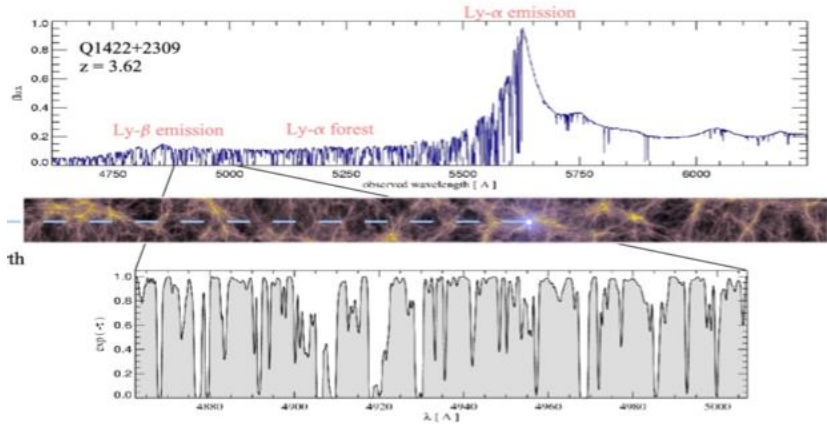
Kovetz+18

Holistic view on HI distribution

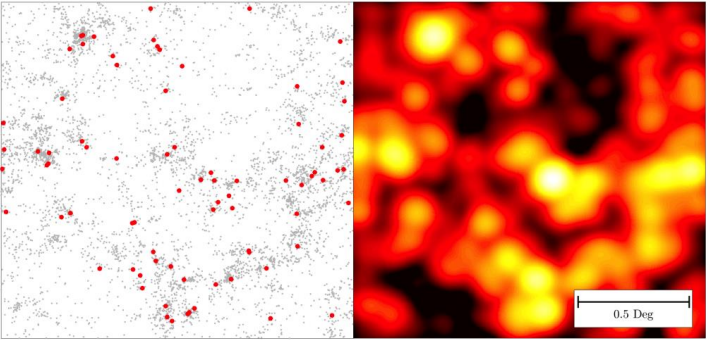
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ABSORPTION Cosmic web

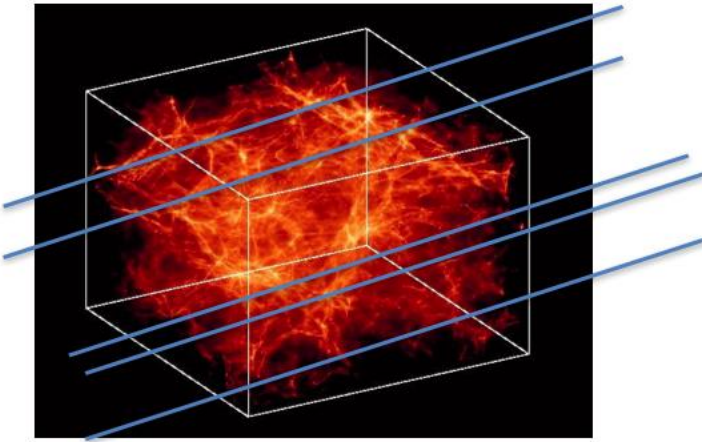


GALAXIES



LIM INTENSITY

Kovetz+18



Quasars

Bolton+18, Puchwein+23

What are the big questions?

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- What is Dark Matter?
- Is evolving Dark Energy real?
- Can we measure neutrino masses?
- Can we probe the matter power spectrum down to the smallest scales?
- Can we test inflation?
- Is there new physics like Primordial Magnetic Fields?

What are the big questions?

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Signature in clustering, decay, annihilation

BAOs

Neutrino free-streaming at large/medium scales

DM acoustic oscillations, suppression of power

Non Gaussianity

Increase of power at small scales

Promises of the post-reionization Universe

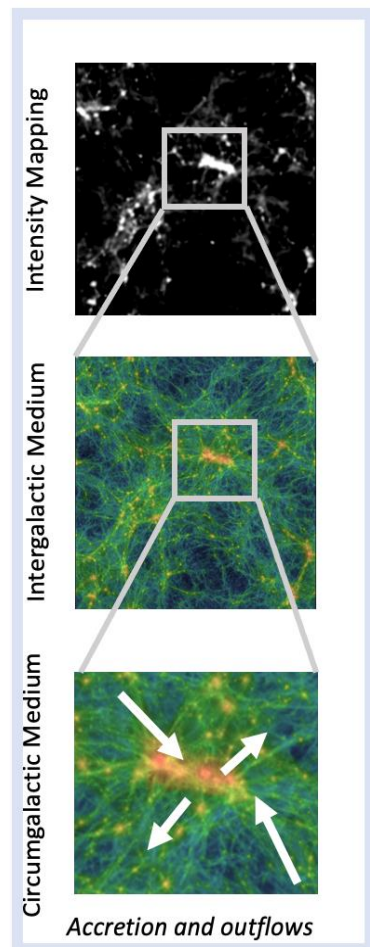
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Long lever arm in terms of scales/redshifts will in turn allow to break degeneracies between astro and cosmo parameters with:

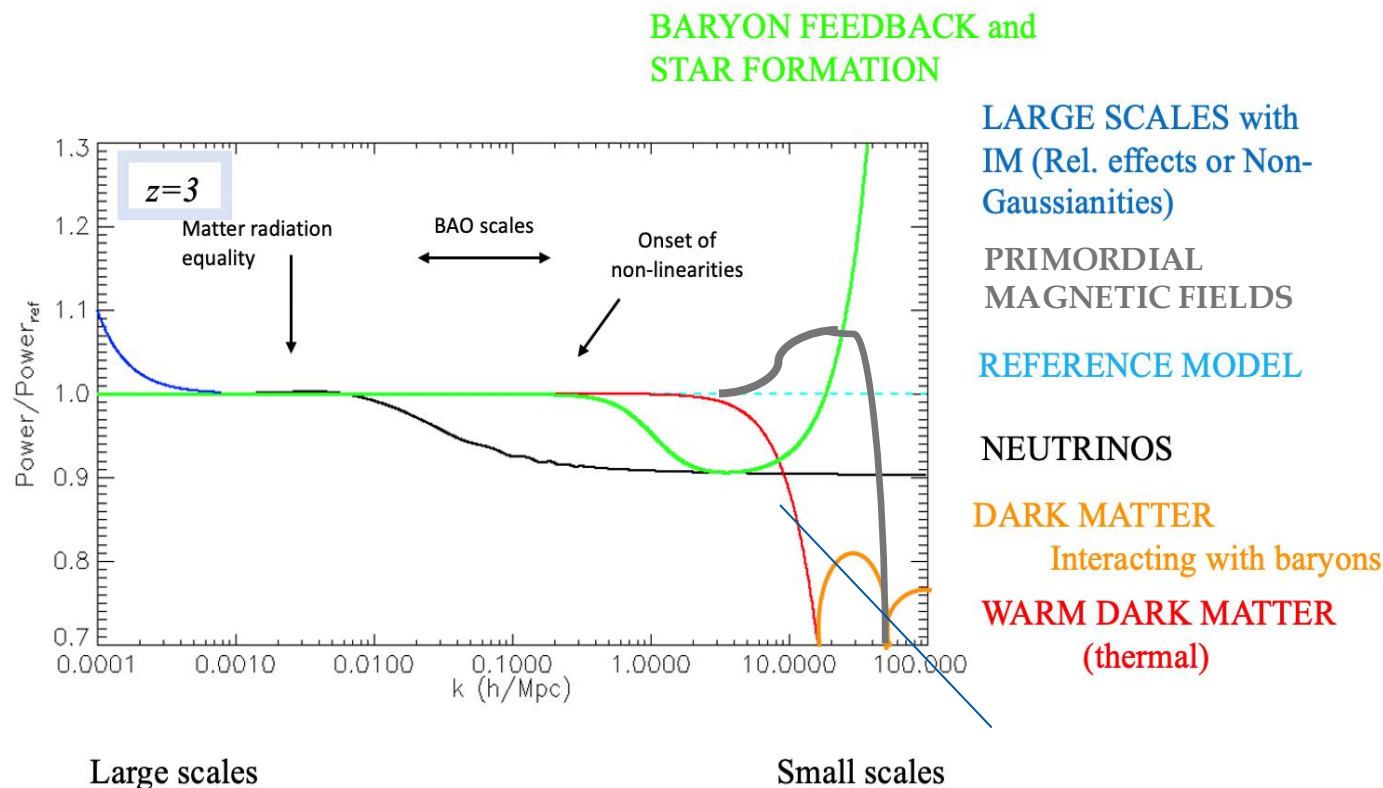
- **Power spectrum**
- **cross-correlations** of different tracers
- **new estimators** (e.g. 1-point function, bispectrum, Machine Learning)

It is an “active phase” of structure formation processes (feedback, star formation, black holes, cosmic bayron cycle etc.)

Environments



Physical Scales

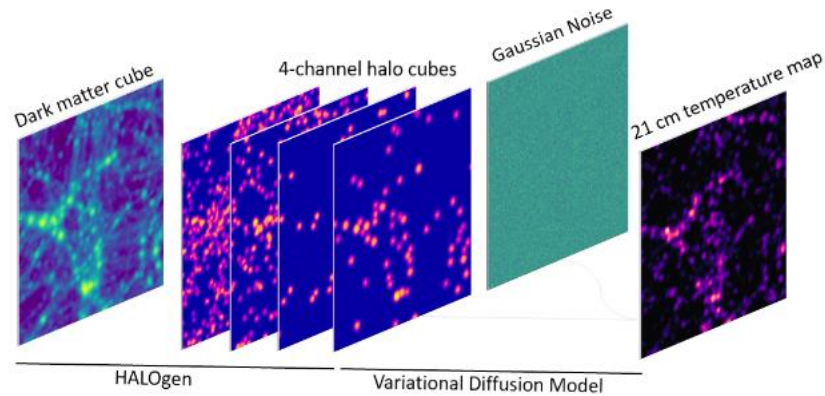


HI measures density perturbations in a matter dominated regime!

INTENSITY MAPPING

- 1) Modelling
- 2) Small scales
- 3) BAO

LODI: Latent Overlap Diffusion for Intensity Mapping



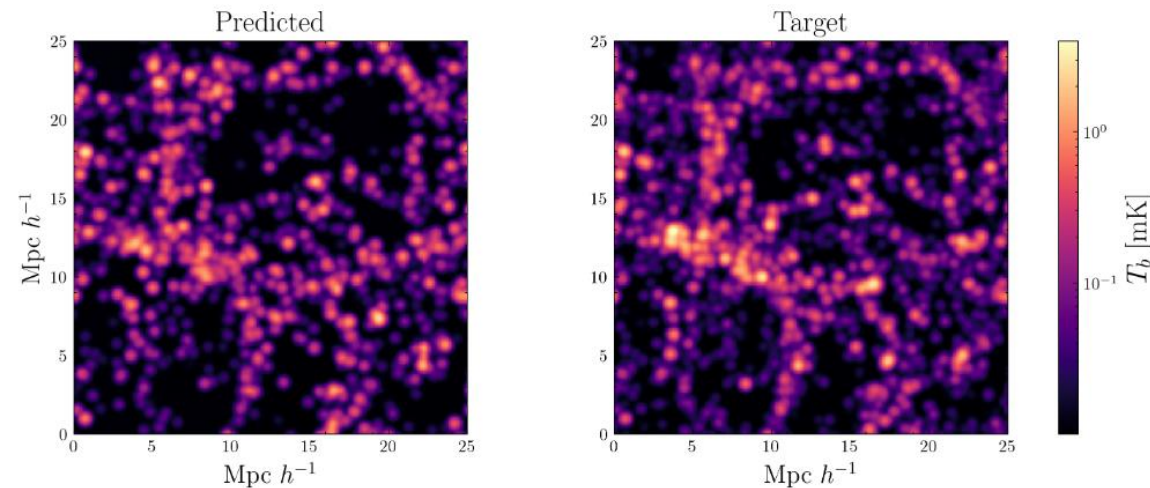
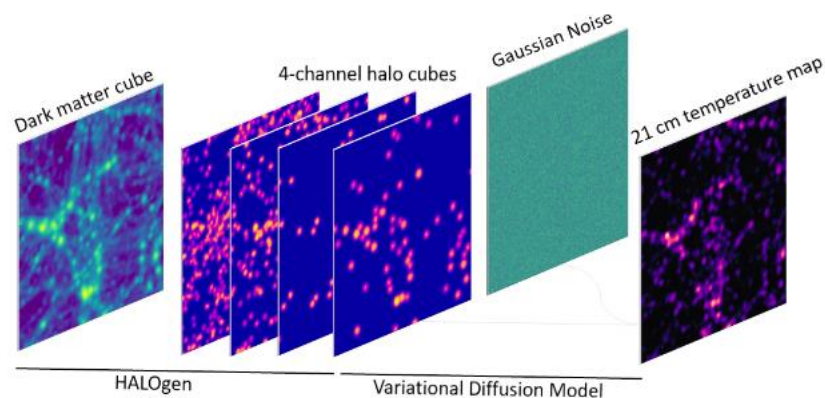
1st step: DM \rightarrow haloes (via U-Net)

2nd step: Haloes \rightarrow Intensity voxel map (via Diffusion model)

Methods: Learning LIM with diffusion models

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LODI: Latent Overlap Diffusion for Intensity Mapping

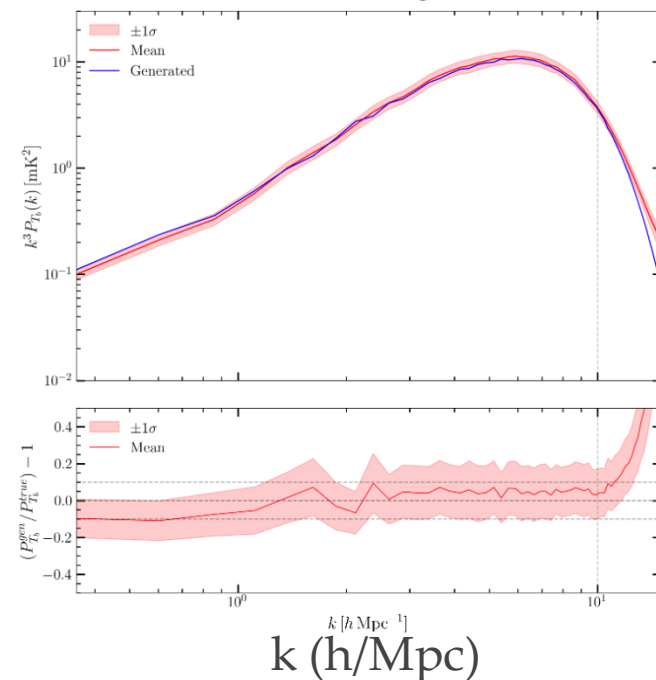


1st step: DM \rightarrow haloes (via U-Net)

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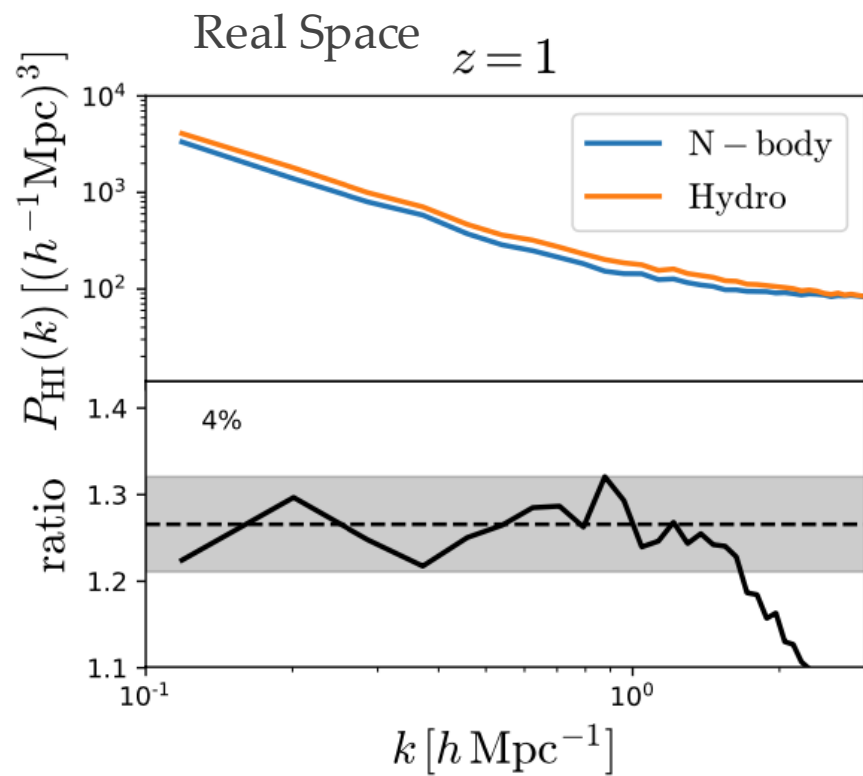
- Trained on CAMELS simulations
- Agreement up to $k = 10 \text{ h/Mpc}$
- Extendable to other LIM lines
- Good for likelihood free field-level inference

21cm power spectrum



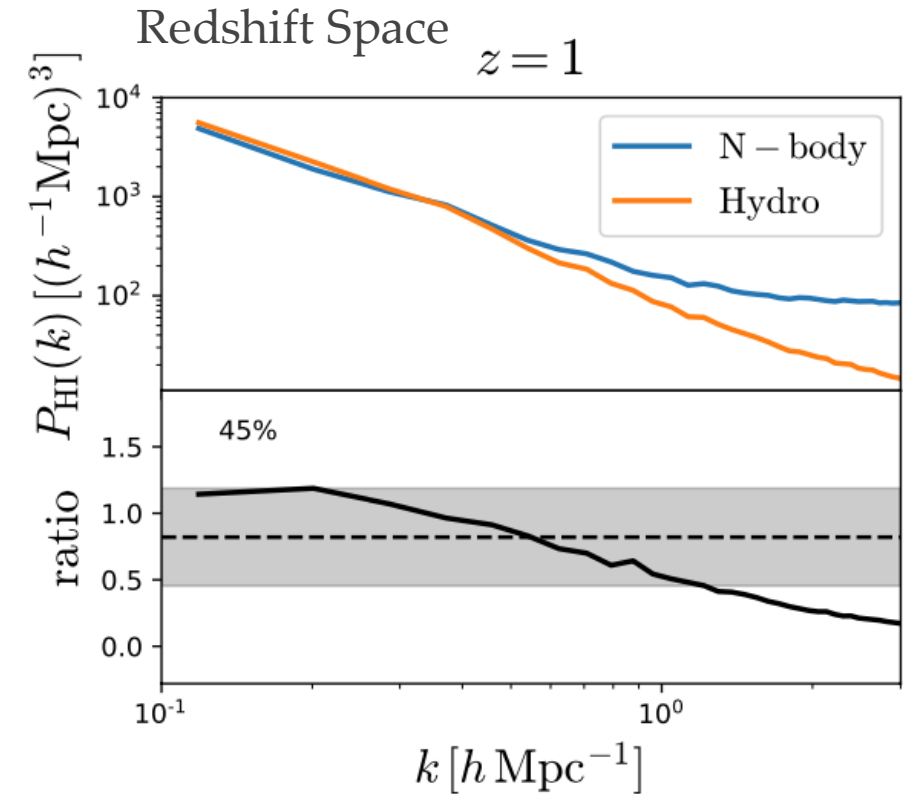
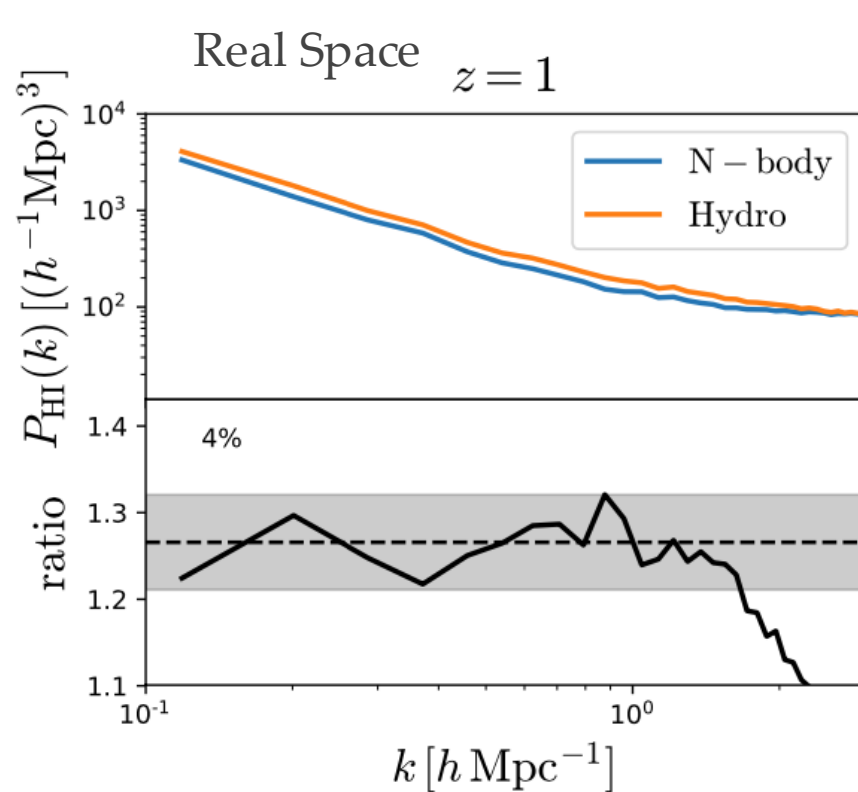
Modelling: populating haloes with HI

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Modelling: populating haloes with HI

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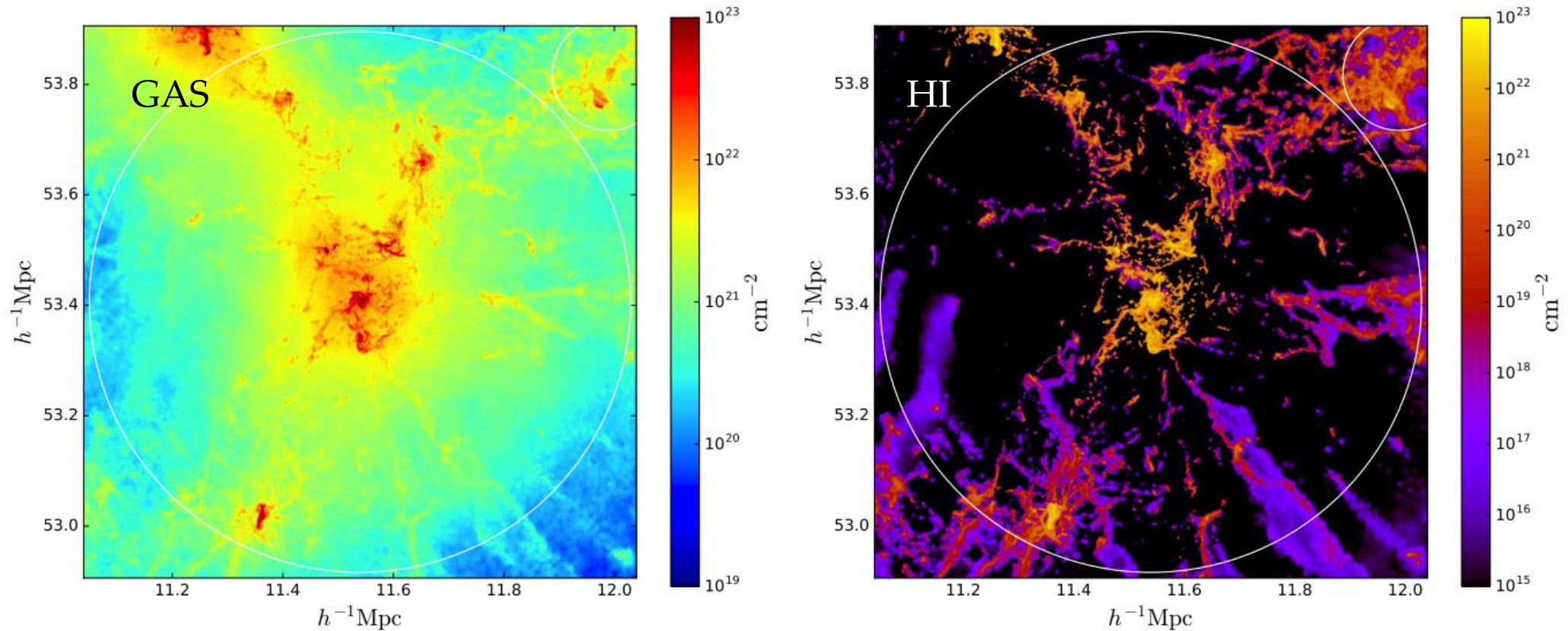


- N-body (all HI is in the center of the halo) vs. full hydro HI power spectrum (effectively there is a 1-halo term). Normalization is quite different but shape is reasonable.
- Kaiser effect (boosts power at large scales) vs. Finger of Gods (suppresses power at small scales but not so small). Behaviour in matter field and HI field is different.....!

How is HI distributed in cold DM haloes?

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$10^{13} M_{\odot}$ halo at $z=3$



Feedback/star formation is shaping the properties of HI...
...unfortunately this above cannot be directly observed

How is HI distributed in non-cold DM haloes at $z=0$?

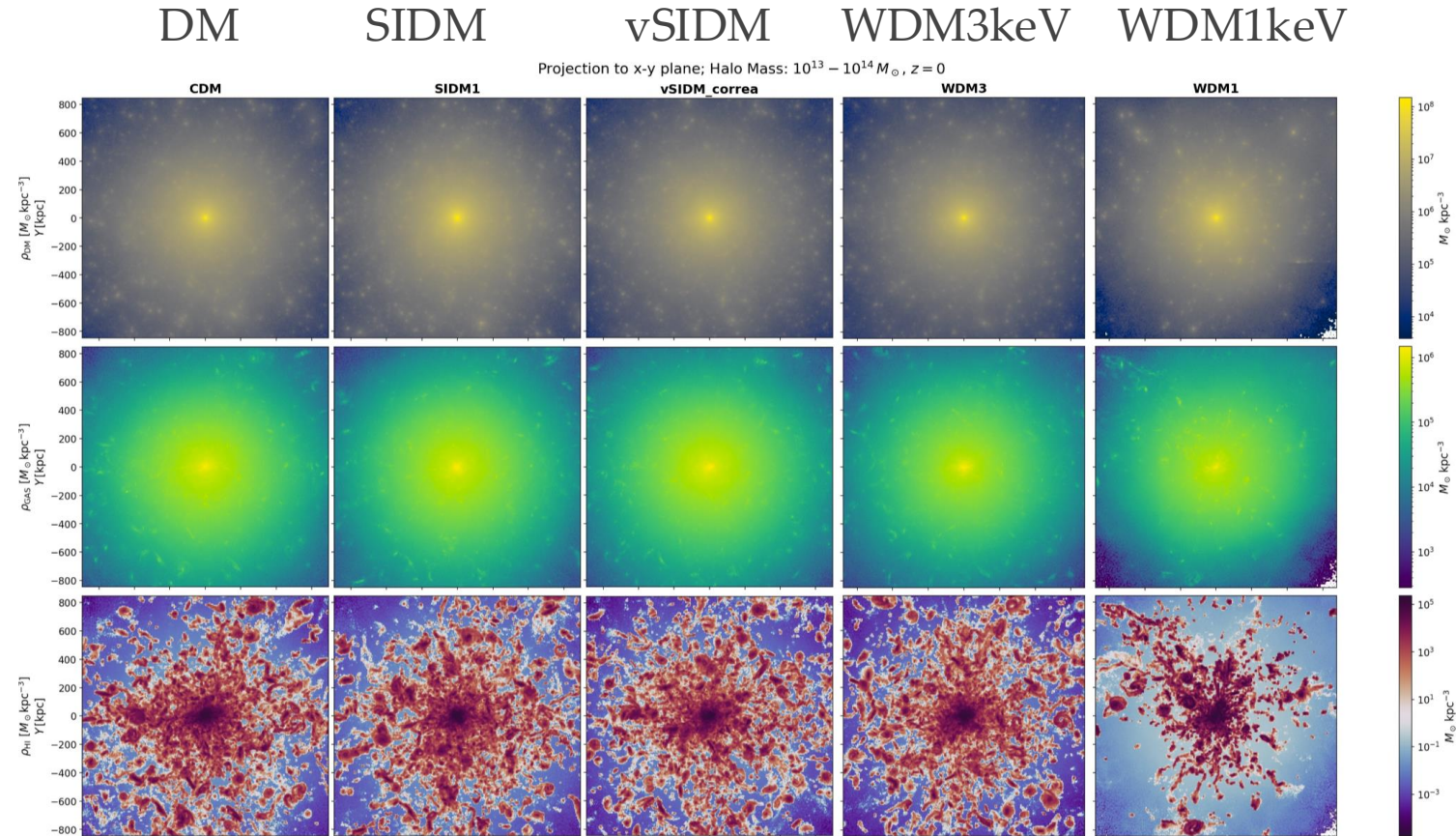
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Redshift $z=0$
stacking of 100
haloes $10^{13}-10^{14} M_{\odot}$

DM

GAS

HI



How is HI distributed in non-cold DM haloes at $z=3$?

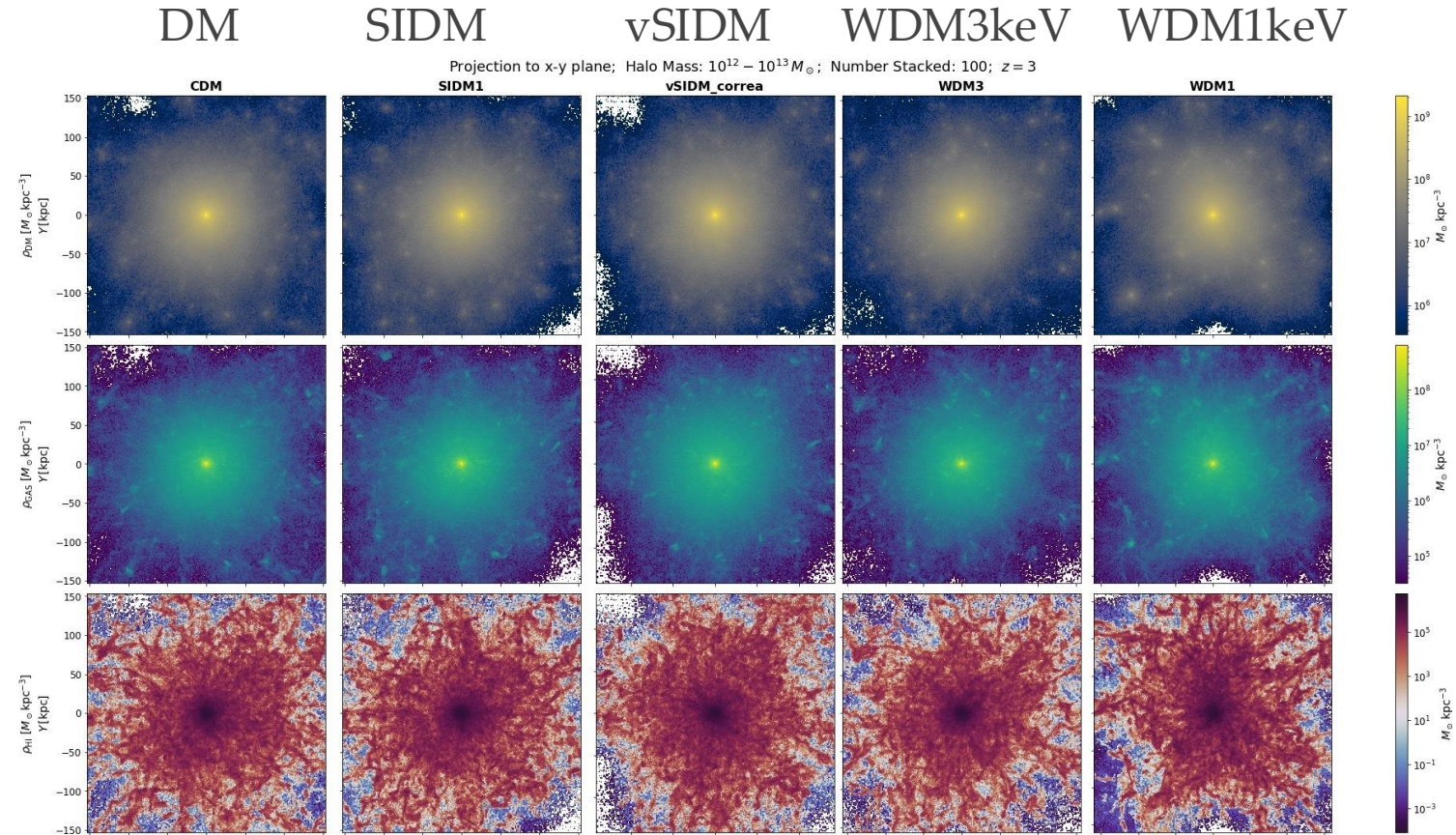
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Redshift $z=3$
stacking of 100
haloes $10^{12}-10^{13} M_{\odot}$

DM

GAS

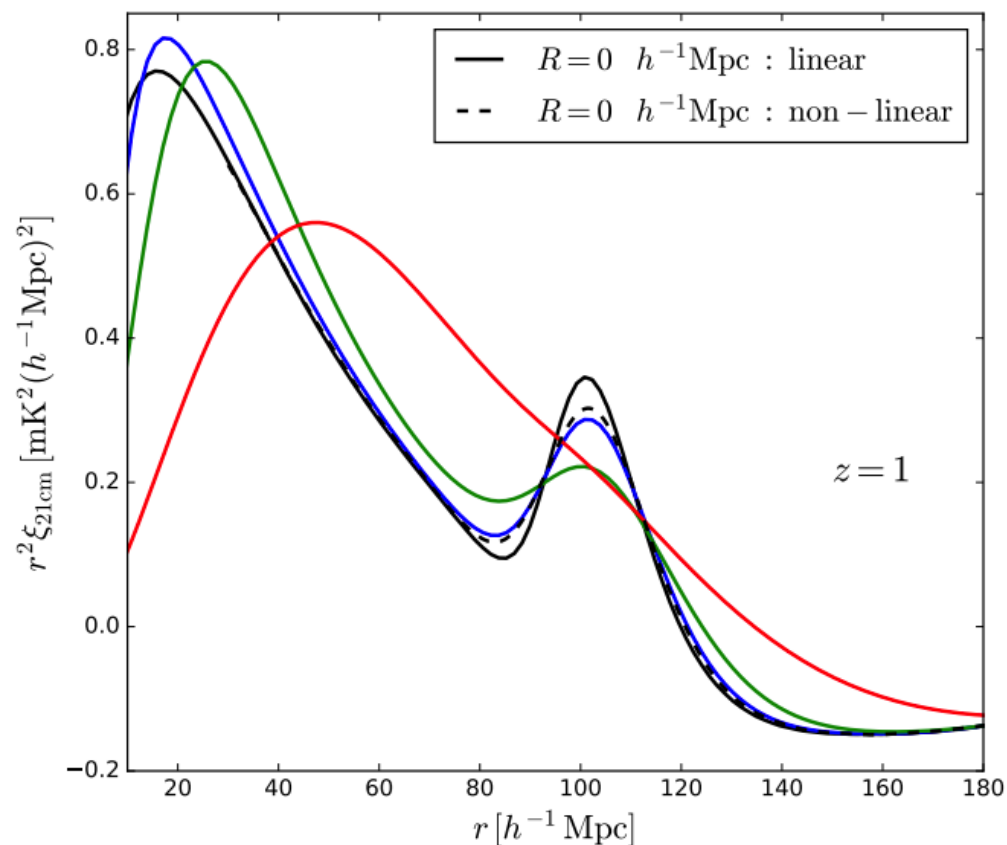
HI



Baryonic Acoustic Oscillations in 21cm IM

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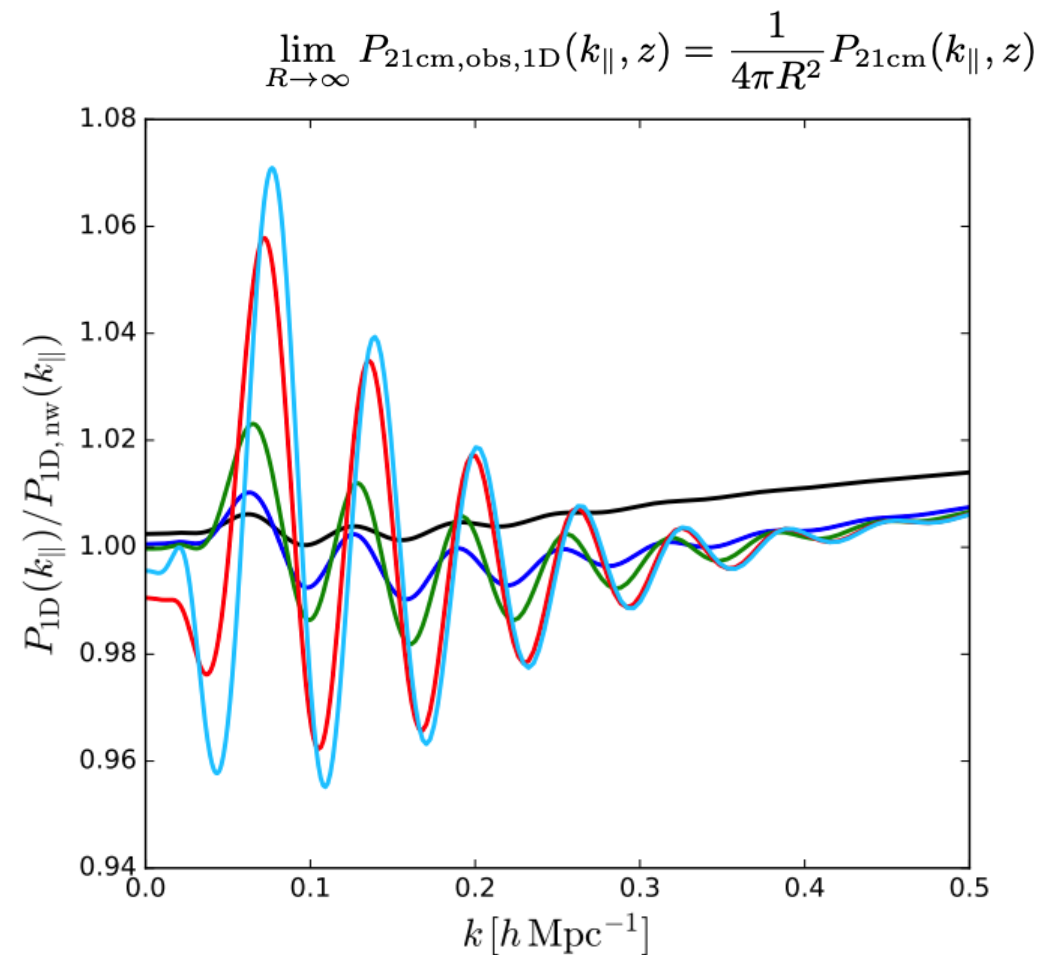
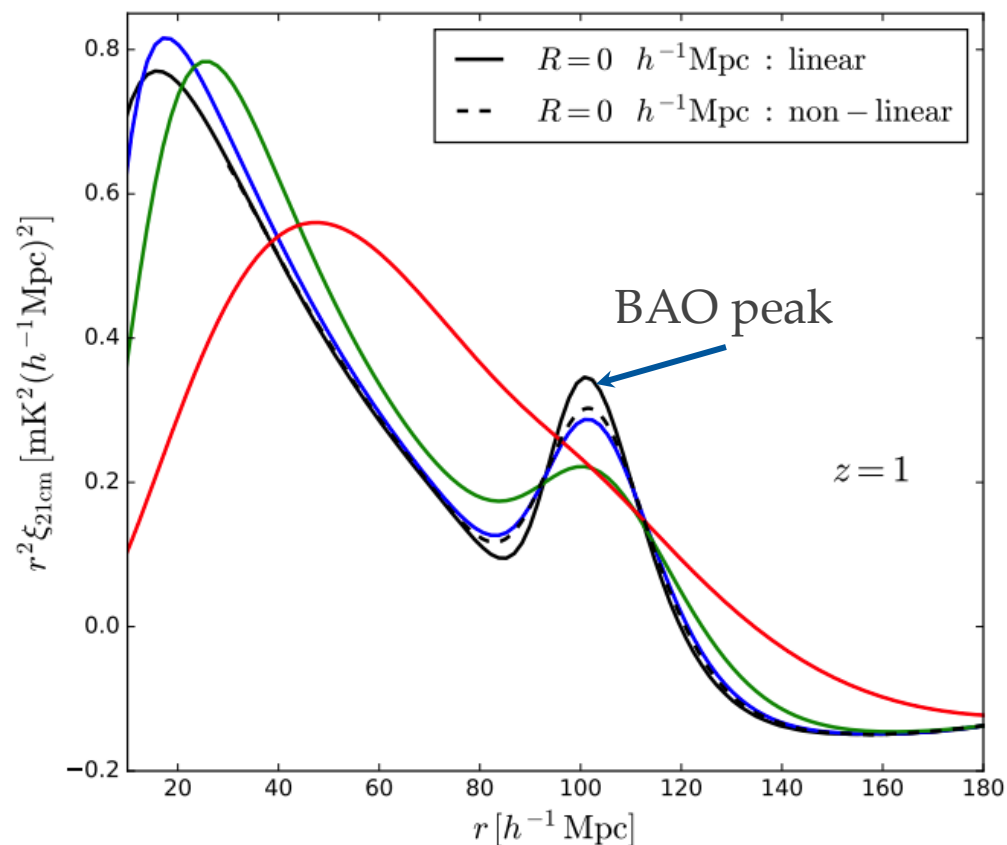
- Poor angular resolution, will smooth BAO feature
- But in the k parallel direction, frequency resolution is very high \rightarrow radial BAO
- 1D power is reduced in amplitude compared to 3D but wiggles are prominent



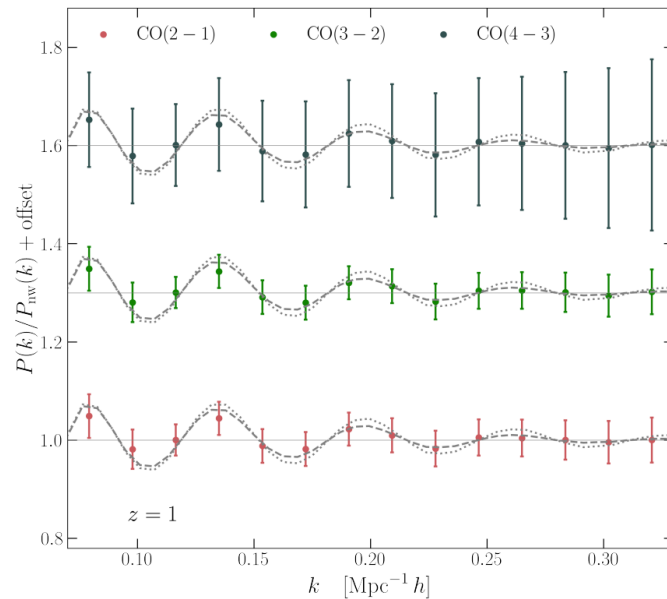
Baryonic Acoustic Oscillations in 21cm IM

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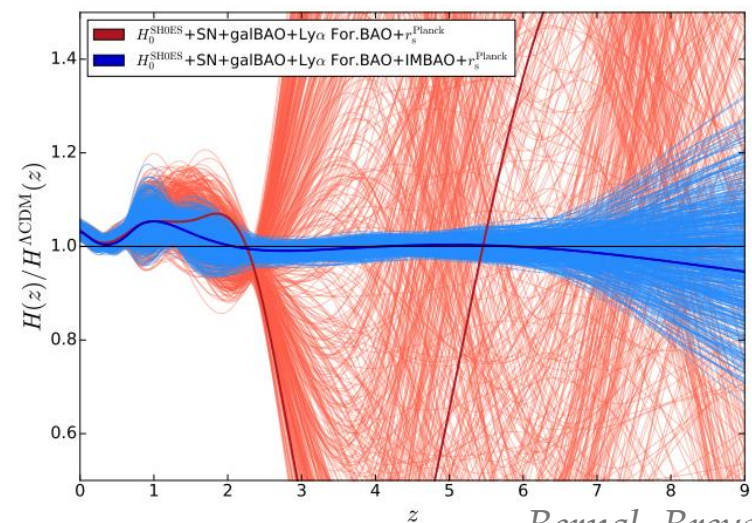
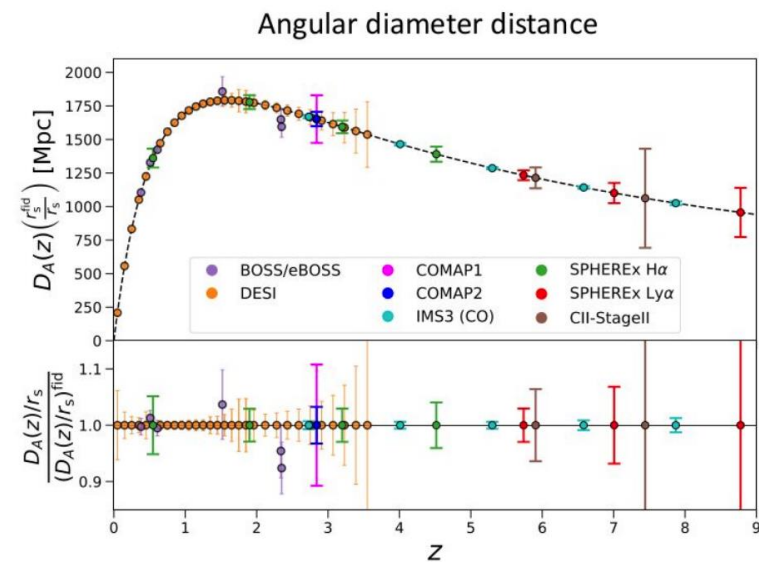
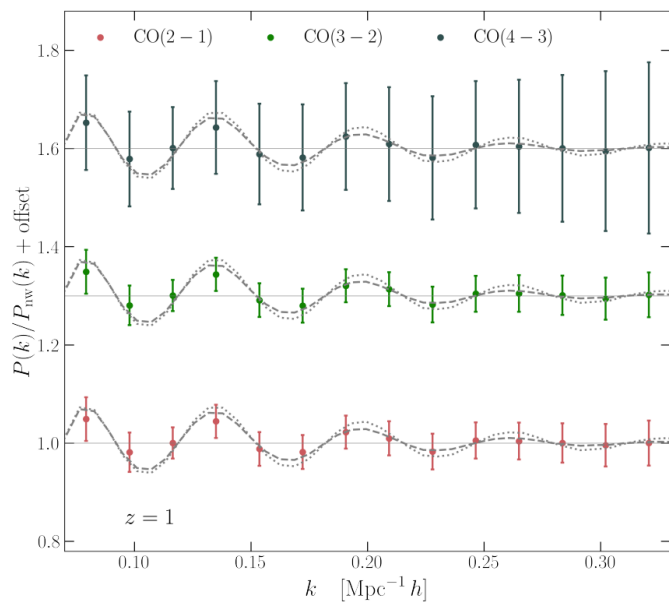
- 3% error on BAO peak position at $z=2.5$ with noise and foregrounds for SKA
- This could improve to 1.7% with BAO pixel-reconstruction
- And other IM lines could be used to (different systematics)



BAO with other LIM experiments

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- 3% error on BAO peak position at $z=2.5$ with noise and foregrounds for SKA
- This could improve to 1.7% with BAO pixel-reconstruction
- And other IM lines could be used to (different systematics)



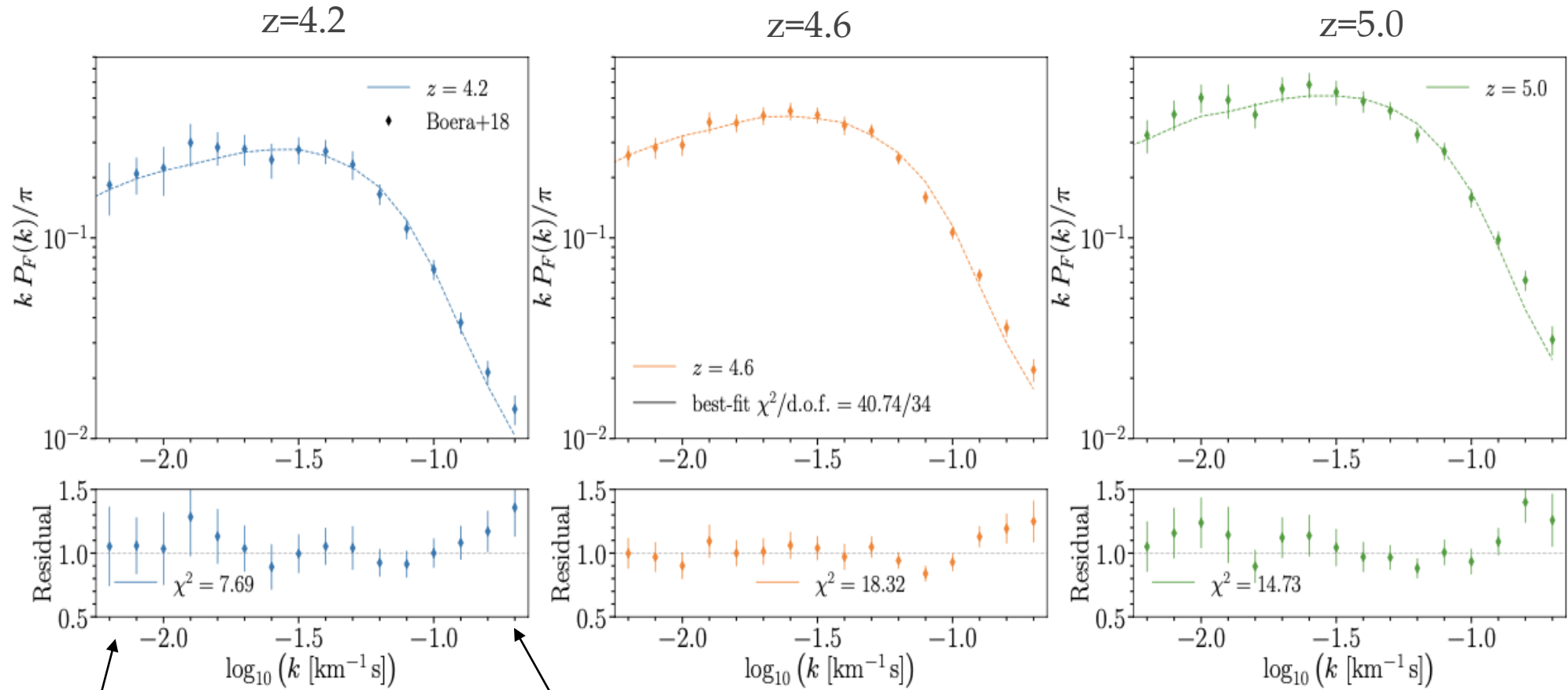
LYMAN- α FOREST

- 1) Modelling relics of reionization
- 2) Warm Dark Matter
- 3) Primordial Magnetic Fields
- 4) Heat injections

The data

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1D FLUX POWER



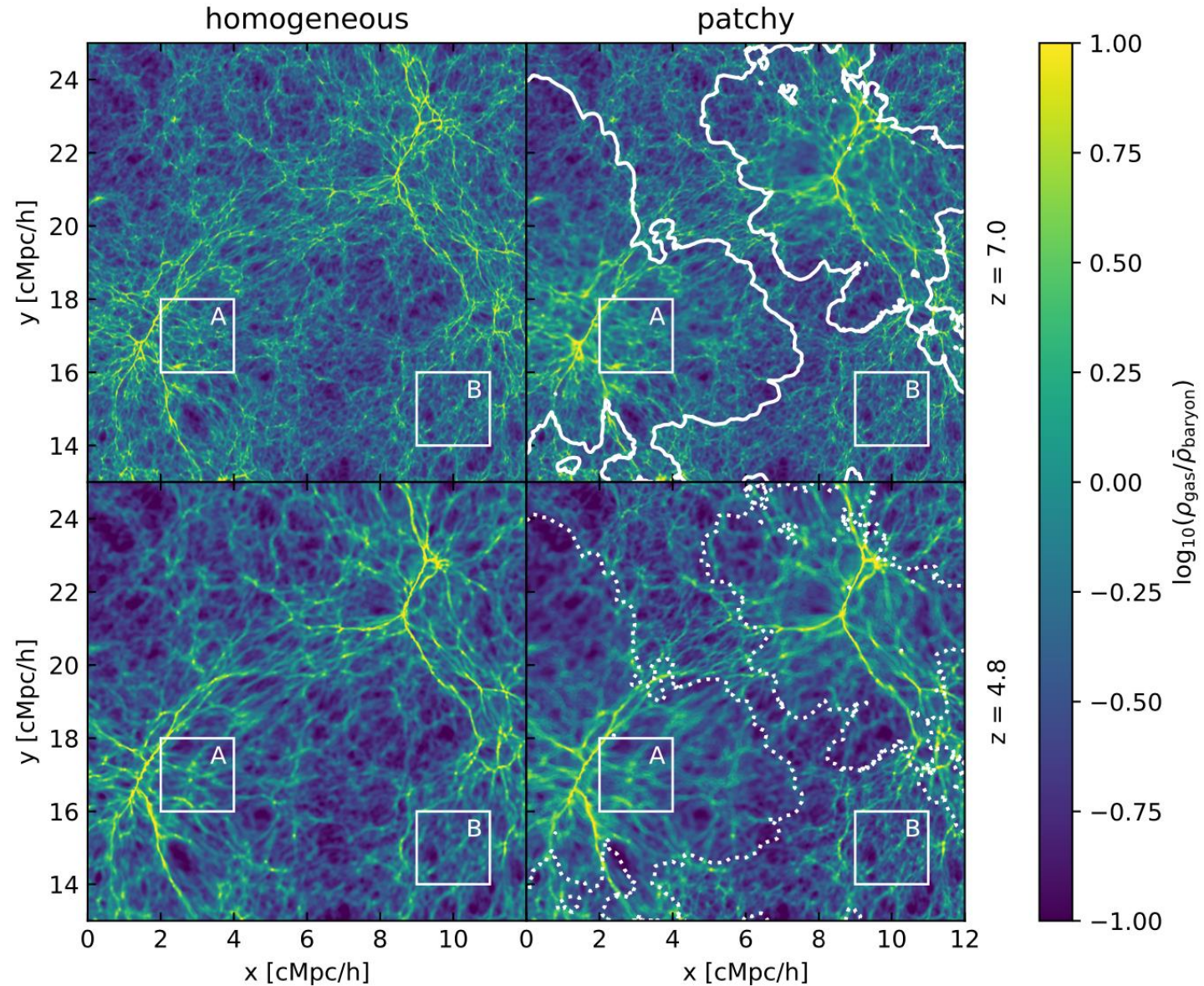
$k=0.4 h/\text{Mpc}$

$k=30 h/\text{Mpc}$

Boera+19, Irsic+23

Theory: patchy reionization

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During reionization $z=7$

*Note:
Reionization ends
at $z=5.4$*

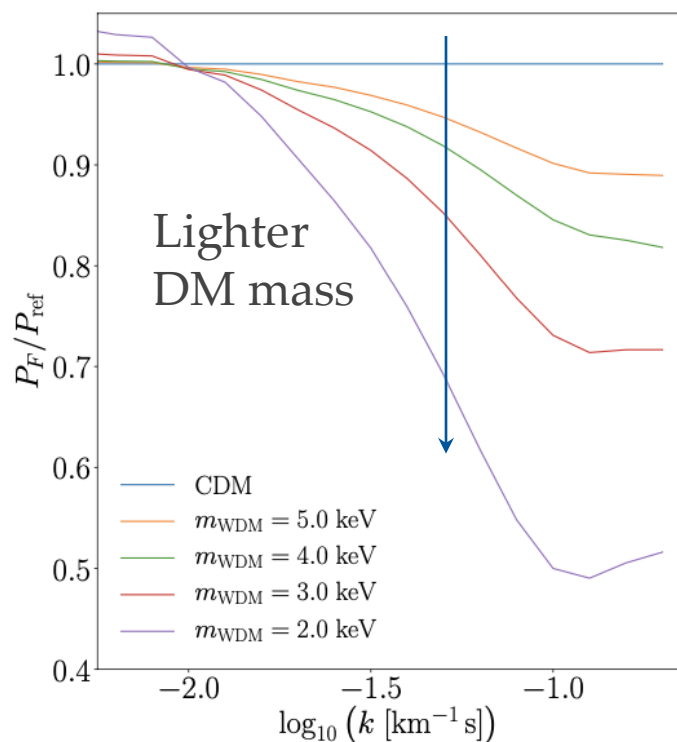
After reionization $z=4.8$

Puchwein+23

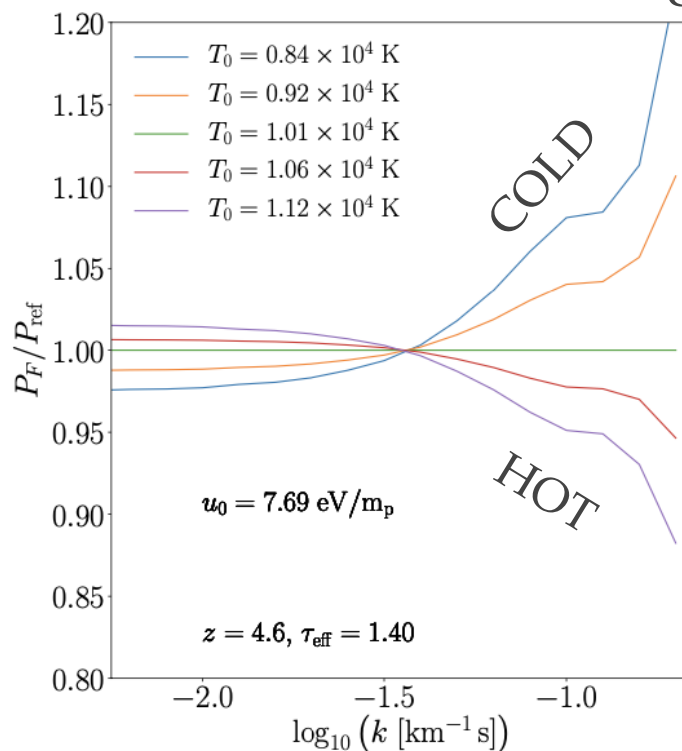
Unveiling Dark Matter free-streaming at the smallest scales with high redshift Lyman-alpha forest

Vid Iršič^{1,2}, Matteo Viel^{3,4,5,6,7}, Martin G. Haehnelt^{1,8}, James S. Bolton⁹, Margherita Molaro⁹, Ewald Puchwein¹⁰, Elisa Boera^{5,6}, George D. Becker¹¹, Prakash Gaikwad¹², Laura C. Keating¹³, Girish Kulkarni¹⁴
¹Kavli Institute for Cosmology, University of Cambridge

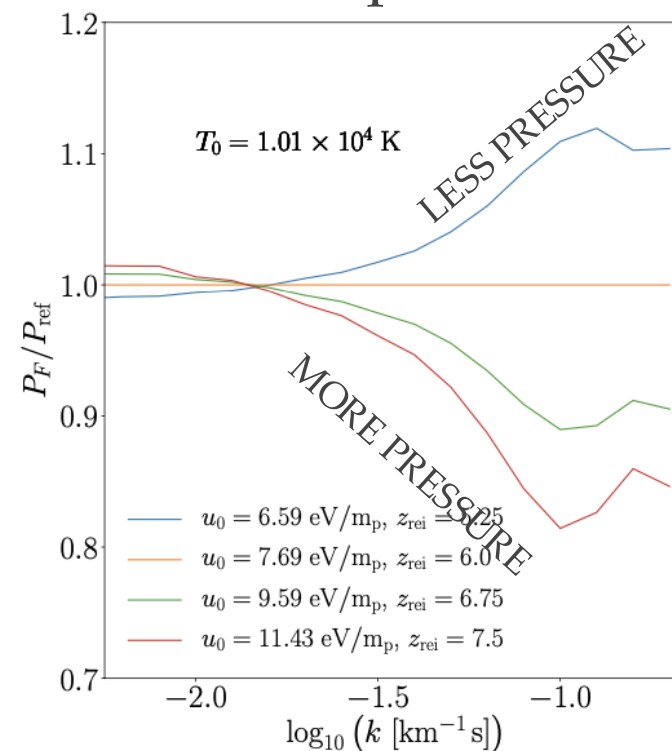
WDM free streaming



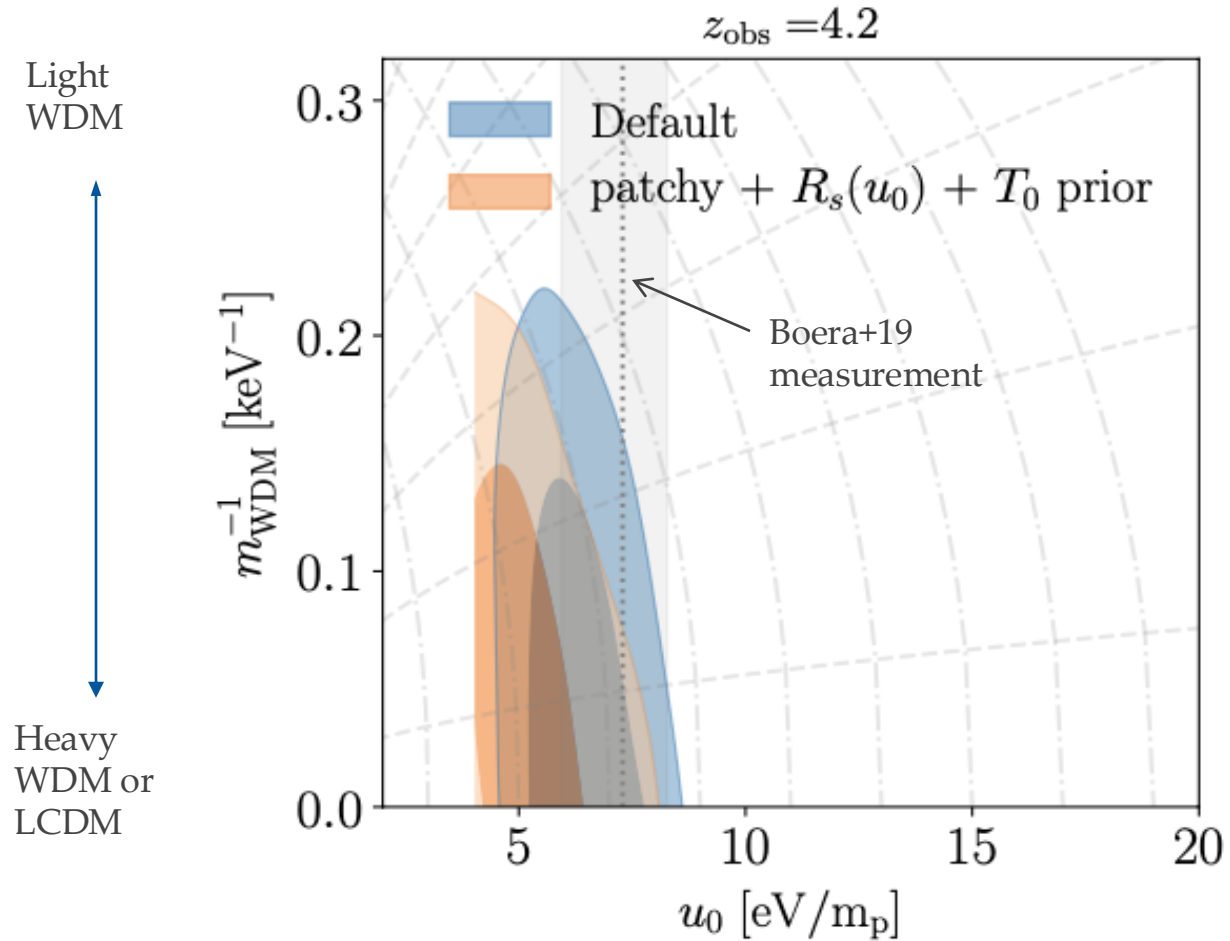
Thermal broadening



Gas pressure



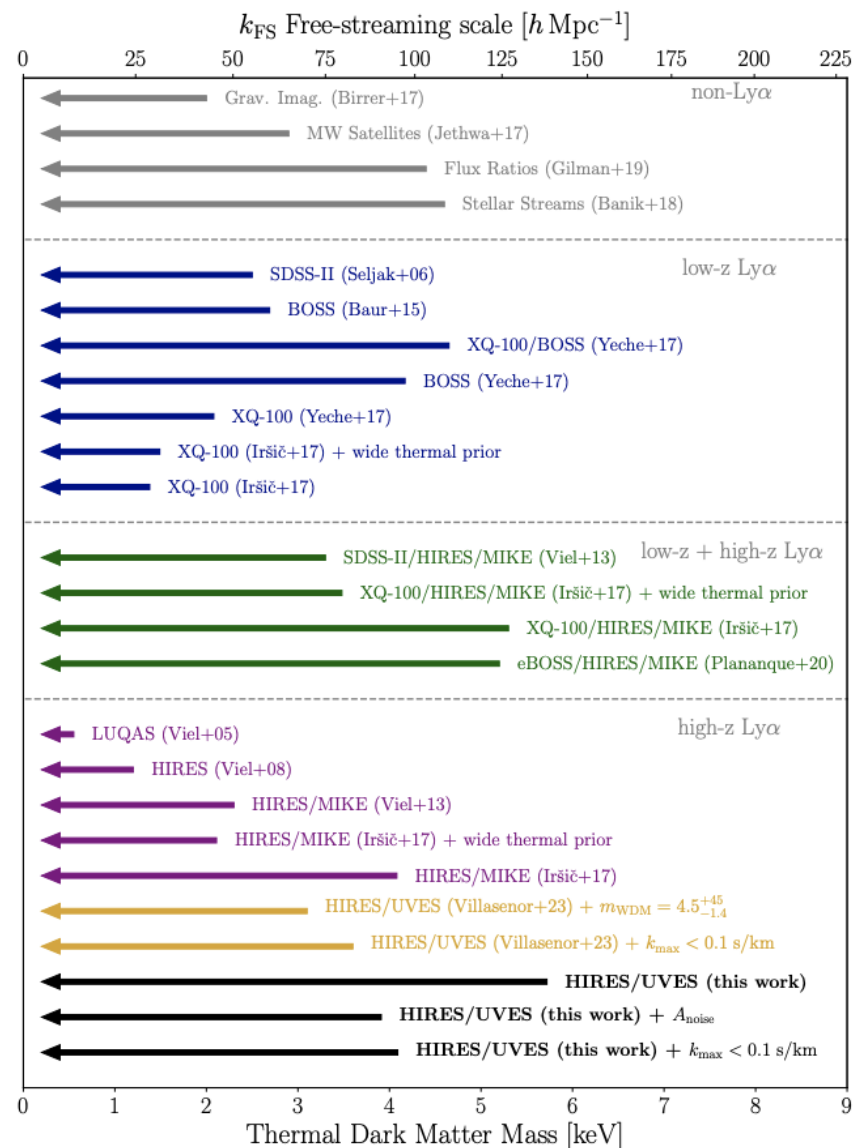
$$u_0(t) = \int_0^t dt \frac{\mathcal{H}}{\bar{\rho}_m} \frac{3k_B}{2\mu} \quad H \text{ is heating rate}$$



Injected heat proxy for GAS pressure

Thermal Warm Dark Matter Constraints

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Tests made:

Cut small scales

Marginalize over data noise

Assume/Remove T_0 priors

Correct for a model dependent resolution

Patchy reionization models

Name	m_{WDM} [keV] (2σ)
Default	> 5.72
$k_{\text{max}} < 0.1 \text{ km}^{-1} \text{ s}$	> 4.10
A_{noise}	> 3.91
T_0 prior	> 5.85
$R_s(u_0)$ mass resolution	> 4.44
patchy reion.	> 5.10
$R_s(u_0) + T_0$ prior	> 4.24
patchy + $R_s(u_0) + T_0$ prior	> 5.90

If f_{WDM} is allowed to vary
for $m_{\text{WDM}}=3 \text{ keV}$
 $f_{\text{WDM}} < 0.5$

Irsic, MV+23-
Garcia-Gallego, Irsic+25

[arXiv:2504.06367](https://arxiv.org/abs/2504.06367)

- Dark photon Dark Matter: simple extension of the SM of particle physics

$$\mathcal{L}_{\gamma A'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F^{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A')^2$$

- Dark photon converts into standard photon when a resonance condition is met

$$E_{A' \rightarrow \gamma} \sim 2.5 \text{ eV} \left(\frac{\epsilon - 14}{0.5} \right)^2 \left(\frac{3}{1 + z_{\text{res}}} \right)^{3/2} \left(\frac{m_{-13}}{0.8} \right)$$

The IGM as a thermometer (low redshift)

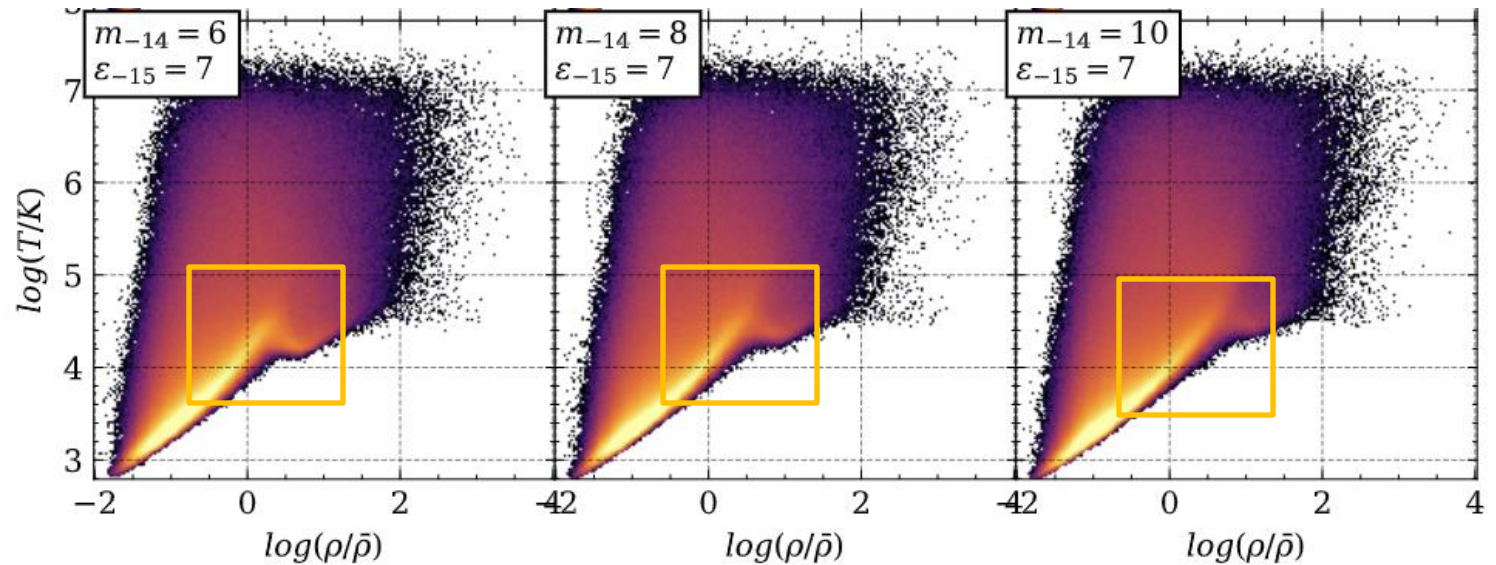
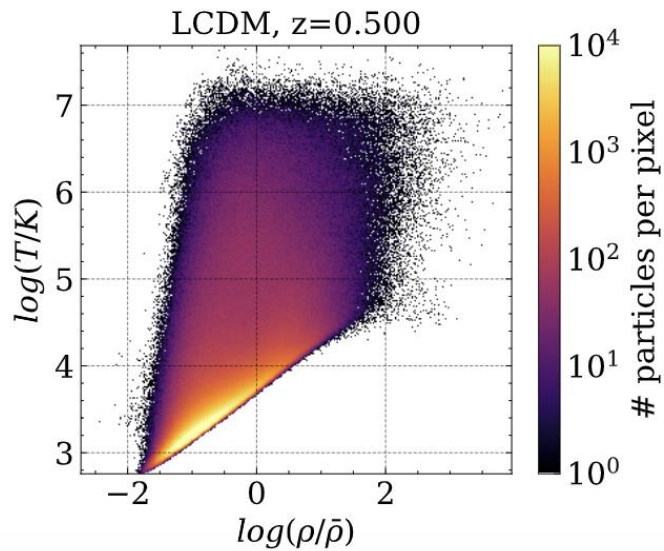
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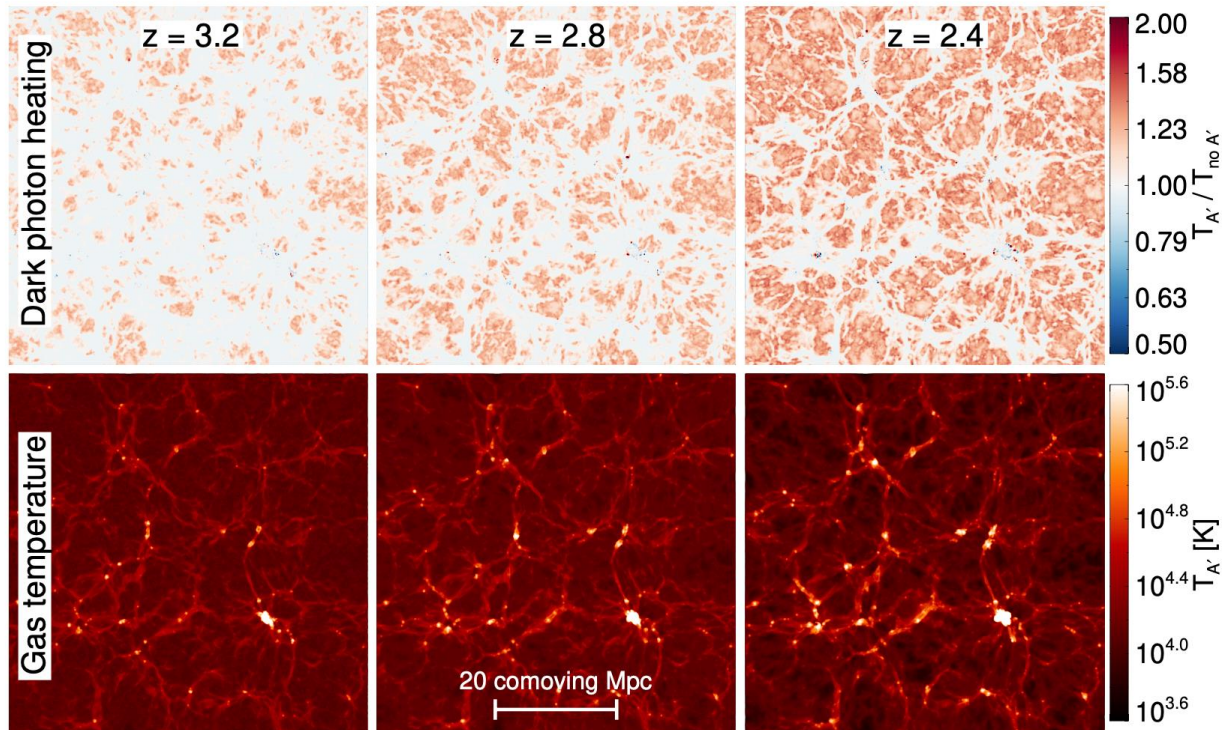
$$E_{A' \rightarrow \gamma} \sim 2.5 \text{ eV} \left(\frac{\epsilon_{-14}}{0.5} \right)^2 \left(\frac{3}{1+z_{\text{res}}} \right)^{3/2} \left(\frac{m_{-13}}{0.8} \right)$$



The IGM as a thermometer (high redshift)

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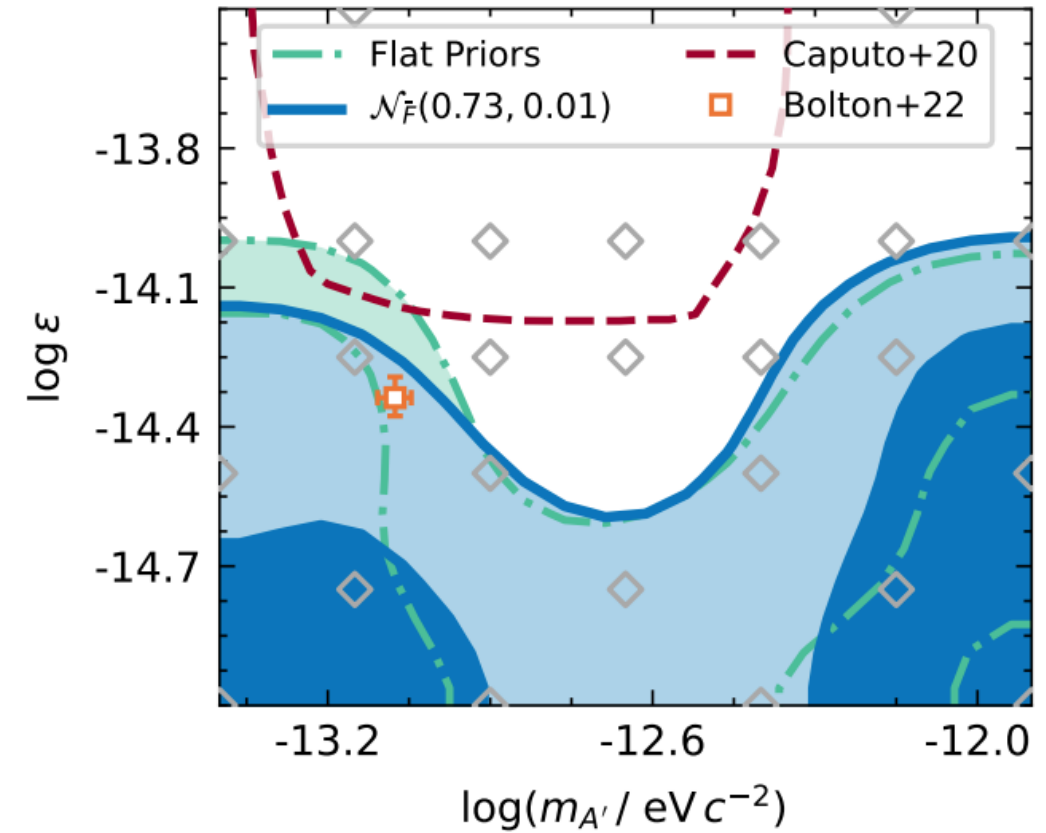
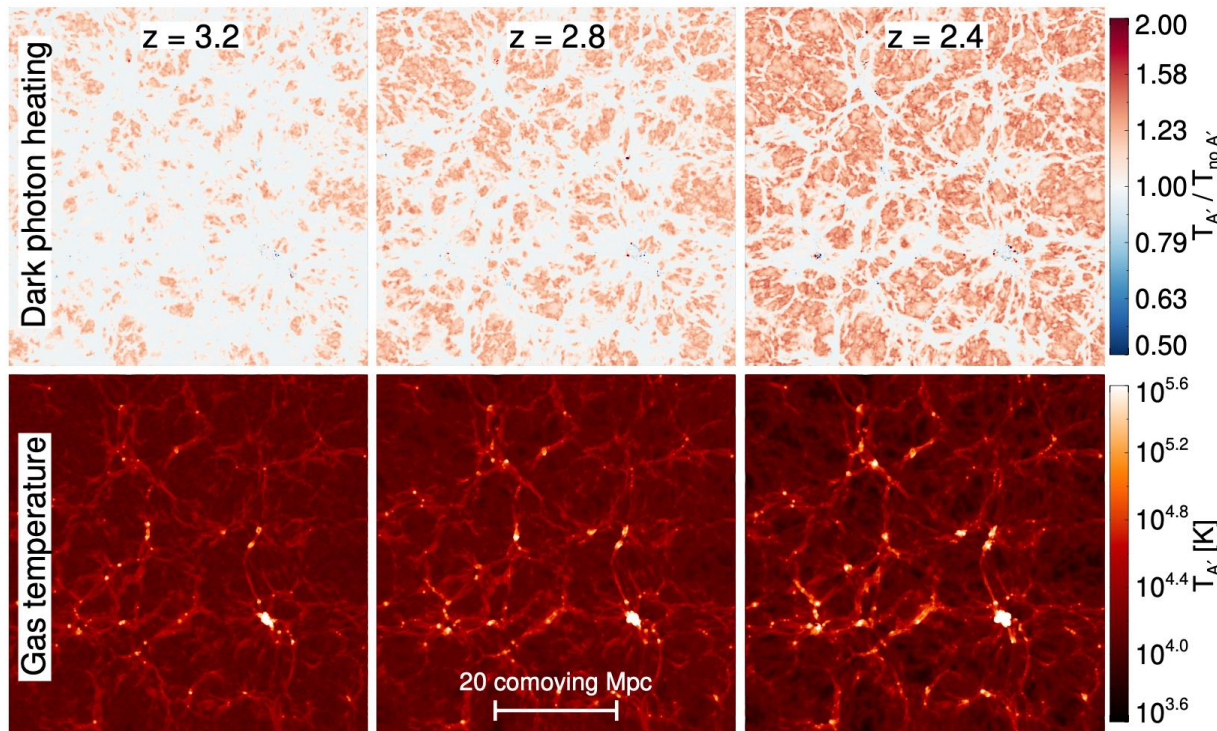
Distinctive heating mechanism happening far away from complex astrophysics



The IGM as a thermometer (high redshift)

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Distinctive heating mechanism happening far away from complex astrophysics



Ideal MHD in the postrecombination Universe

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$$\frac{\partial (\vec{B})}{\partial t} = 0$$

Comoving Magnetic field is conserved

$$\frac{\partial^2 \delta_b}{\partial a^2} + \frac{3}{2} \frac{\partial \delta_b}{a \partial a} = -\frac{\nabla \cdot (\nabla \times \vec{B}) \times \vec{B}}{(4\pi a^3 \rho_b) a^5 H^2} + \frac{\nabla^2 \phi}{(a^2 H)^2}$$

Baryon perturbations driven by magnetic field and gravity

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

Gravity has the usual form

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \frac{3}{2} \frac{\partial \delta_{DM}}{a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

Ideal MHD in the postrecombination Universe

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$$\frac{\partial (\vec{B})}{\partial t} = 0$$

$$\frac{\partial^2 \delta_b}{\partial a^2} + \frac{3}{2} \frac{\partial \delta_b}{a \partial a} = - \frac{\nabla \cdot (\nabla \times \vec{B}) \times \vec{B}}{(4\pi a^3 \rho_b) a^5 H^2} + \frac{\nabla^2 \phi}{(a^2 H)^2}$$

$S_0/a^3 H^2$

$$S_0 = \frac{\nabla \cdot [(\nabla \times \vec{B}) \times \vec{B}]}{4\pi a^3 \rho_b}$$

Key ingredient is the S_0 source term

$$\nabla^2 \phi = \frac{a^2}{2M_{Pl}^2} (\rho_b \delta_b + \rho_{DM} \delta_{DM})$$

$$\frac{\partial^2 \delta_{DM}}{\partial a^2} + \frac{3}{2} \frac{\partial \delta_{DM}}{a \partial a} = \frac{\nabla^2 \phi}{(a^2 H)^2}$$

Ideal MHD in the postrecombination Universe

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$$\begin{aligned} a^2 \frac{\partial^2 \delta_b}{\partial a^2} + a \frac{3}{2} \frac{\partial \delta_b}{\partial a} - \frac{3}{2} \frac{\Omega_b}{\Omega_m (1 + a_{\text{eq}}/a)} \delta_b &= -\frac{S_0}{a^3 H^2} + \frac{3}{2} \frac{\Omega_{\text{DM}}}{\Omega_m (1 + a_{\text{eq}}/a)} \delta_{\text{DM}} \\ a^2 \frac{\partial^2 \delta_{\text{DM}}}{\partial a^2} + a \frac{3}{2} \frac{\partial \delta_{\text{DM}}}{\partial a} - \frac{3}{2} \frac{\Omega_{\text{DM}}}{\Omega_m (1 + a_{\text{eq}}/a)} \delta_{\text{DM}} &= \frac{3}{2} \frac{\Omega_b}{\Omega_m (1 + a_{\text{eq}}/a)} \delta_b. \end{aligned}$$

DM

baryons

Coupled differential equations

$$\delta_b^{\text{PMF}} = -\xi_b(a) \frac{S_0}{a^3 H^2}$$

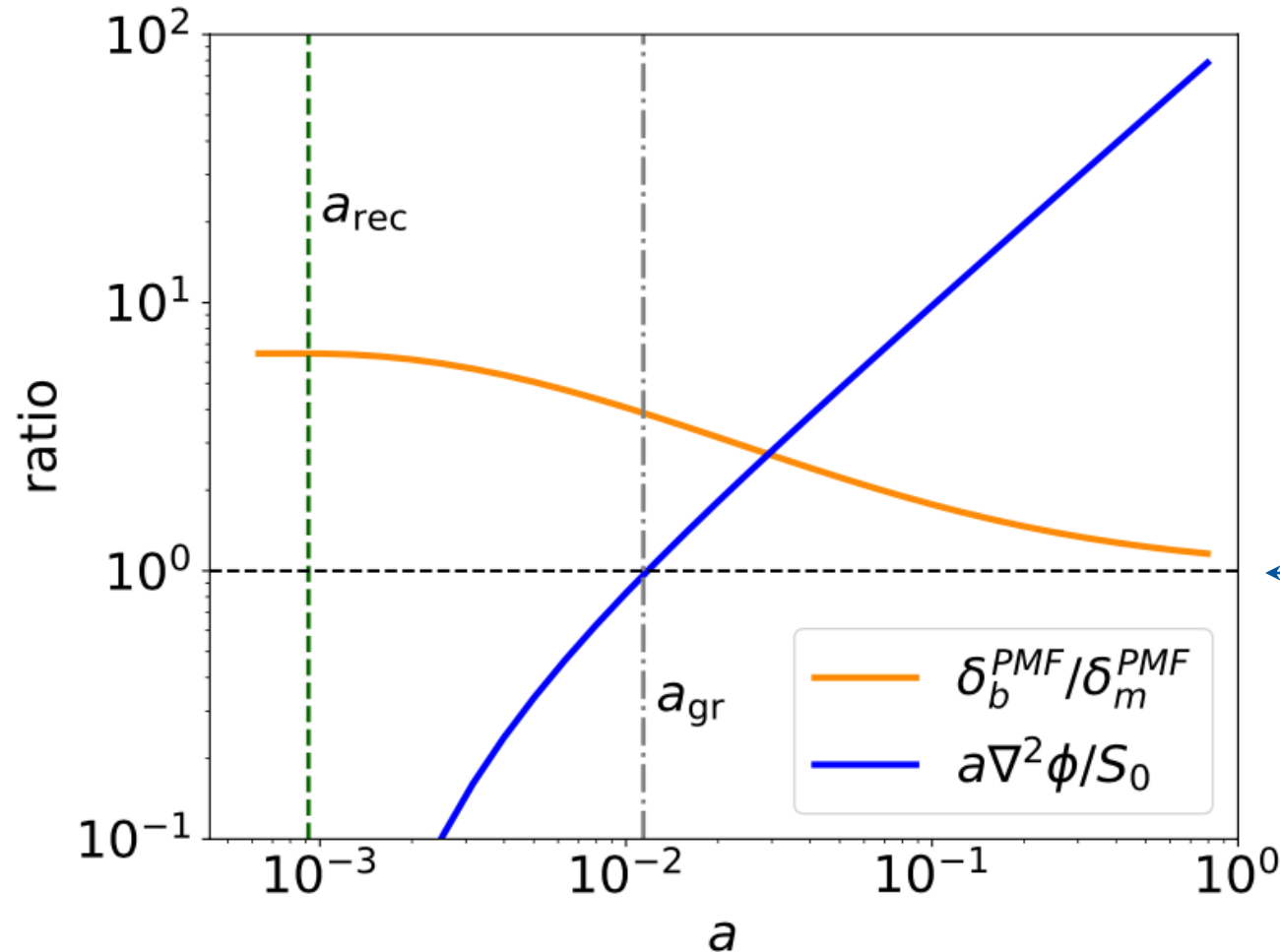
$$\delta_{\text{DM}}^{\text{PMF}} = -\xi_{\text{DM}}(a) \frac{S_0}{a^3 H^2}.$$

$$P_b^{\text{PMF}} \propto P_{S_0}$$

Power spectrum of Lorentz force
For $n_b \sim -3$ (scale invariant) this returns
 $P_{\text{matter}} \sim k$

Ideal MHD in the postrecombination Universe

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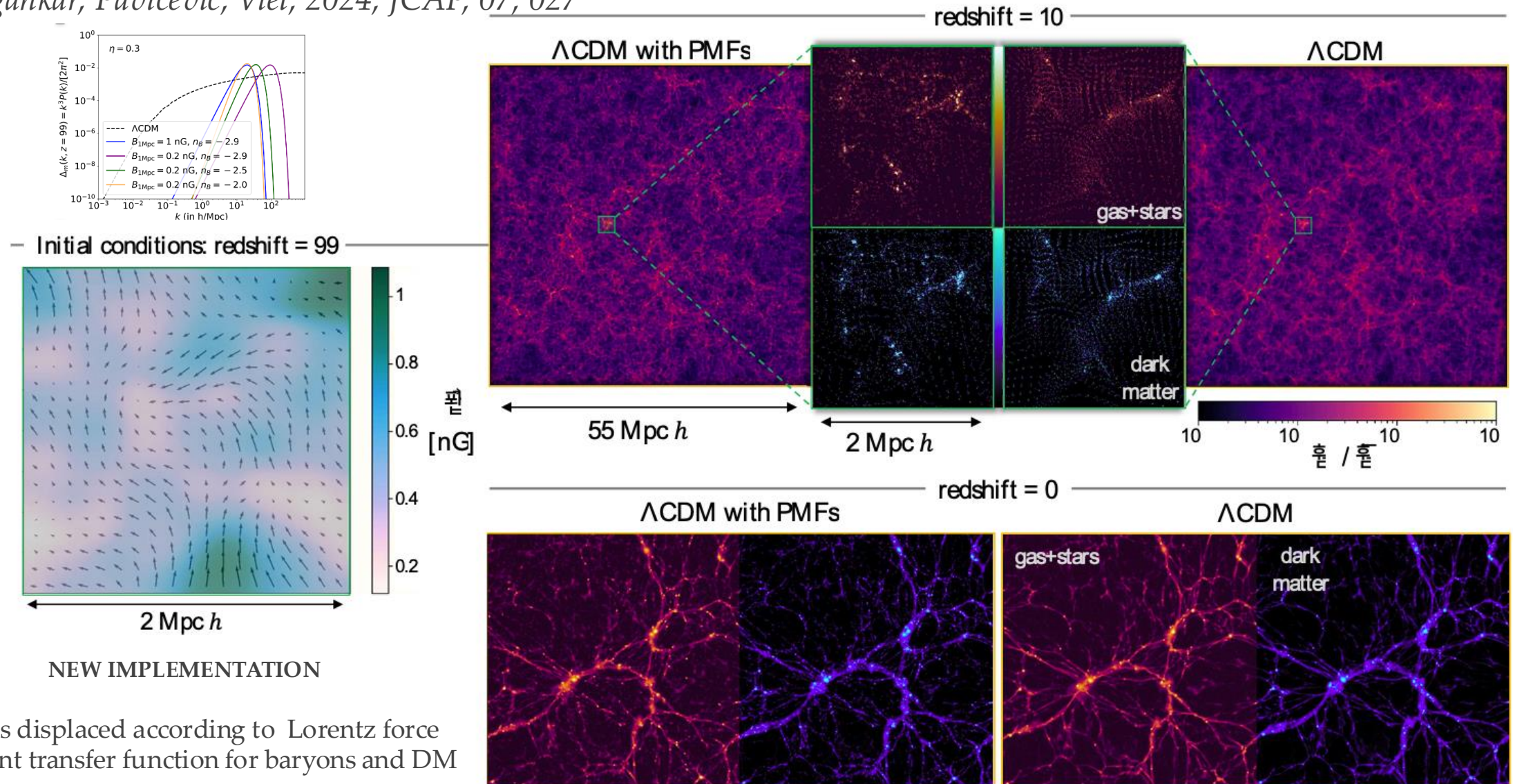


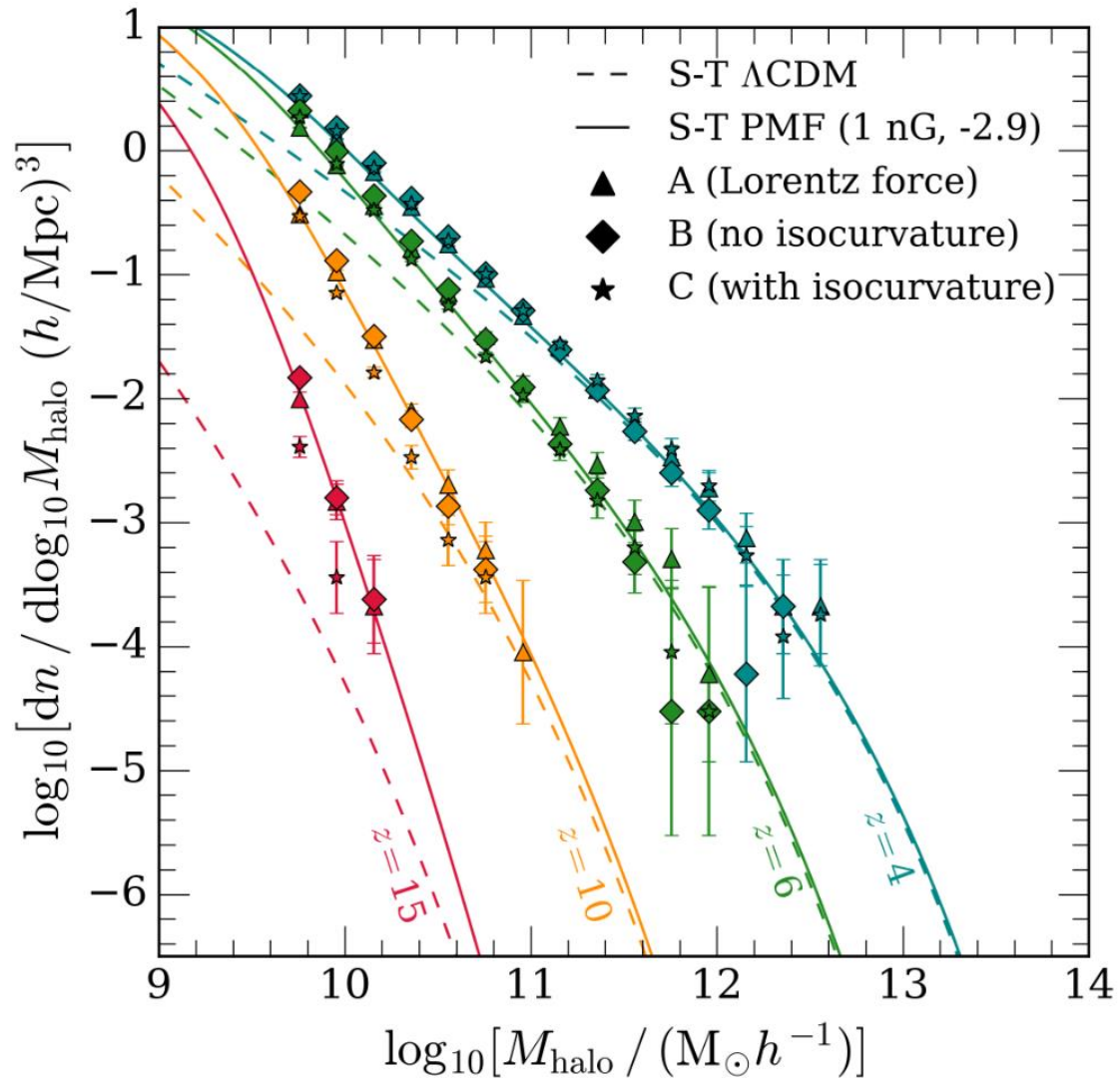
- Ratio of perturbations is equivalent to baryon fraction and starts very high and only now reaches the cosmic mean 0.17 value

Hydro sims with extra PMFs-induced power

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Ralegankar, Pavicevic, Viel, 2024, JCAP, 07, 027

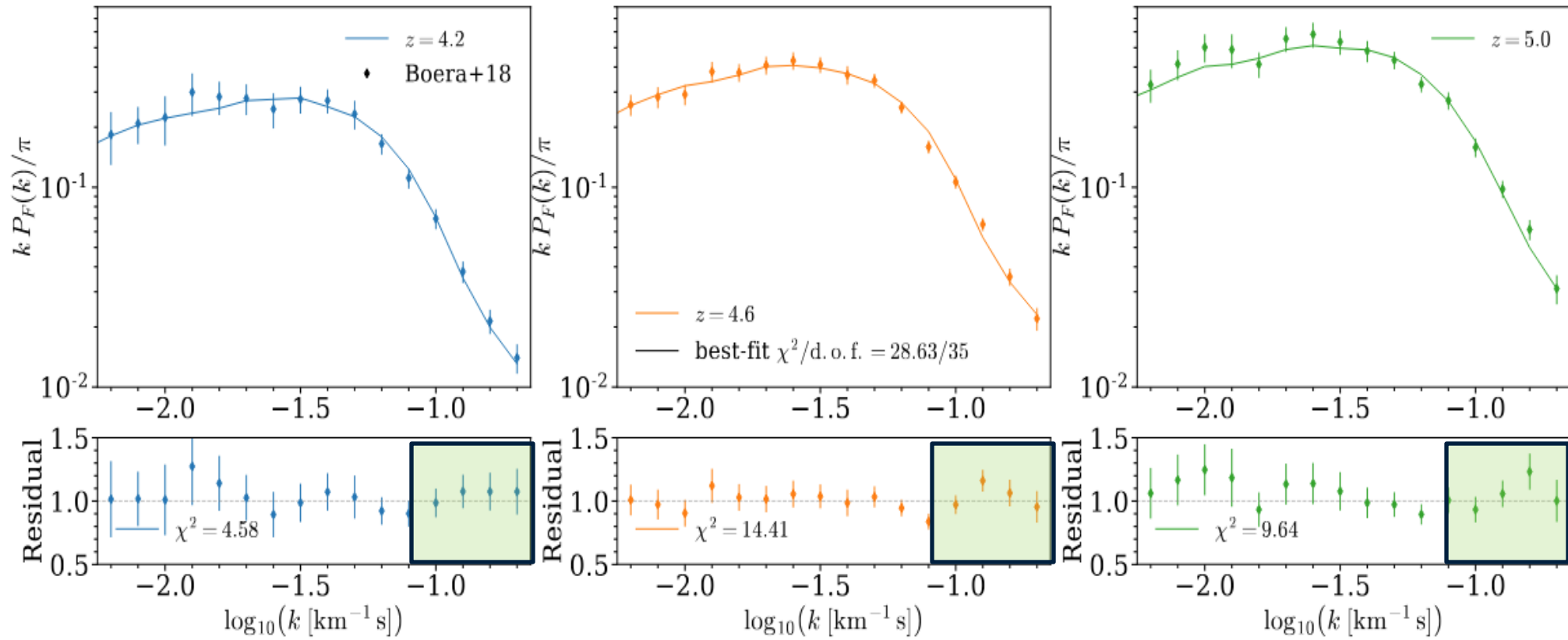




- Extra PMFs power produces more haloes, at “low” mass
- With lower B values (< 1 nG) the enhancement will move to lower masses
- Below 0.05 nG effect is probably too small at any scale

Best fit PMF models

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$$\chi^2_{\Lambda\text{CDM}} = 40.8 \text{ for } 36 \text{ d.o.f.}$$

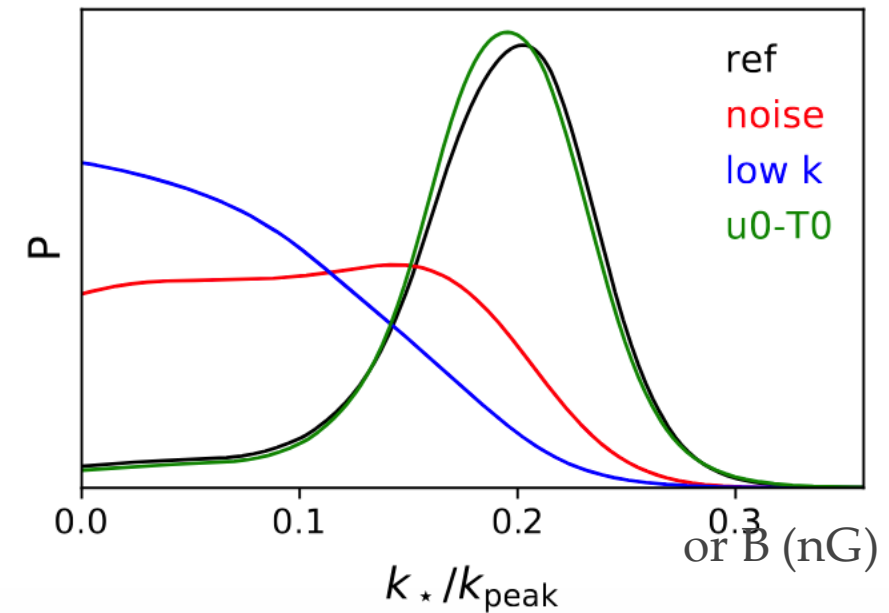
$$\chi^2_{\text{PMF}} = 28.63 \text{ for } 35 \text{ d.o.f.}$$

Constraints on peak position

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$$k_{\text{peak}} = \lambda_D^{-1} \sqrt{\frac{n_B + 5}{2}} \text{ Mpc}^{-1} \quad k_{\star} = 10 \text{ Mpc}^{-1}$$

- Measurement of extra power in the data interpreted in the context of PMFs
- Effectively probing underdense high-redshift regions
- Voids/filaments in the local Universe are also magnetized (see Garg, Durrer, Schober 25)

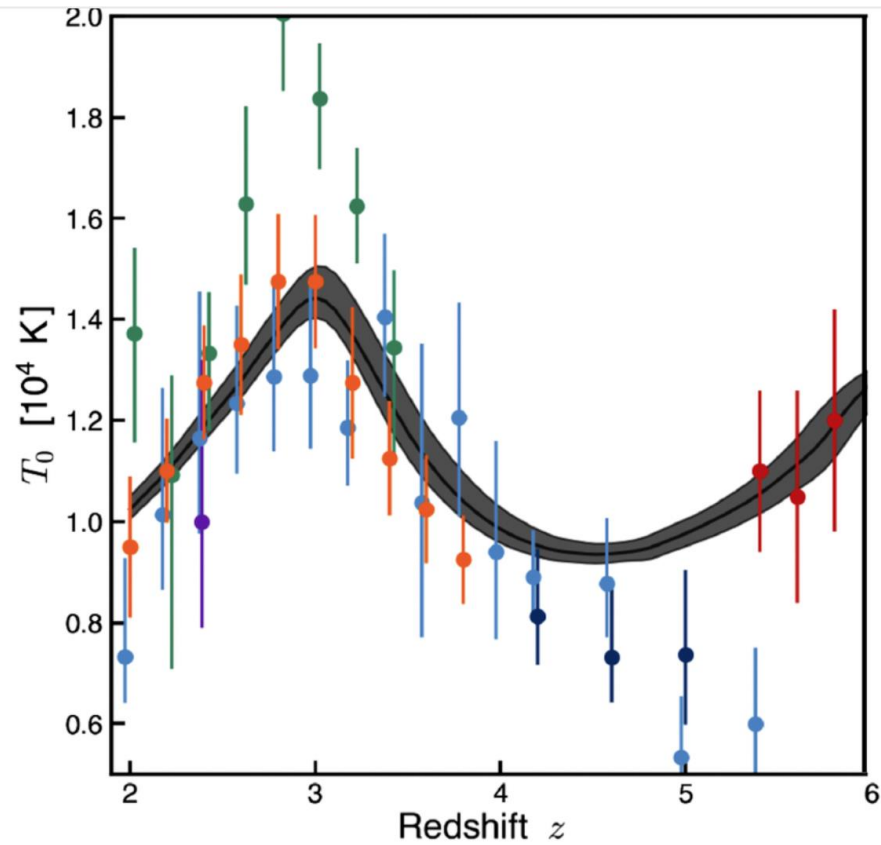


Detection $\rightarrow B = 0.2 \pm 0.05 \text{ nG}$ (1σ)

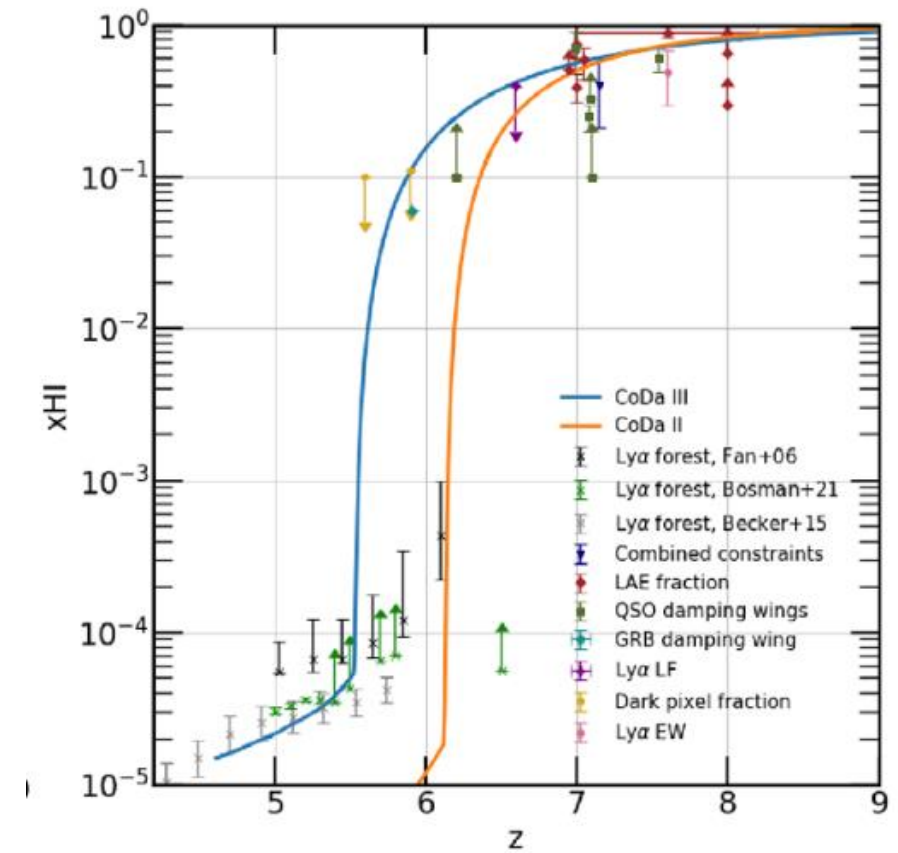
Upper limit $\rightarrow B = 0.3 \text{ nG}$ (3σ)

- Post-reionization Universe: a new place to test structure (and galaxy) formation and probe fundamental physics
- Access to relatively small scales $k \sim 1 \ h/\text{Mpc}$ with IM and $k \sim 30 \ h/\text{Mpc}$ with forest
- With intensity mapping 1D radial BAO will constrain geometry
- Power spectrum will constrain growth (and thus neutrino masses – Autieri's talk)
- Warm dark matter constraints $m > 5.72 \text{ keV}$ – constraints on Cold + Warm DM models too
- Hint of extra power in the data is well fitted with PMF at 0.2 nG , robust upper limit 0.3 nG

Thermal state of baryons at mean density



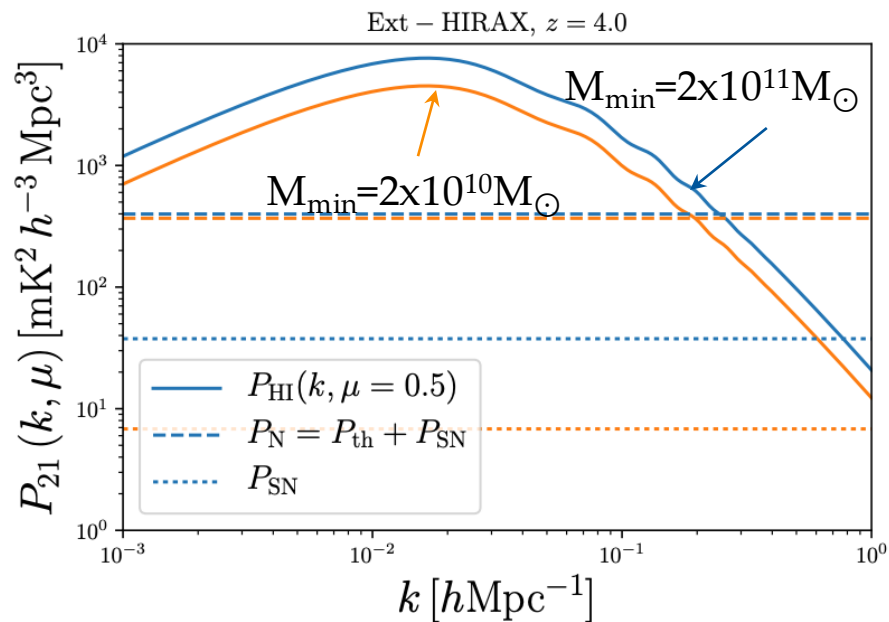
Neutral fraction evolution



Post-reionization cosmology with IM experiments

Matteo Viel

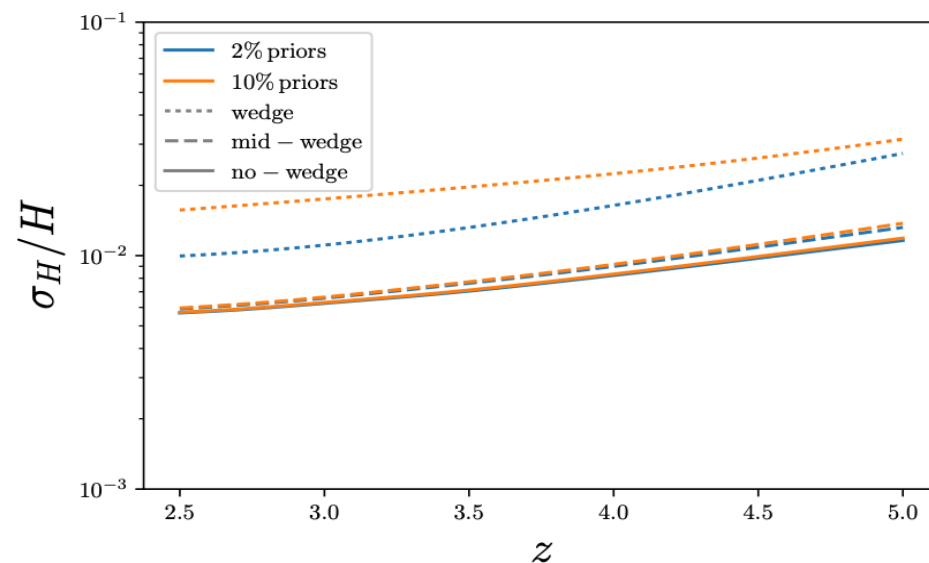
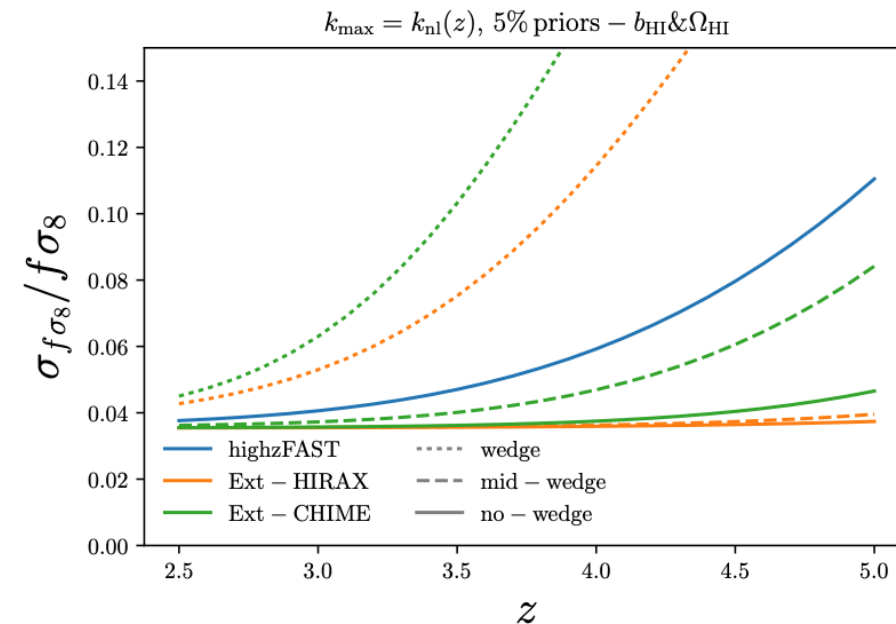
Fisher matrix analysis of future "extended" HIRAX experiment



Considering also the wedge, beam, noise
Shot noise of the tracer

DYNAMICS

GEOMETRY



Modelling LIM power to linear order

Matteo Viel

$$P_{21\text{ cm}}(k, \mu, z) = \bar{T}_b(z)^2 [(b_{\text{HI}}(z) + f(z)\mu^2)^2 P_{\text{m}}(k, z) + P_{\text{SN}}(z)],$$

Linear power (cosmology)

Brightness HI temperature or other lines

$$\bar{T}_b(z) = 189h \left(\frac{H_0(1+z)^2}{H(z)} \right) \Omega_{\text{HI}}(z) \text{ mK}$$

Amount of HI

$$\Omega_{\text{HI}}(z) = \frac{1}{\rho_c^0} \int_0^\infty n(M, z) M_{\text{HI}}(M, z) dM$$

HI bias

$$b_{\text{HI}}(z) = \frac{1}{\rho_c^0 \Omega_{\text{HI}}(z)} \int_0^\infty n(M, z) b(M, z) M_{\text{HI}}(M, z) dM$$

Shot-Noise power spectrum

$$P_{\text{SN}}(z) = \frac{1}{(\rho_c^0 \Omega_{\text{HI}}(z))^2} \int_0^\infty n(M, z) M_{\text{HI}}^2(M, z) dM$$

New physics from $P(k)$ or $n(M)$

Halo mass function (cosmology)

Amount of HI in each DM halo (astrophysics)

Modelling of the LIM power with the halo model

Matteo Viel

- Halo models important for reaching small scales
- Can be easily extended to any IM line
- Profile must be specified
- M_{HI} and Ω_{HI} from sims or from observed HI mass function or DLAs

$$P_{\text{HI}}(k, z) = P_{\text{HI},1\text{h}}(k) + P_{\text{HI},2\text{h}}(k) \quad P_{\text{HI}}^{\text{SN}}(z) = \lim_{k \rightarrow 0} P_{1\text{h},\text{HI}}(k, z) :$$
$$P_{\text{HI},1\text{h}}(k, z) = \frac{1}{(\rho_{\text{c}}^0 \Omega_{\text{HI}}(z))^2} \int_0^\infty dM n(M, z) M_{\text{HI}}^2(M, z) |u_{\text{HI}}(k|M, z)|^2$$
$$P_{\text{HI},2\text{h}}(k, z) = \frac{P_{\text{lin}}(k, z)}{(\rho_{\text{c}}^0 \Omega_{\text{HI}}(z))^2} \left[\int_0^\infty dM n(M, z) b(M, z) M_{\text{HI}}(M, z) |u_{\text{HI}}(k|M, z)| \right]^2$$

COSMOLOGY

Modelling of the LIM power with the halo model

Matteo Viel

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- Can be easily extended to any IM line
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$$P_{\text{HI},2\text{h}}(k, z) = \frac{P_{\text{lin}}(k, z)}{(\rho_{\text{c}}^0 \Omega_{\text{HI}}(z))^2} \left[\int_0^\infty dM n(M, z) b(M, z) M_{\text{HI}}(M, z) |u_{\text{HI}}(k|M, z)| \right]^2$$

COSMOLOGY

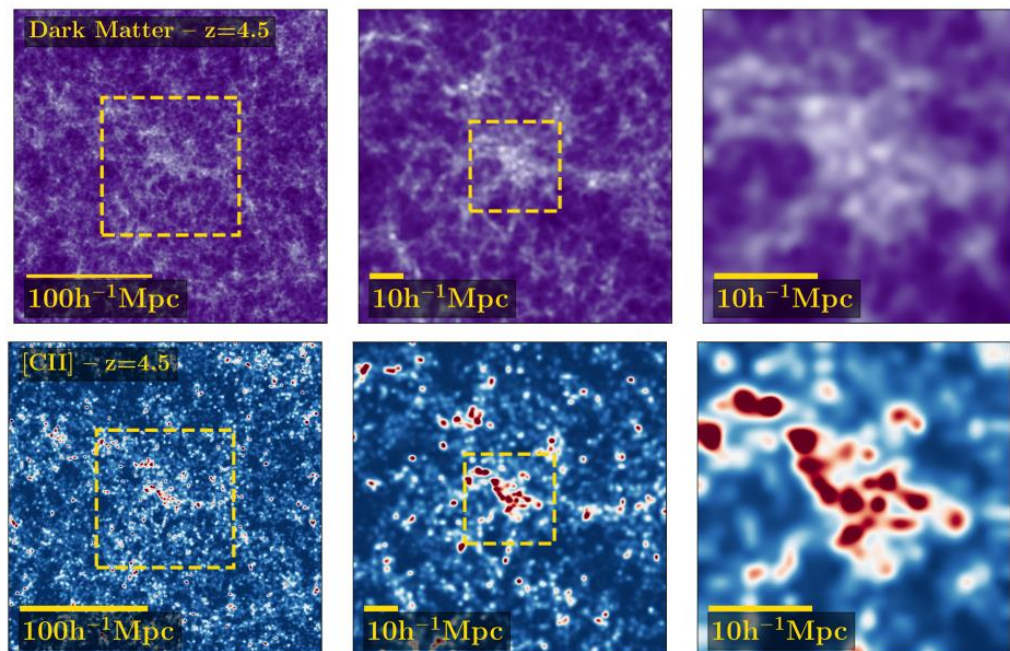
ASTROPHYSICS OF THE HALOES

$$M_{\text{HI}}(M, z) = \alpha f_{\text{H,c}} M \left(\frac{M}{10^{11} h^{-1} M_{\odot}} \right)^{\beta} \exp \left[- \left(\frac{v_{\text{c}0}}{v_{\text{c}}(M, z)} \right)^3 \right] \quad \rho_{\text{HI}}(r; M, z) = \rho_0 \exp(-r/r_{\text{s,HI}})$$

Physically-rich modelling: involves a set of parameters that are calibrated on sims to fit observations

New modelling of LIM power with the halo model

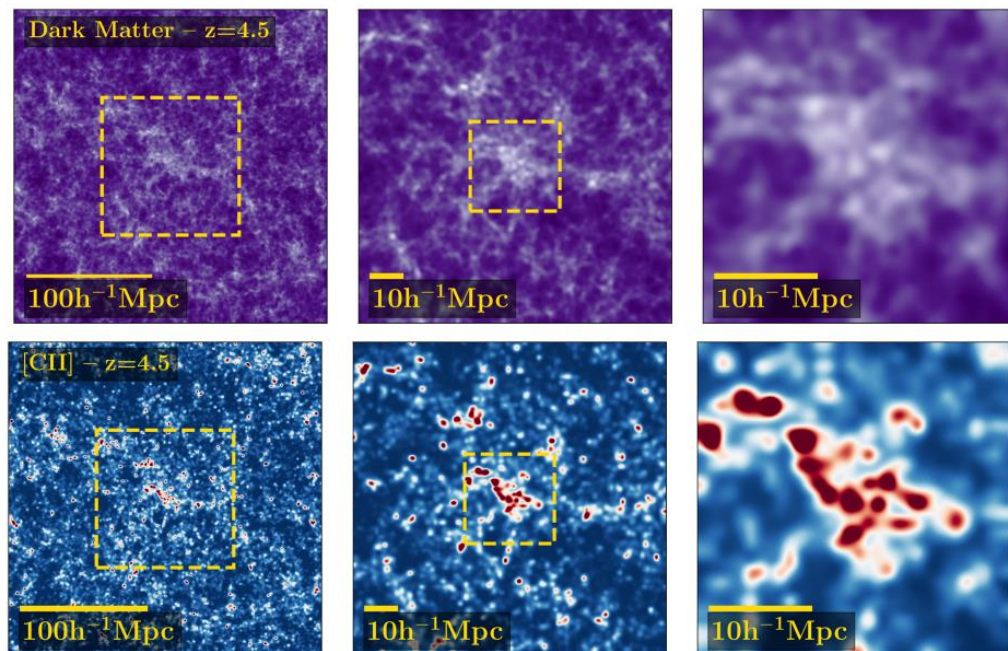
Matteo Viel



- Alcock-Paczynski parameters and BAO wiggles
- State-of-the-art treatment of (non) Poisson shot noise
- Bias of the different lines [CII] and CO
- EFT inspired perturbation theory at 1-loop
- Comparison with large scale/high res. (DM only) mocks

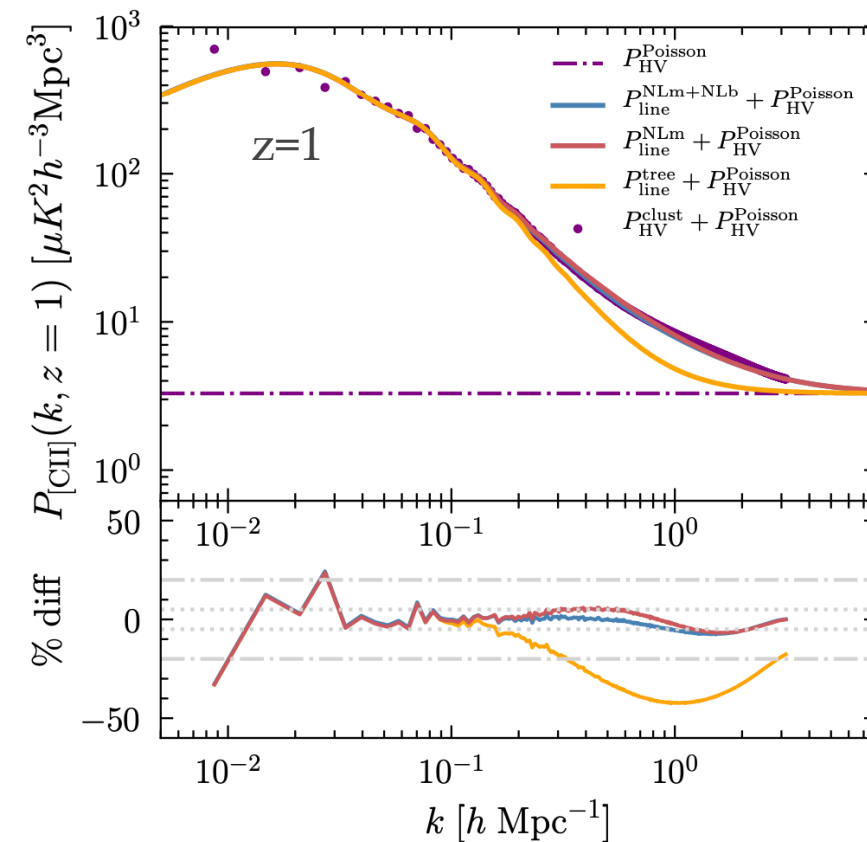
New modelling of LIM power with the halo model

Matteo Viel



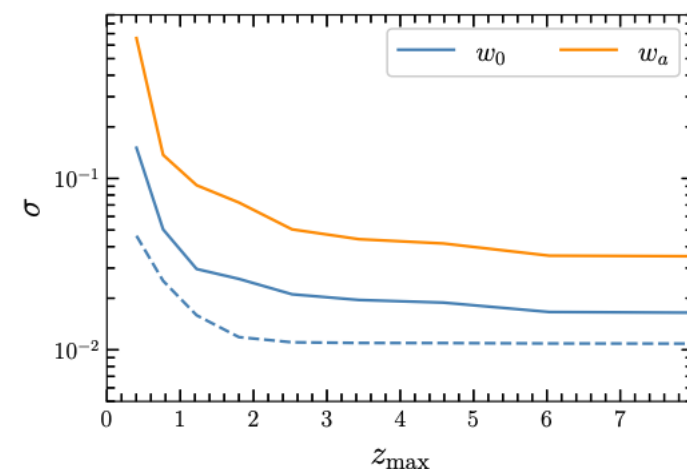
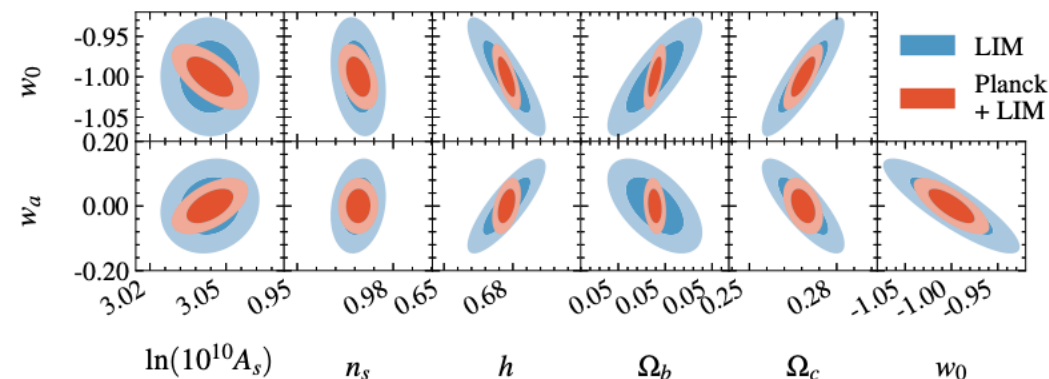
- Alcock-Paczynski parameters and BAO wiggles
- State-of-the-art treatment of (non) Poisson shot noise
- Bias of the different lines [CII] and CO
- EFT inspired perturbation theory at 1-loop
- Comparison with large scale/high res. (DM only) mocks
- Range of validity $k \sim 1 h/\text{Mpc}$ at $z=1$ (5% agreement)

Different treatment of non linearities in matter and bias



- Testing GR and DE with LIM (Horndeski, Brans-Dicke, early dark energy models)
- Fisher matrix analysis for CO and [CII] on $P(k)$ including modelling of the interlopers; scatter in $L(M)$; shot noise; instrumental noise
- Effectively a linear model, which is sensitive to geometry and dynamics

CPL parameterization



Beyond P(k): Voxel Intensity Distribution

Matteo Viel

cosmology

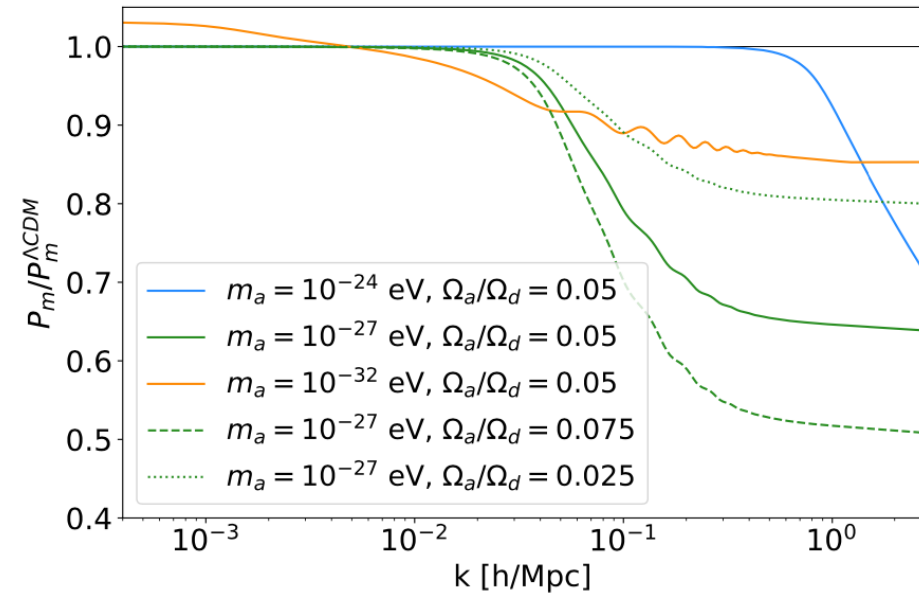
$$\mathcal{T}^2(k) \equiv \frac{P_{\text{nCDM}}(k)}{P_{\text{CDM}}(k)} = \begin{cases} 1 & \text{if } k \leq k_{\text{cut}} \\ \left(\frac{k}{k_{\text{cut}}}\right)^{-n} & \text{if } k > k_{\text{cut}} \end{cases}$$

astrophysics

$$\frac{L_{\text{CO}}}{L_{\odot}}(M) = 4.9 \times 10^{-5} \frac{C}{(M/M_*)^A + (M/M_*)^B}$$

$$\mathcal{P}(T) = \sum_{N=0}^{\infty} \mathcal{P}_N(T) \mathcal{P}(N)$$

- Pheno model that captures axion + cold DM
- COMAP-Y5 experiment ($z=2.9$)
- Monopole of the power spectrum + Voxel intensity distribution (VID) – this is important to capture non gaussian nature of signal

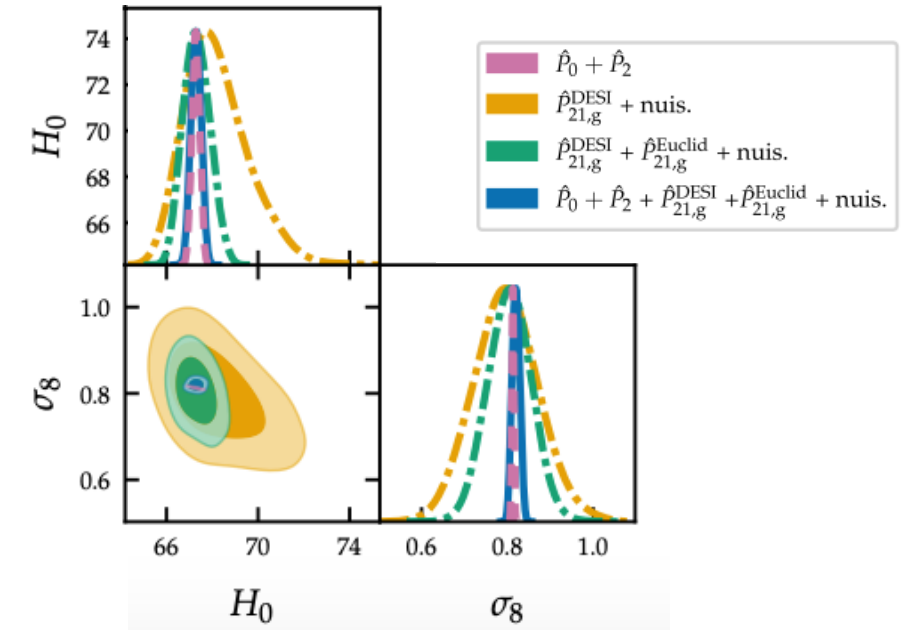


Modelling the 21cm cross-power: SKA-mid X DESI/Euclid

Matteo Viel

$$P_g(z, k, \mu) = \left(b_g(z) + f(z) \mu^2 \right)^2 P_m(z, k) + \frac{1}{\bar{n}_g(z)}$$

- Cross-correlation with spectroscopic samples
- Nuisance parameters to bracket instrumental and astrophysical uncertainties
- Very constraining (similar to auto-correlation)

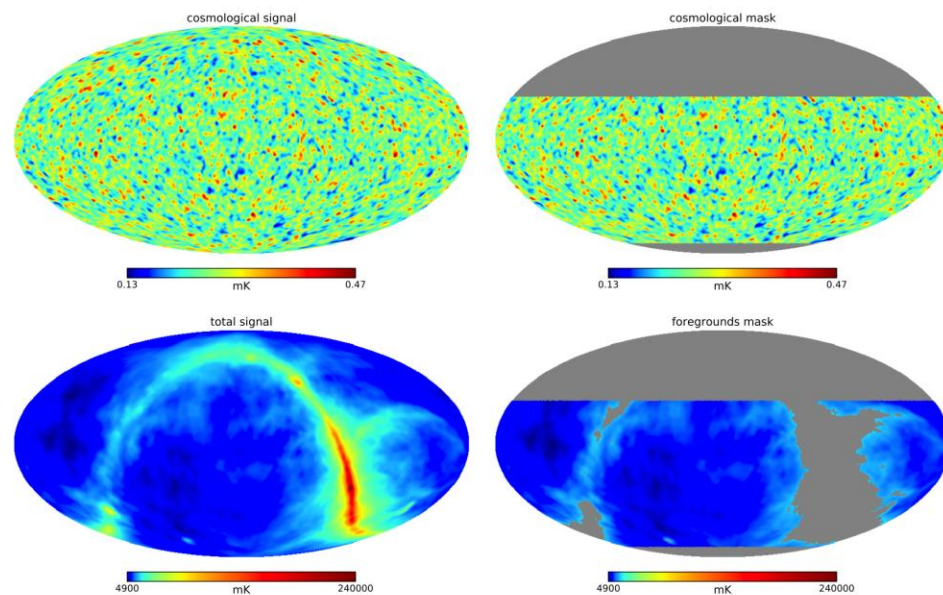


Parameter	$\hat{P}_0 + \hat{P}_2$	$\hat{P}_{21,g}^{\text{DESI}}$	$\hat{P}_{21,g}^{\text{DESI}} + \text{nuis.}$	$\hat{P}_{21,g}^{\text{Euclid}}$	$\hat{P}_{21,g}^{\text{Euclid}} + \text{nuis.}$	$\hat{P}_{21,g}^{\text{DESI}} + \hat{P}_{21,g}^{\text{Euclid}} + \text{nuis.}$	$\hat{P}_0 + \hat{P}_2 + \hat{P}_{21,g}^{\text{DESI}} + \hat{P}_{21,g}^{\text{Euclid}} + \text{nuis.}$
H_0	0.25%	0.69%	1.96%	0.49%	1.07%	0.87%	0.33%

Baryonic Acoustic Oscillations in 21cm IM

Matteo Viel

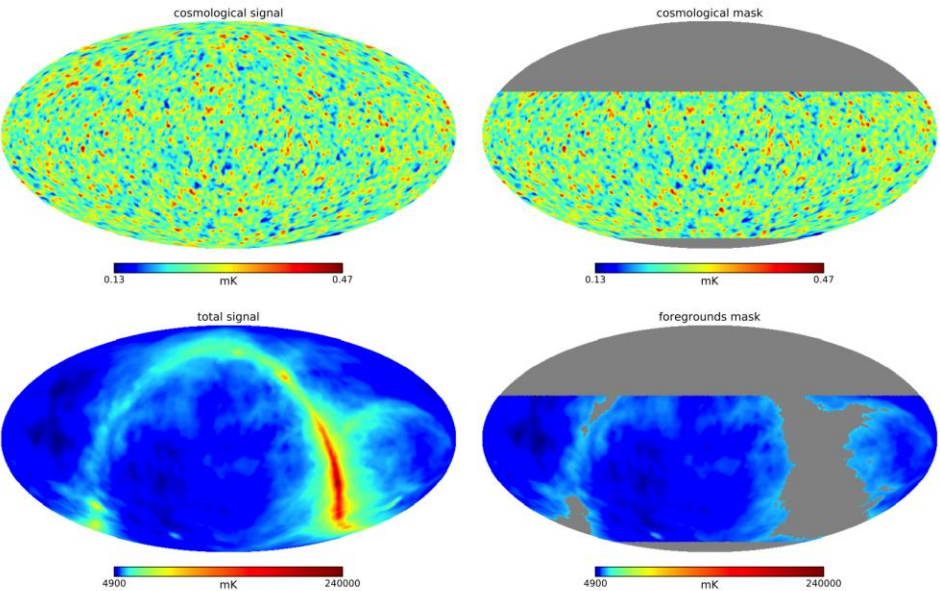
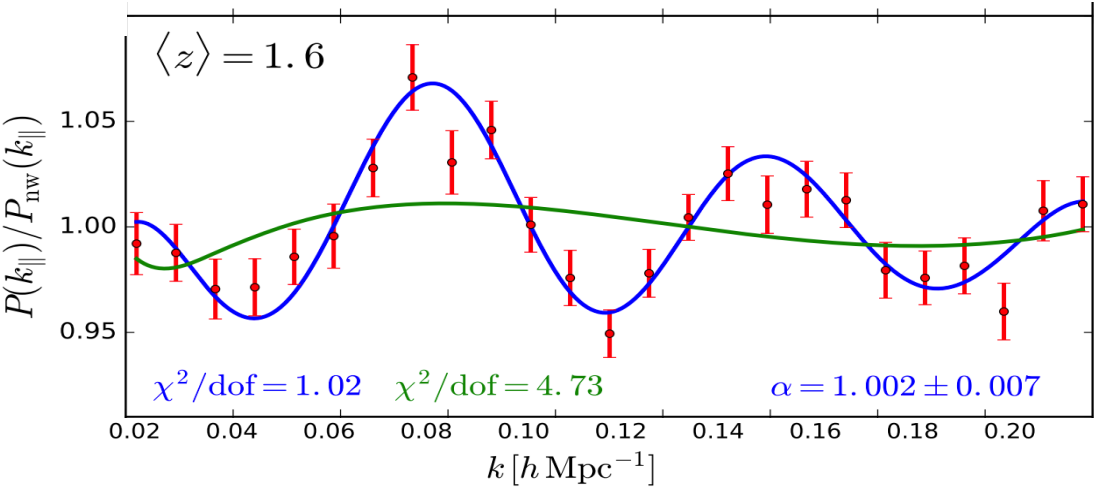
- SKA estimate: $H(z)$ measured at sub-percent level up to $z=2.5$
- Made with mask, and foregrounds removal
- And realistic treatment of instrument noise



Baryonic Acoustic Oscillations in 21cm IM

Matteo Viel

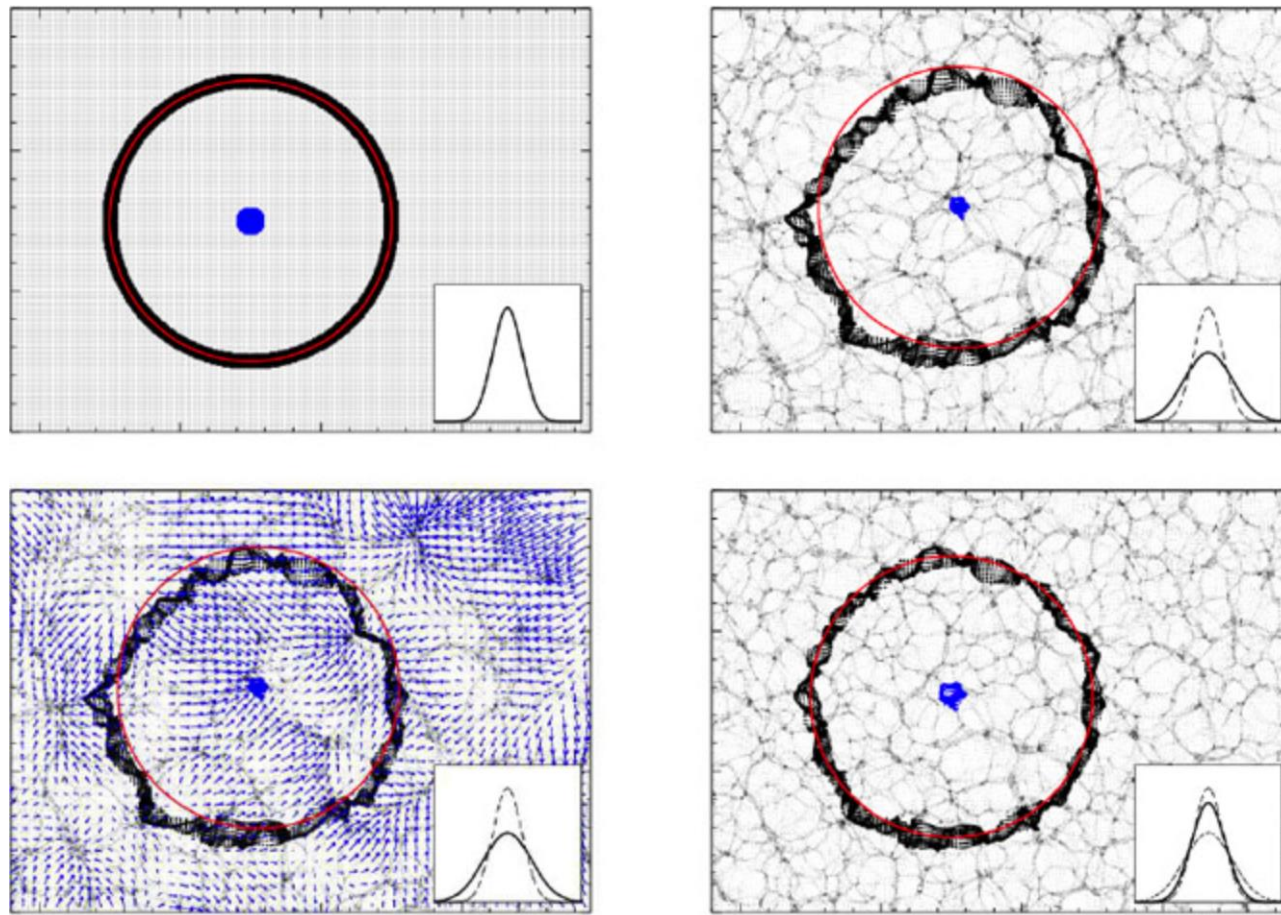
- SKA estimate: $H(z)$ measured at sub-percent level up to $z=2.5$
- Made with mask, and foregrounds removal
- And realistic treatment of instrument noise



z range	$\langle z \rangle$	mask	σ_α	σ_α	σ_α
			(C)	(C+N)	(C+N+FG)
[0.36-0.75]	0.6	no	1.008 ± 0.016	1.008 ± 0.016	1.007 ± 0.016
		yes	1.006 ± 0.020	1.006 ± 0.021	1.006 ± 0.024
[0.75-1.26]	1.0	no	0.996 ± 0.010	0.997 ± 0.011	0.996 ± 0.011
		yes	0.997 ± 0.012	0.997 ± 0.013	0.998 ± 0.015
[1.26-1.98]	1.6	no	1.001 ± 0.011	1.004 ± 0.014	1.003 ± 0.014
		yes	1.000 ± 0.013	1.003 ± 0.016	1.004 ± 0.019
[1.98-3.05]	2.5	no	1.004 ± 0.013	1.003 ± 0.021	1.000 ± 0.021
		yes	1.004 ± 0.016	1.002 ± 0.026	1.002 ± 0.031

Reconstruction of the BAO peak with pixels

Matteo Viel



- For galaxies, undoing non-linearities with Zeldovich to sharpen the BAO peak
- How would this work for pixels?
- Even less computationally demanding, pixels are moved... and local density estimates (grid based) are for free

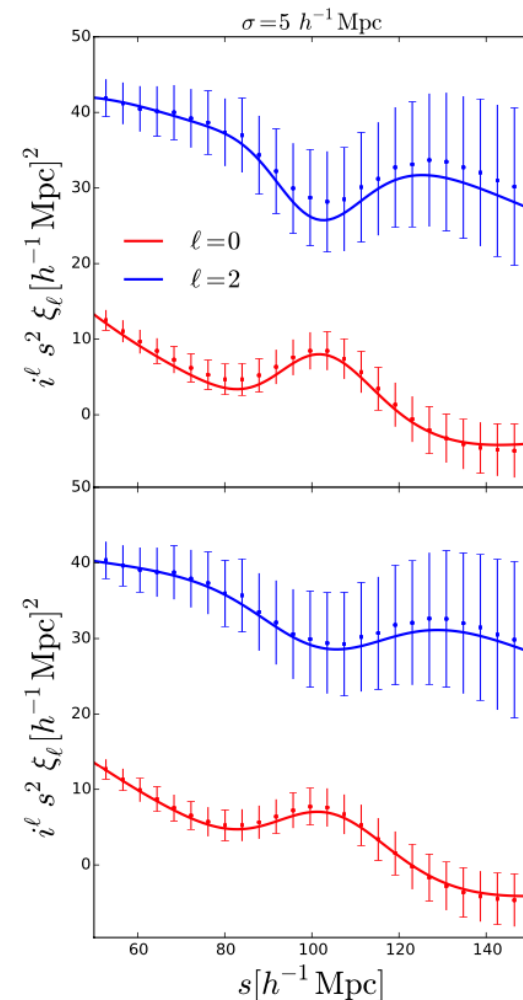
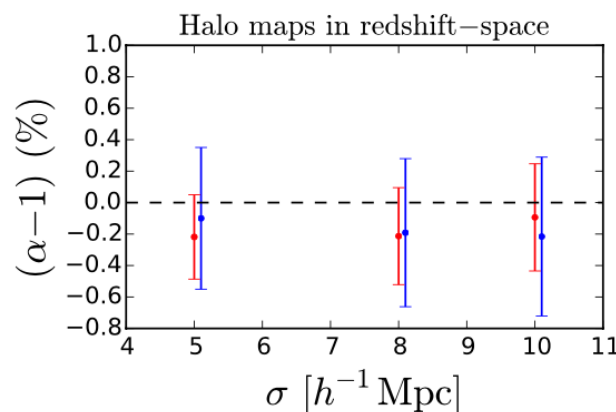
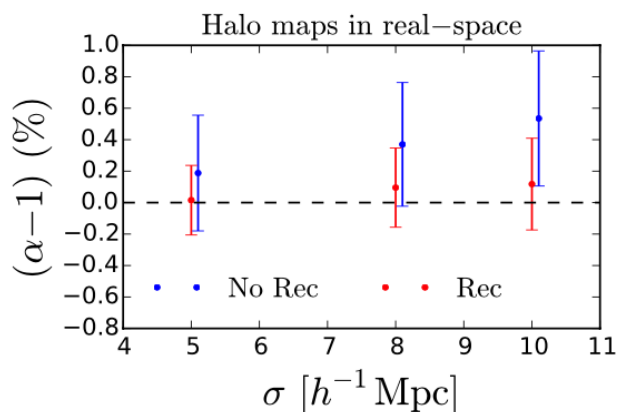
Reconstruction of the BAO peak with pixels

Matteo Viel

- Errors on α decreases by 40% after reconstruction, and this depends on the angular resolution

$$\alpha_{\parallel} \equiv \frac{H_f r_{d,f}}{H r_d} \quad \text{and} \quad \alpha_{\perp} \equiv \frac{D_A r_{d,f}}{D_A r_d}$$

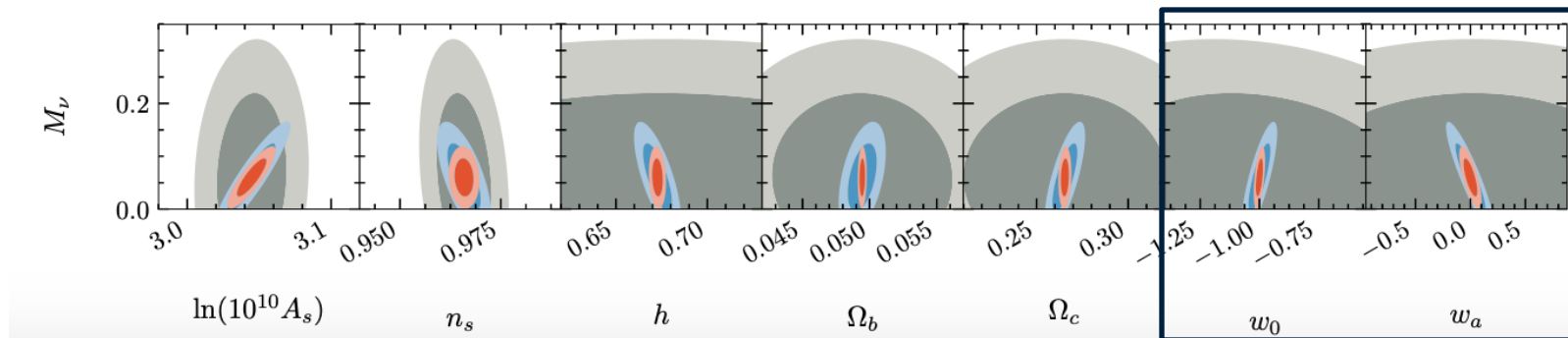
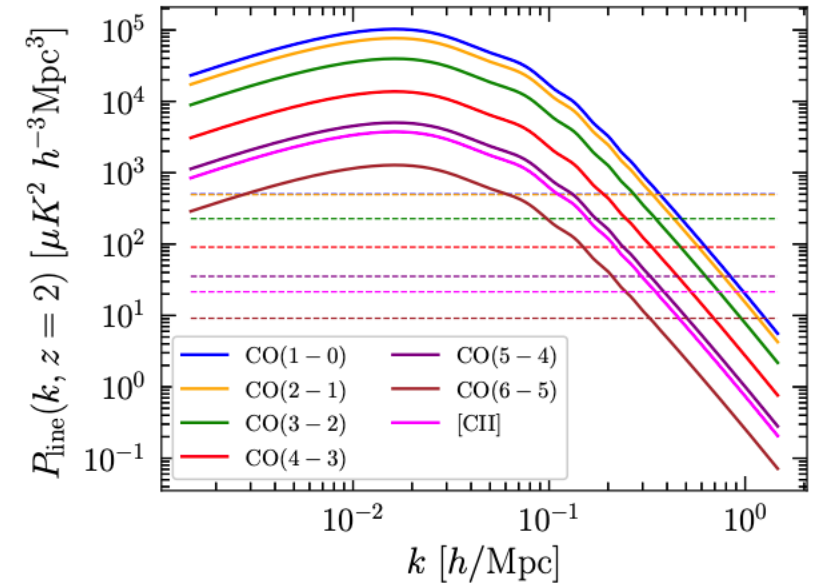
$$\alpha = \alpha_{\parallel}^{1/3} \alpha_{\perp}^{2/3}$$



Post-reionization cosmology: neutrino masses from CO

Matteo Viel

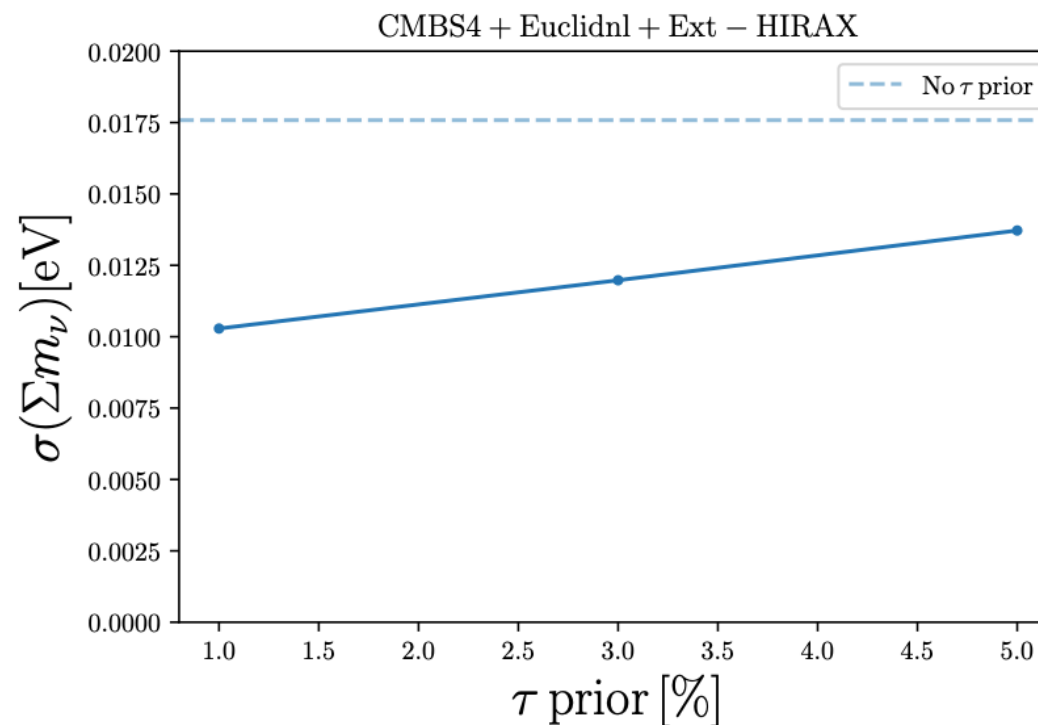
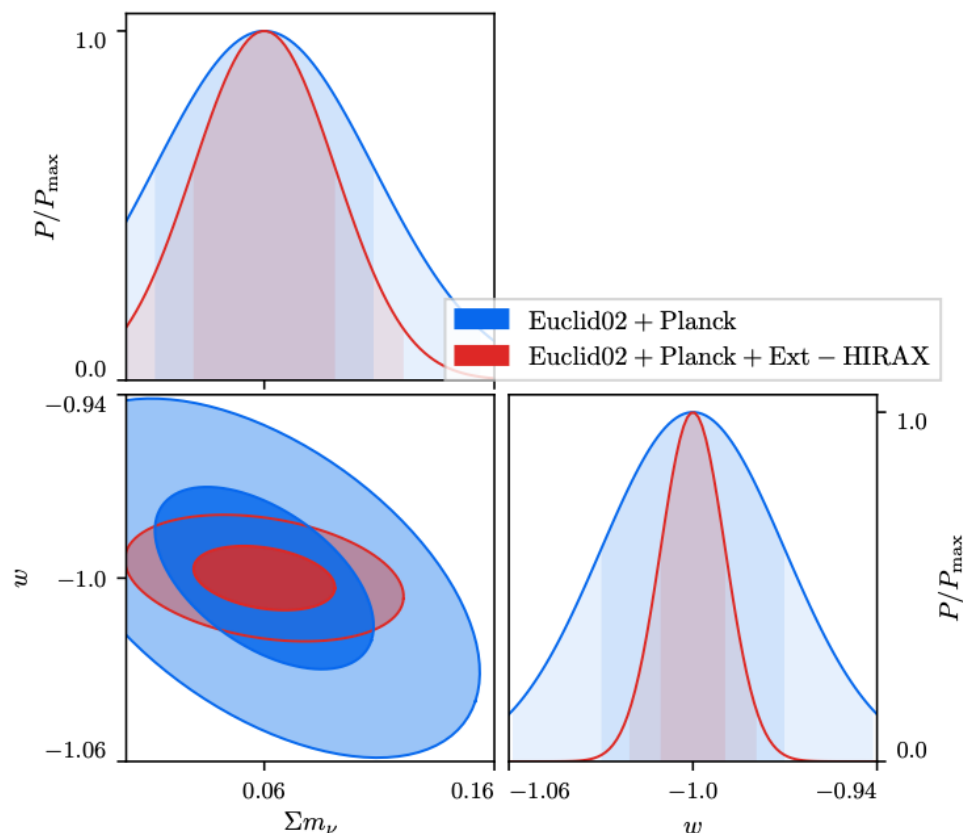
- Realistic Fisher-matrix based forecasts for CO and [CII] in a wide redshift range $z=[0,12]$
- Crucial different degeneracies pattern for LIM w.r.t. CMB data
- Especially true in the extended M_ν – CPL model
- Very promising: 40% of the sky, with 10^8 spectrometer hours and no removal of interlopers could provide $\sigma(N_{\text{eff}})\sim 0.023$ and $\sigma(M_\nu)\sim 13$ meV



Post-reionization cosmology: neutrino masses

Matteo Viel

Obuljen, Castorina, Villaescusa-Navarro, MV 2018

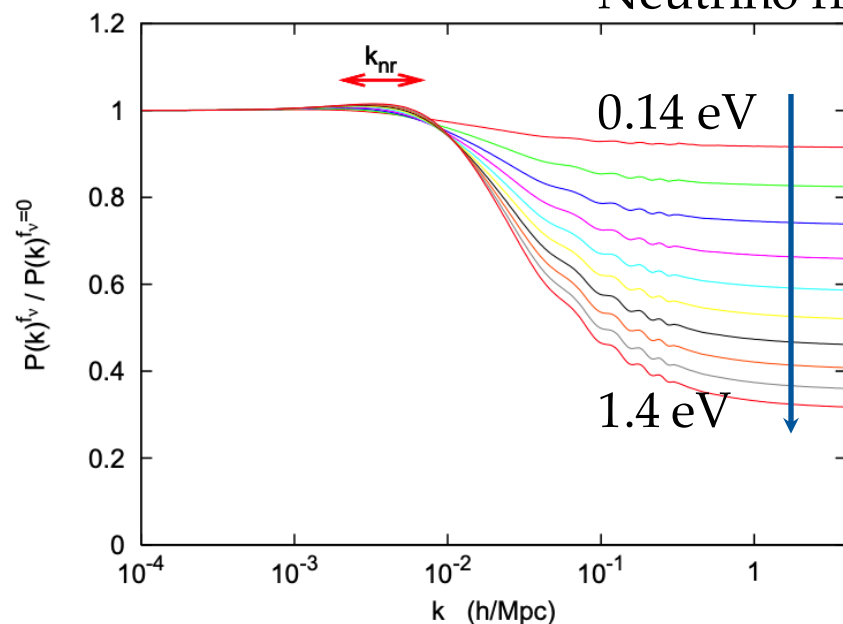


Prior on the CMB optical depth somehow fixes large scale amplitude inferred from the CMB... and helps measuring neutrino free streaming

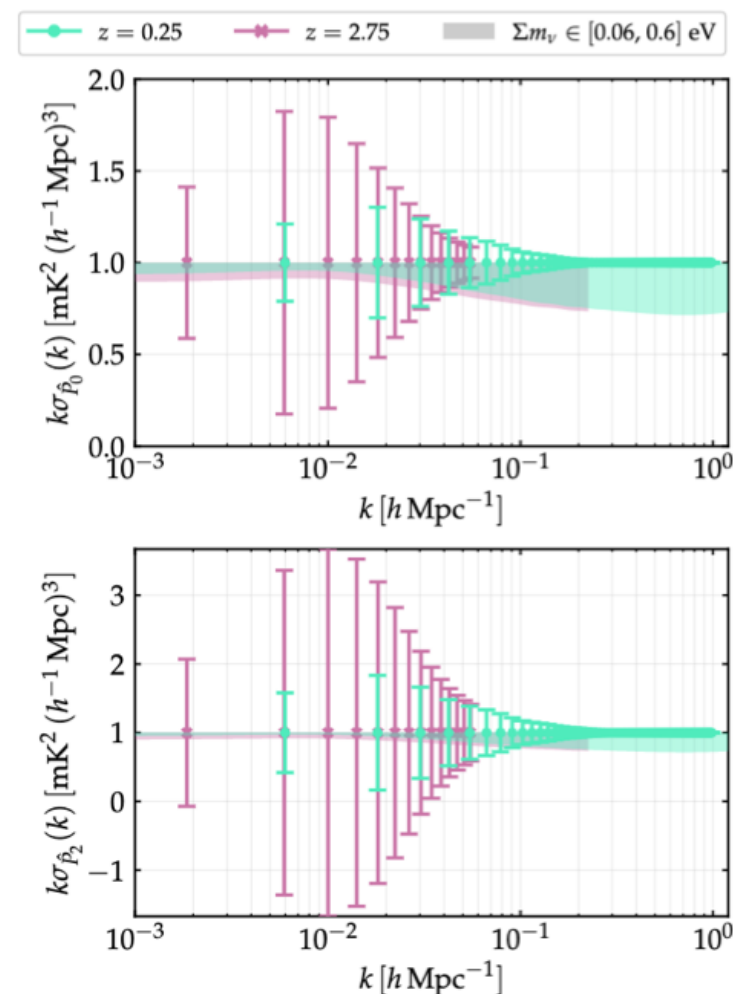
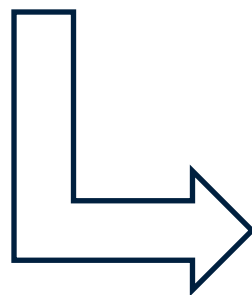
Post-reionization cosmology: neutrino masses - II

Matteo Viel

Neutrino free streaming



21cm power
SKA mid
Monopole and quadrupole

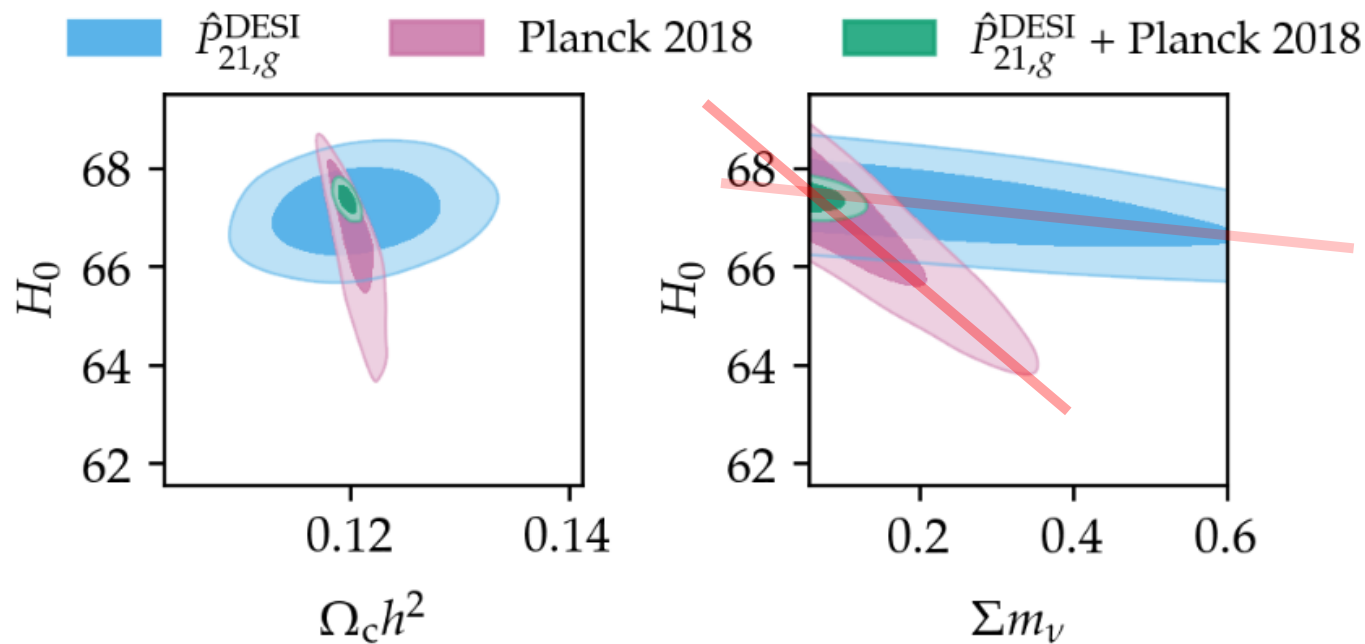


Post-reionization cosmology: neutrino masses - III

Matteo Viel

$$P_{21,g}(z, k, \mu) = \bar{T}_b(z) \left(b_{\text{HI}}(z) + f_{\text{CDM+b}}(k, z) \mu^2 \right) \left(b_g(z) + f_{\text{CDM+b}}(k, z) \mu^2 \right) P_{\text{CDM+b}}(z, k, \mu),$$

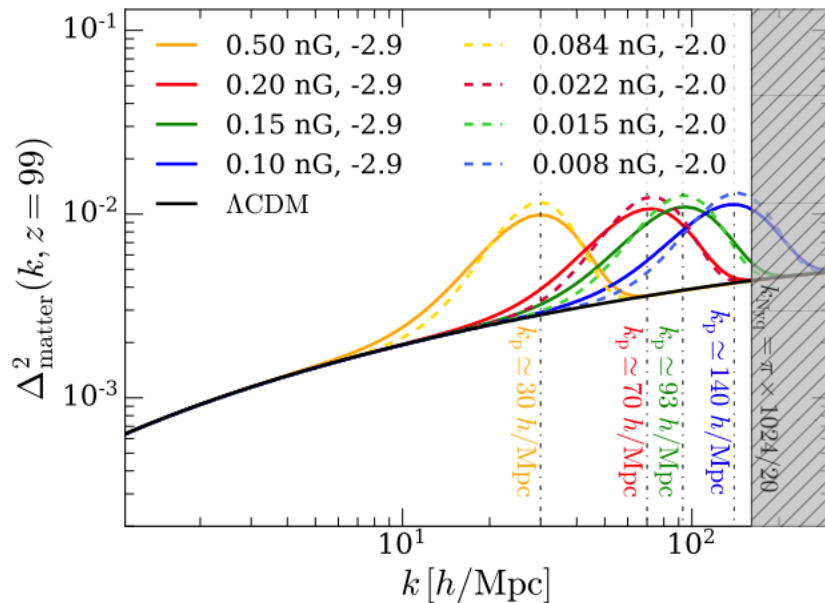
Note the different degeneracies



- Cross-correlation data alone cannot constrain neutrino masses.
- When combined with CMB data, gives constraints competitive to the ones obtained with auto-power.

Likelihoods	$\Sigma m_\nu^{\text{fid}} = 0.06 \text{ eV}$	$\Sigma m_\nu^{\text{fid}} = 0.1 \text{ eV}$
$\hat{P}_0 + \hat{P}_2$	< 0.287	< 0.317
+ nuisances	< 0.425	< 0.452
Planck 2018		
+ $\hat{P}_0 + \hat{P}_2$	< 0.105	0.098 ± 0.022
+ nuisances	< 0.126	< 0.151
Planck 2018		
+ $\hat{P}_{21,g}^{\text{DESI}}$	< 0.116	$0.099^{+0.020}_{-0.033}$
+ nuisances	< 0.155	< 0.177
Planck 2018		
+ $\hat{P}_{21,g}^{\text{Euclid}}$	< 0.117	$0.100^{+0.021}_{-0.032}$
+ nuisances	< 0.156	< 0.180

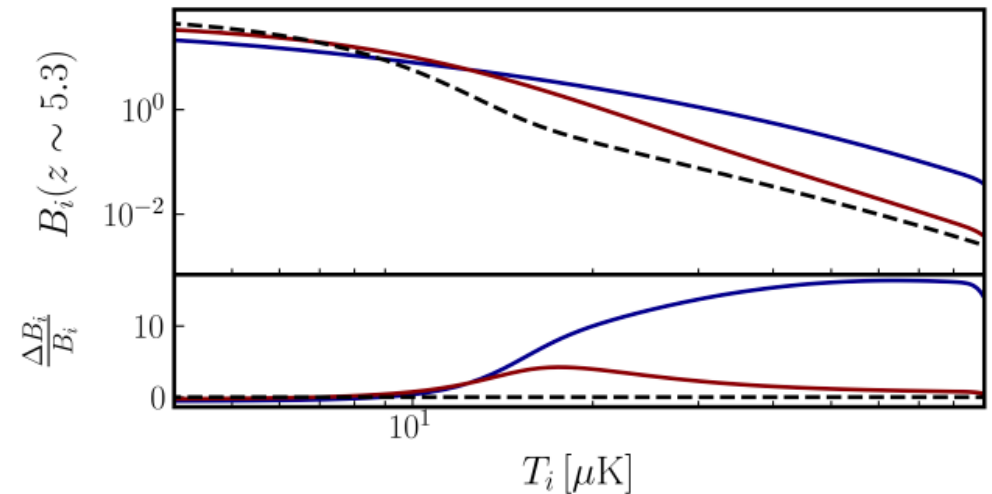
Increase of power in total matter power Spectrum due to Lorentz force affecting Baryons' clustering



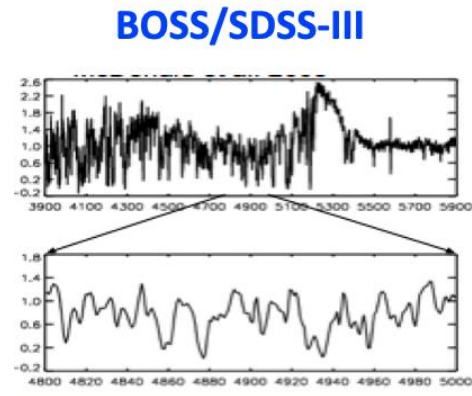
From the forest: $B \sim 0.2$ nG “hint”

Impact on VID
COMAP EoR survey + other instruments at $z=2-3$

$$B_i = N_{\text{vox}} \int_{T_i}^{T_i + \Delta T_i} \mathcal{P}_{\text{tot}}(T) dT,$$



From CO: IM $B \sim 0.006$ nG can be probed

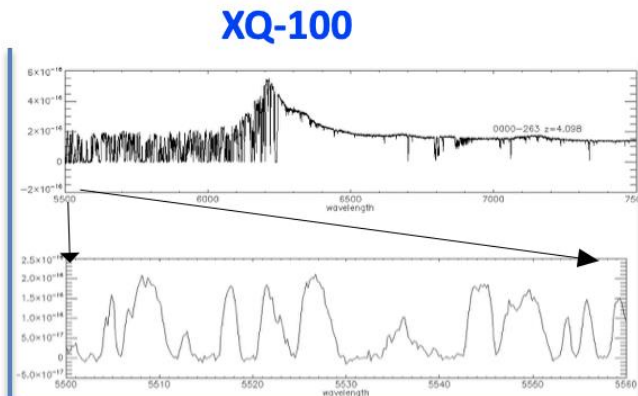


Low resolution BOSS
and SDSS-III spectra
S/N~2-3 - 160,000
spectra

Used to detect BAOs at
 $z=2.3$ and correlations
in the transverse
direction

Used to place
stringent constraints
on neutrino masses
 <0.12 eV

*Busca+13, Slosar+14, Font-Ribera+14
Palanque-Delabrouille+15
Seljak+06, Baur+16, Yèche+17 etc.*

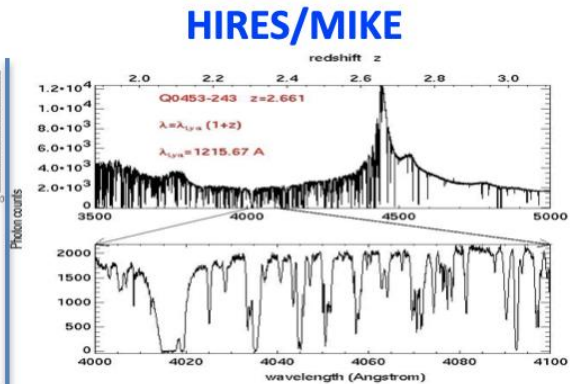


Medium resolution X-
Shooter VLT spectra
S/N ~ 30

100 spectra at $z>3.5$

Used to place
stringent constraints
on Warm Dark Matter in
combination with high
res. spectra

*Irsic, MV+ 17a,17b
Lopez+16, Irsic+16*



High resolution VLT
or Keck spectra S/N
~100 - ~hundreds of
spectra

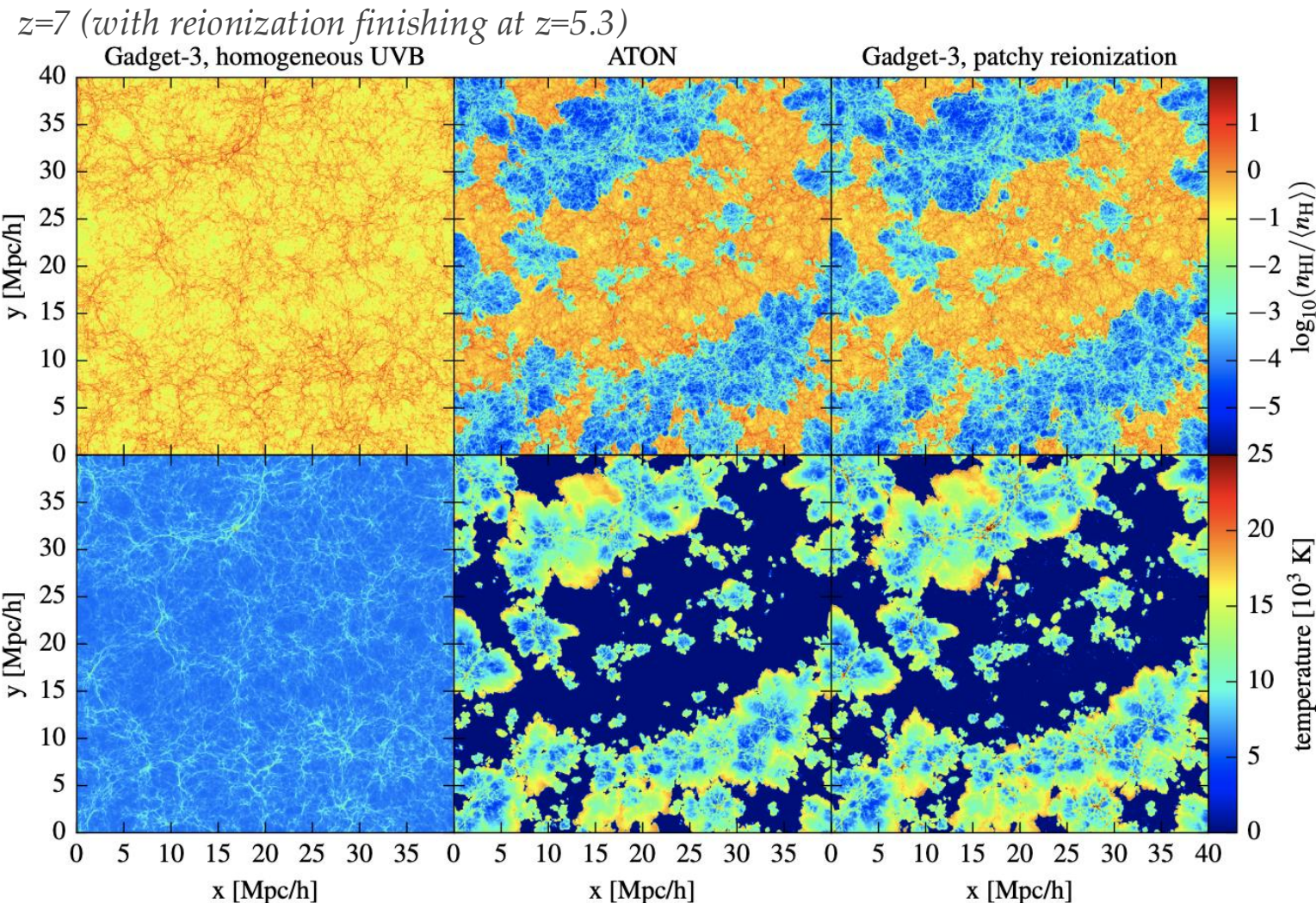
Used for WDM,
astrophysics of the
IGM and galaxy
formation, variation
of fundamental
constants

*MV+05,08,13, Becker+11
Yèche+17, Garzilli+18,
Bosman+18*

The simulations - I

Matteo Viel

<https://www.nottingham.ac.uk/astronomy/sherwood/>



Bolton+17

Puchwein, Bolton+



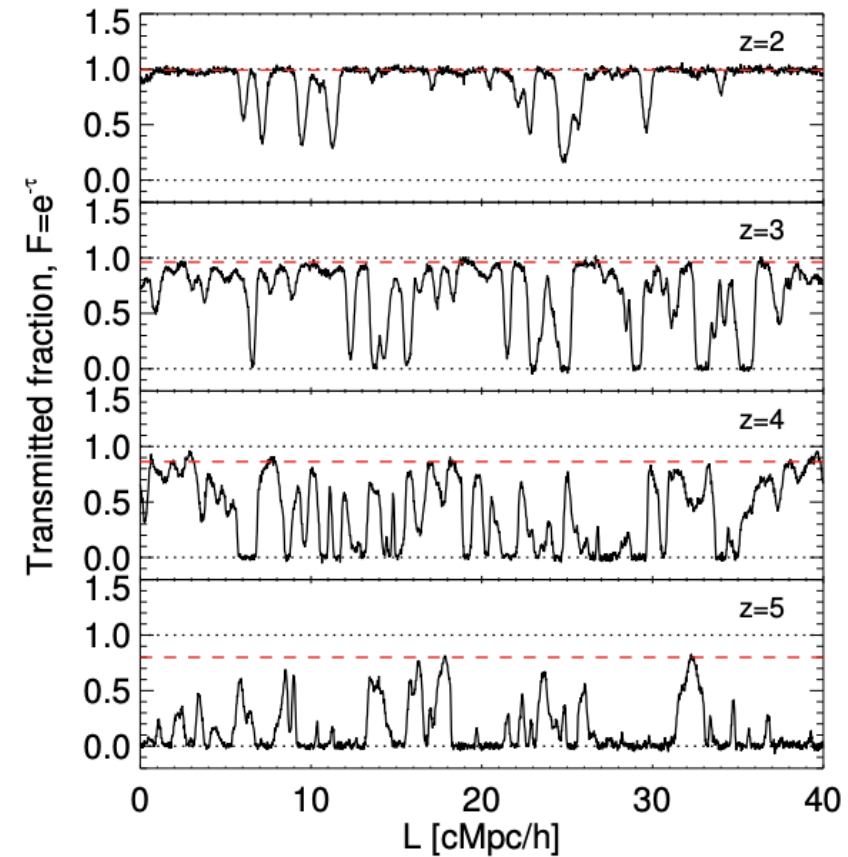
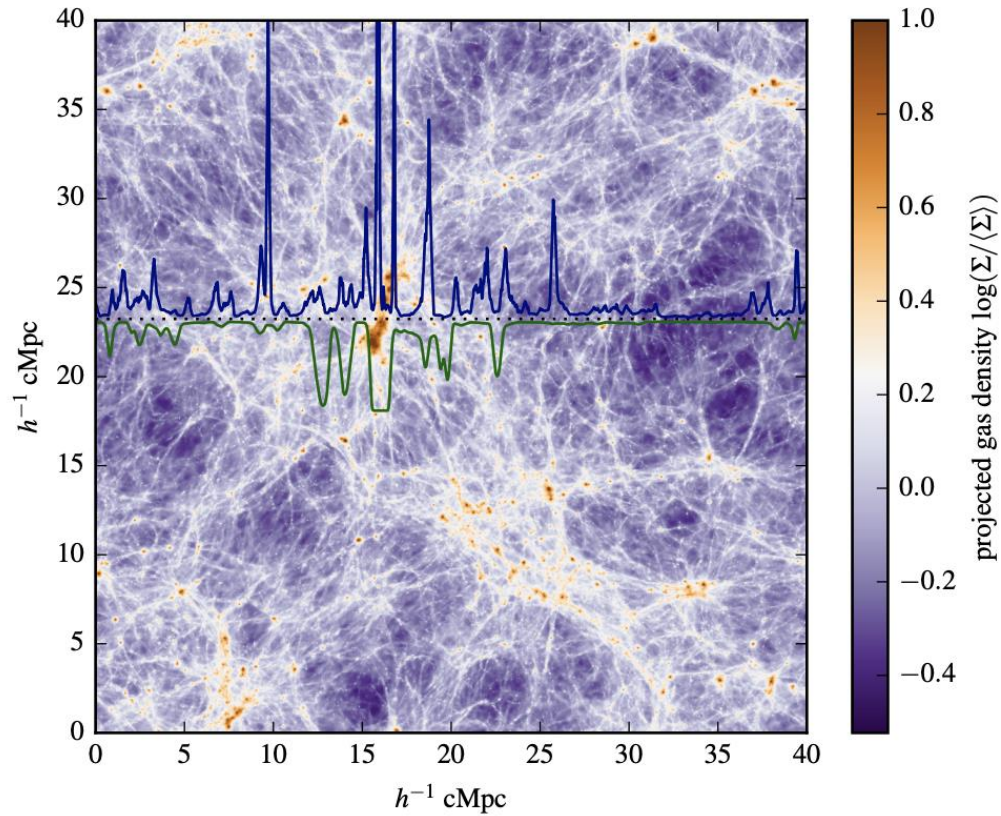
J. Bolton

E. Puchwein

- **Sherwood-Relics suite**
(>200 simulations: boxes 5-160 cMpc/h;
 $M_{\text{gas}}=3.7\text{e}3\text{-}6.4\text{e}6 M_{\odot}$) –
about 75 Million CPU hrs
(2017-now)
- G3 code + ATON to
perform radiative
transfer for patchy
reionization

The simulations - III

Matteo Viel



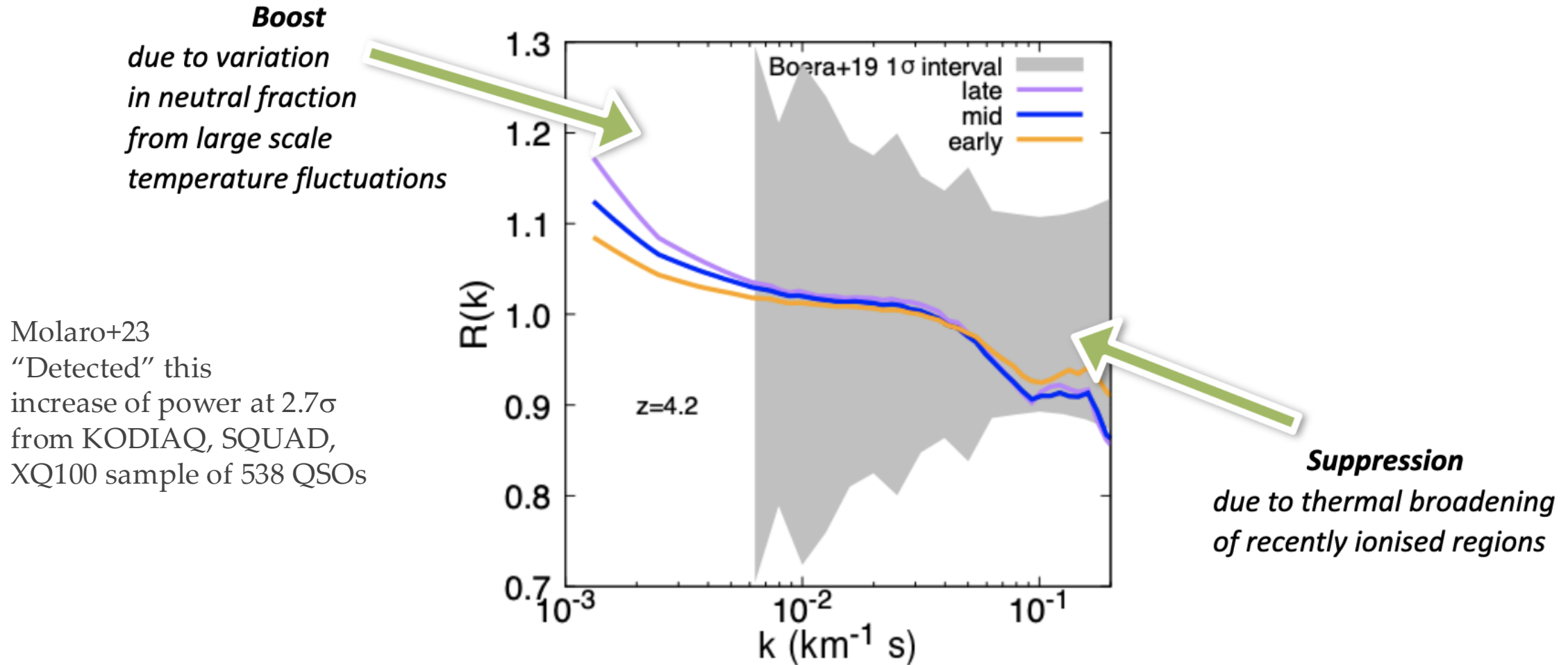
- Most of the flux statistics are in agreement with Λ CDM – 216,000 flux models fed into MCMC analysis

Increasing $z \rightarrow$ increasing HI \rightarrow more absorption

eld

Patchy Reionization - II

Matteo Viel

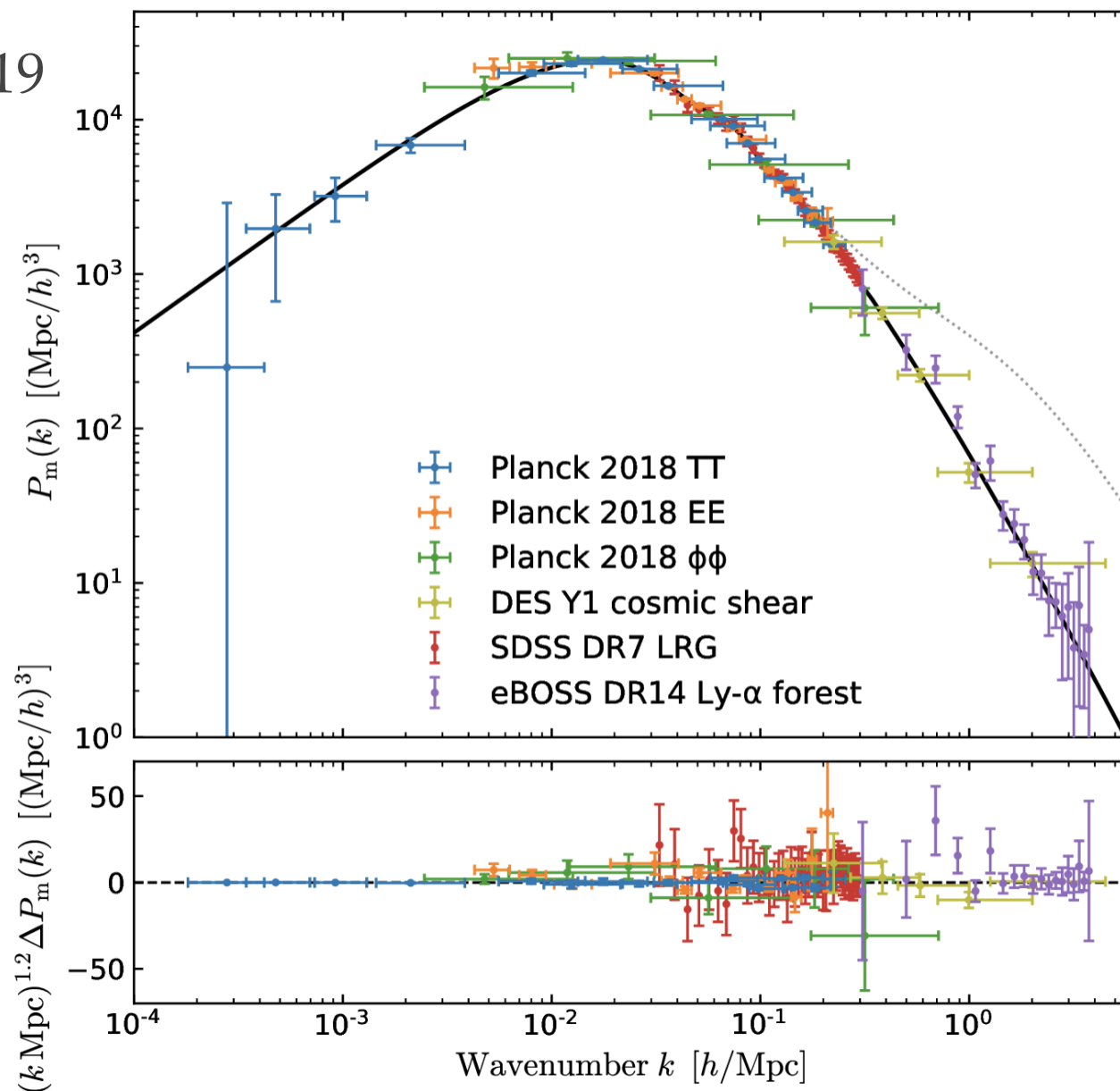


Molaro, Bolton, Irsic,... MV 2021&

Long lever arm of the linear power spectrum

Matteo Viel

Chabanier+19



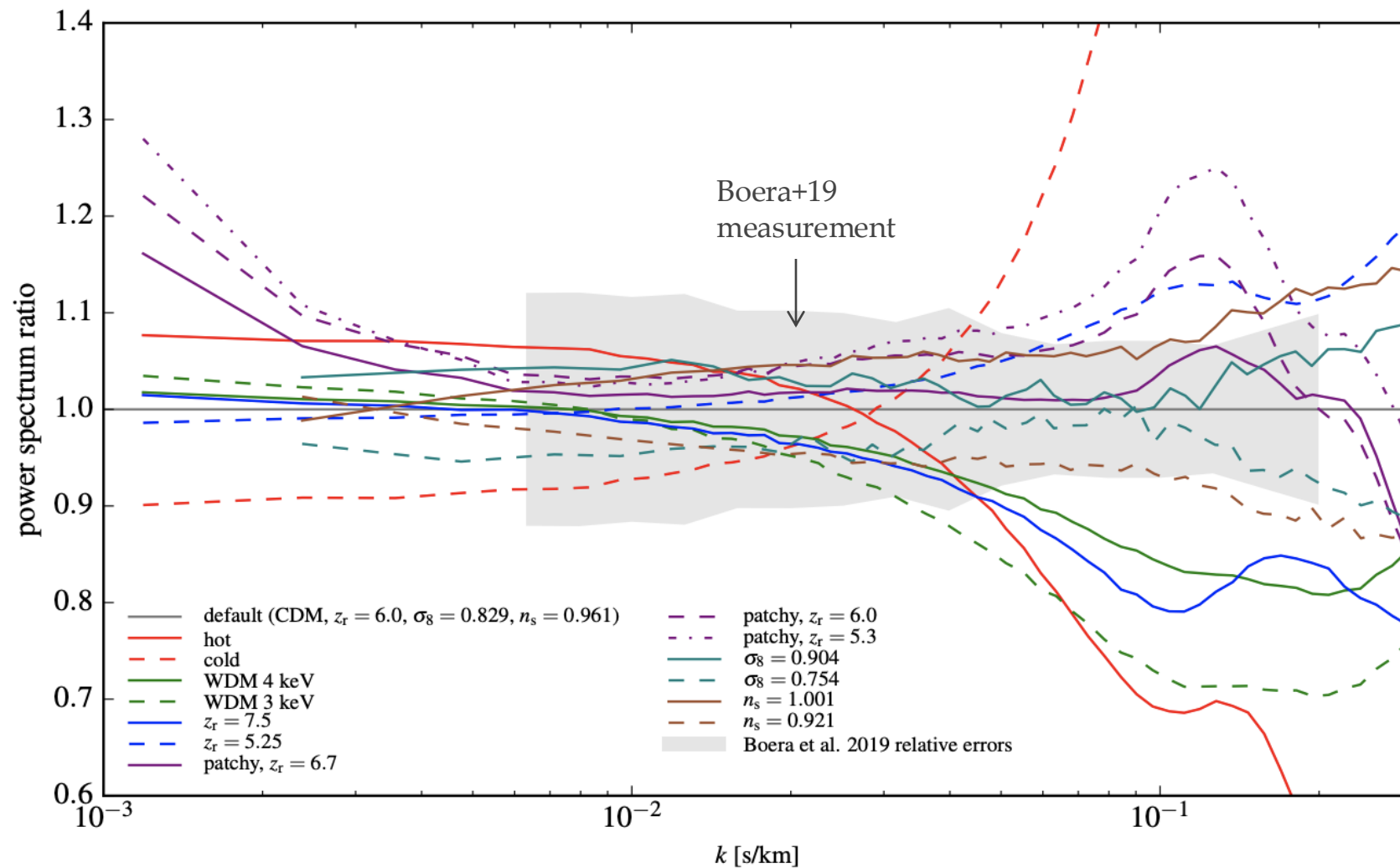
Two reasons for why Ly α is so constraining:

- 1) 1D is projected power
- 2) We are at high- z possibly closer to linear regime.

Impact on 1D flux power

Matteo Viel

Simulated 1D flux power @ $z=4.6$



More power

Less power

Large scales

Small scales

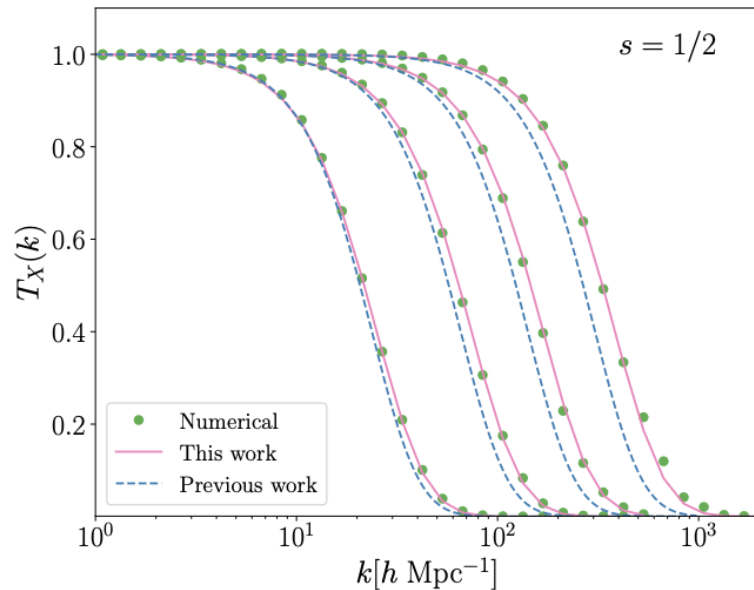
Matter power spectrum and WDM

Matteo Viel

$$T(k) \equiv [1 + (k/k_{break})^p]^{-10/p} \quad \text{with } p = 2.24$$

$$k_{break} = \frac{1}{0.24} X^{0.83} \left(\frac{\omega_X}{0.25 \times 0.7^2} \right)^{0.16} \text{Mpc}^{-1} \quad \text{with } X \equiv \frac{m_X/T_X}{1 \text{ keV}} T_\nu^a$$

Important: unlike active neutrinos this depends on both DM density and X
Because free streaming horizon depends on those

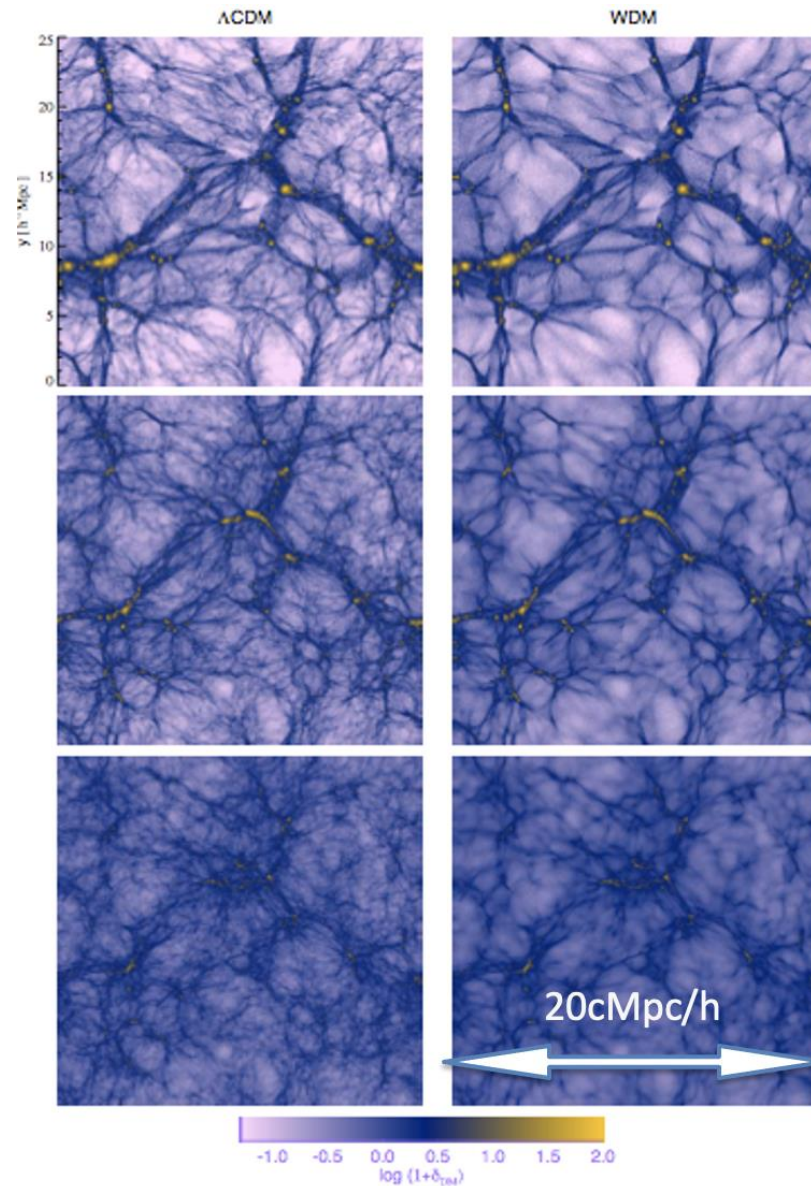


Viel+05;

Vogel&Abazajian <https://arxiv.org/abs/2210.10753>

A warm cosmic web?

Matteo Viel



$z=0$

$$k_{\text{FS}} \sim 15.6 \frac{h}{\text{Mpc}} \left(\frac{m_{\text{WDM}}}{1 \text{ keV}} \right)^{4/3} \left(\frac{0.12}{\Omega_{\text{DM}} h^2} \right)^{1/3}$$

$z=2$

Free streaming scale
of thermal warm dark
matter

$z=5$

Viel et al 2005

The smoothing scales

Matteo Viel

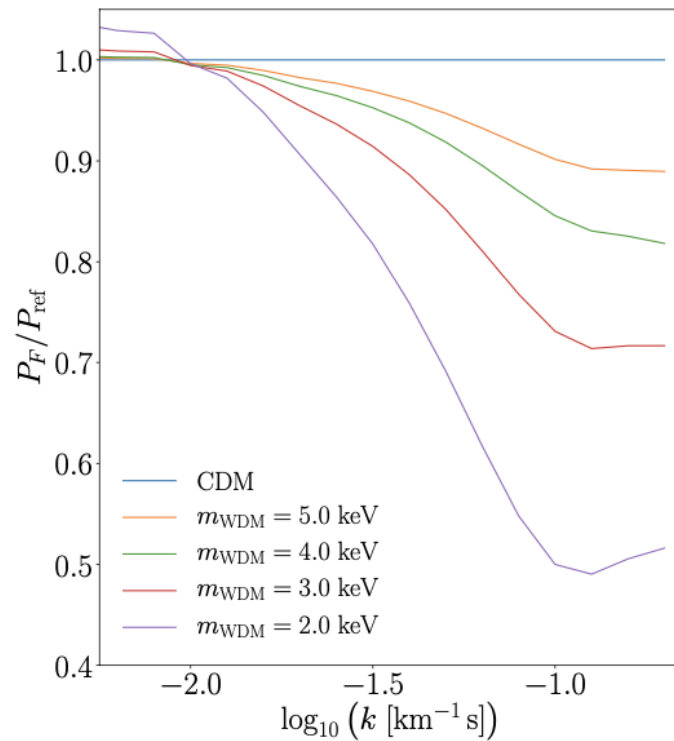
Vid Irsic



Unveiling Dark Matter free-streaming at the smallest scales with high redshift Lyman-alpha forest

Vid Irsi^{1,2}, Matteo Viel^{3,4,5,6,7}, Martin G. Haehnelt^{1,8}, James S. Bolton⁹, Margherita Molaro⁹, Ewald Puchwein¹⁰, Elisa Boera^{5,6}, George D. Becker¹¹, Prakash Gaikwad¹², Laura C. Keating¹³, Girish Kulkarni¹⁴
¹Kavli Institute for Cosmology, University of Cambridge

WDM free streami



The smoothing scales

Matteo Viel

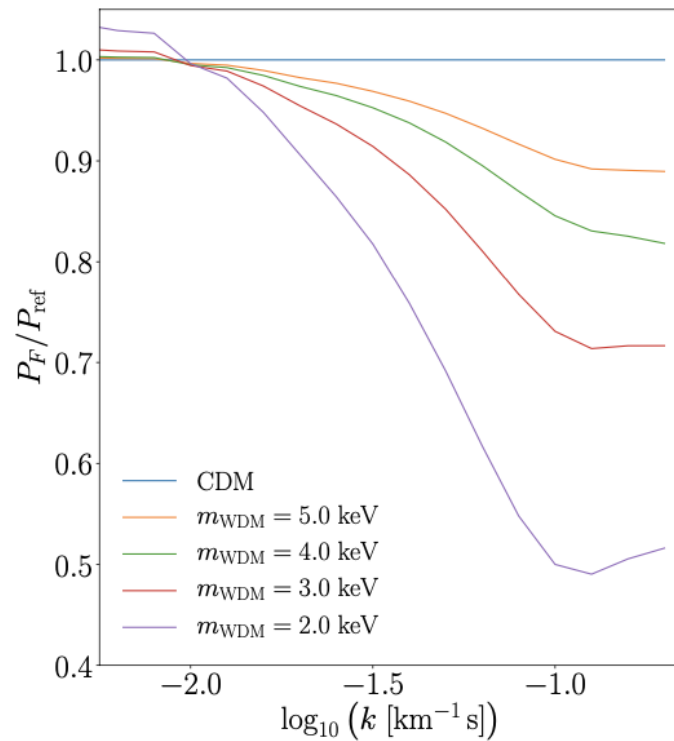
Vid Iršič



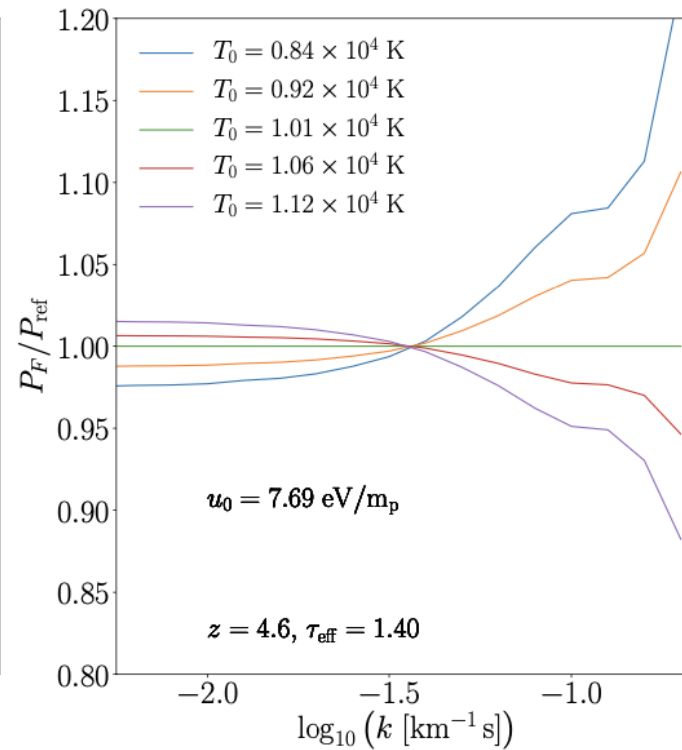
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WDM free streaming



Thermal broadening



The smoothing scales

Matteo Viel

Vid Irsic

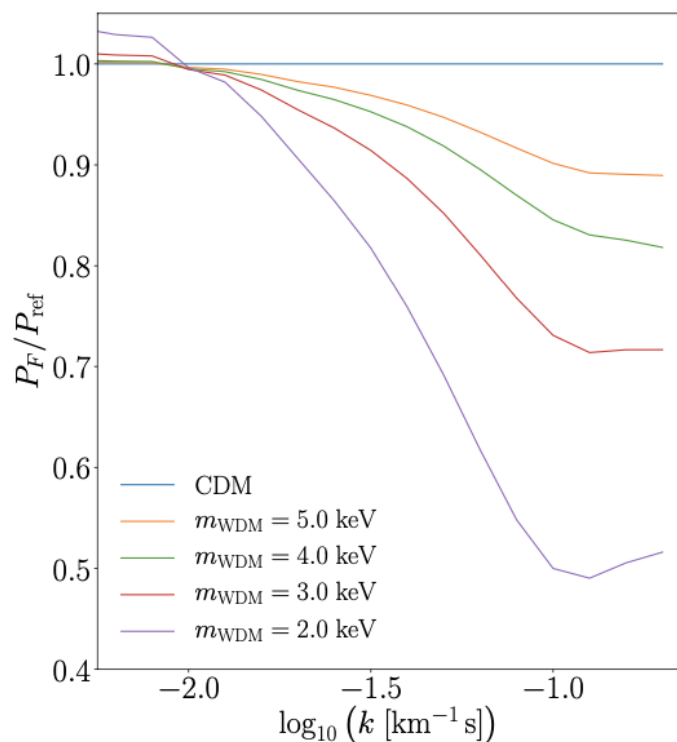


Unveiling Dark Matter free-streaming at the smallest scales with high redshift Lyman-alpha forest

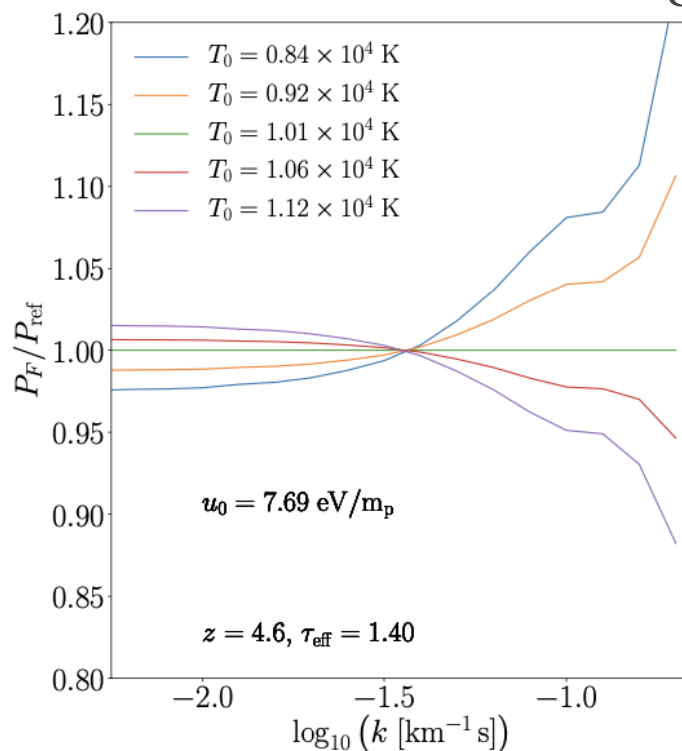
Vid Irsic^{1,2}, Matteo Viel^{3,4,5,6,7}, Martin G. Haehnelt^{1,8}, James S. Bolton⁹, Margherita Molaro⁹, Ewald Puchwein¹⁰, Elisa Boera^{5,6}, George D. Becker¹¹, Prakash Gaikwad¹², Laura C. Keating¹³, Girish Kulkarni¹⁴

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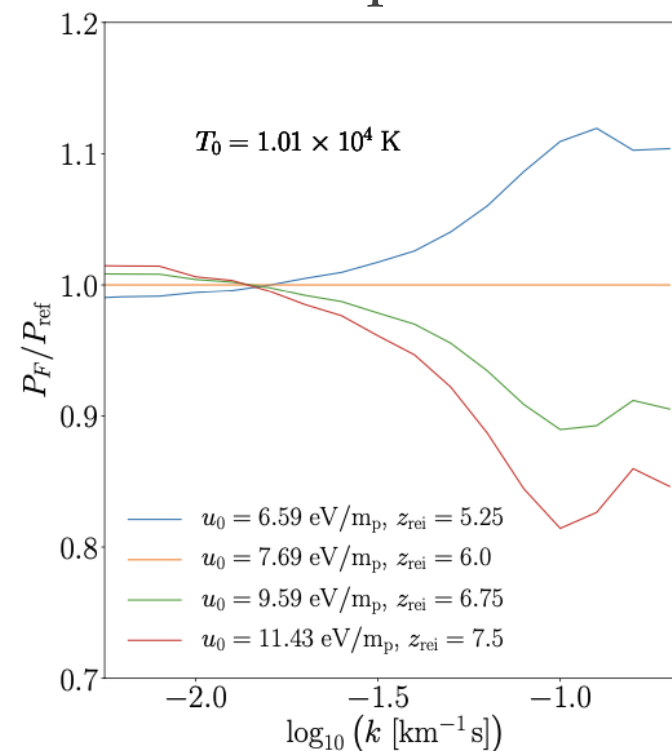
WDM free streaming



Thermal broadening

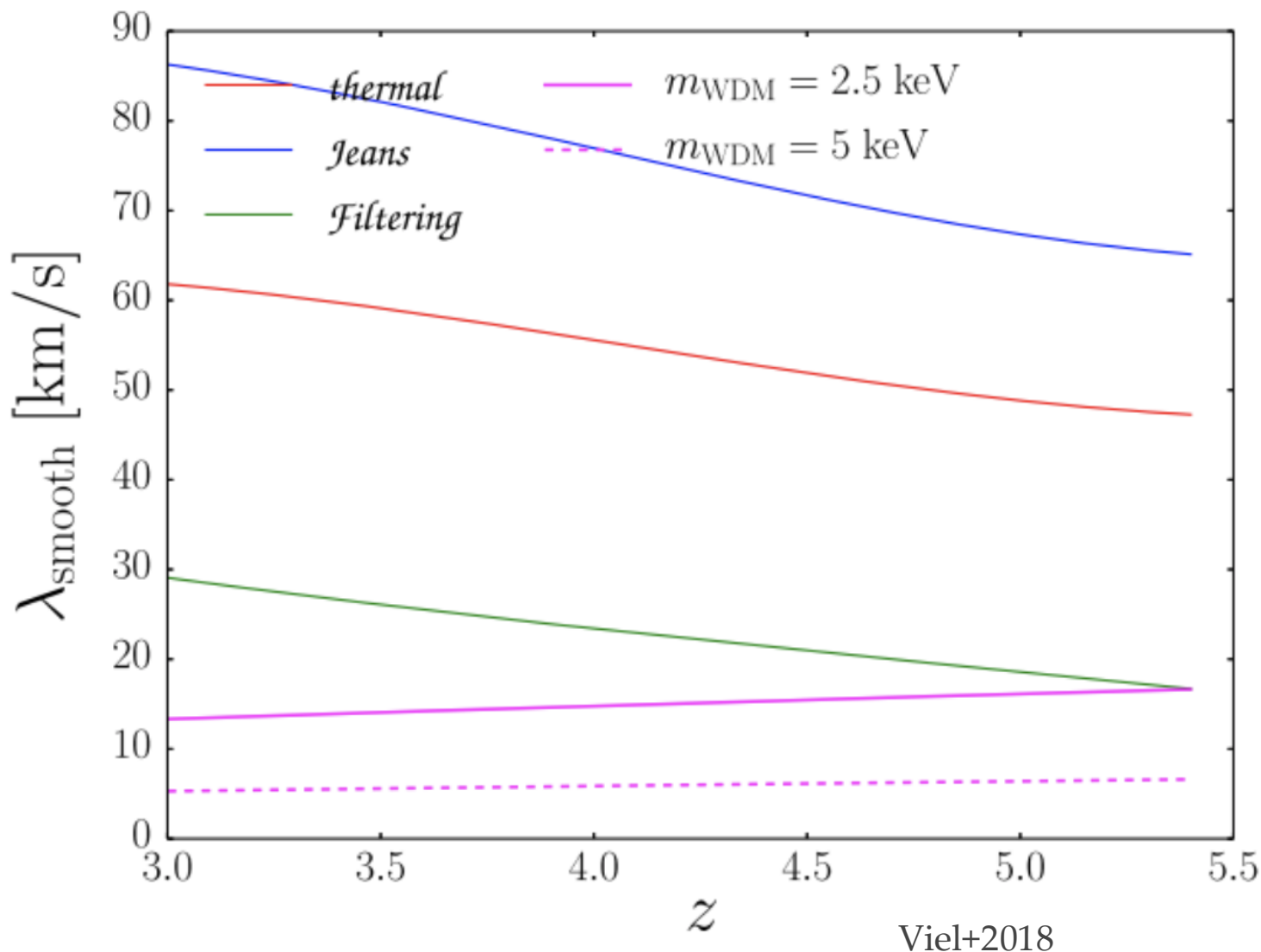


Gas pressure



$$u_0(t) = \int_0^t dt \frac{\mathcal{H}}{\bar{\rho}_m} \frac{3k_B}{2\mu} \quad H \text{ is heating rate}$$

The smoothing scales - II

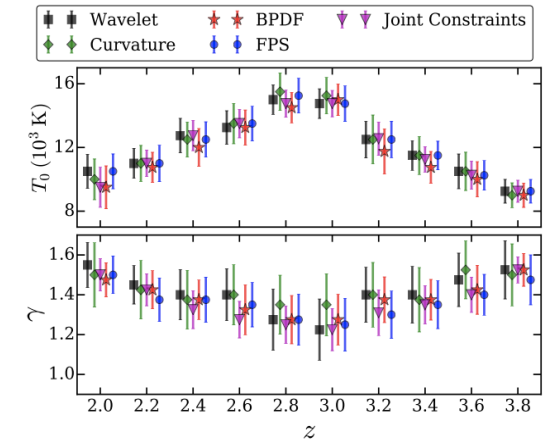
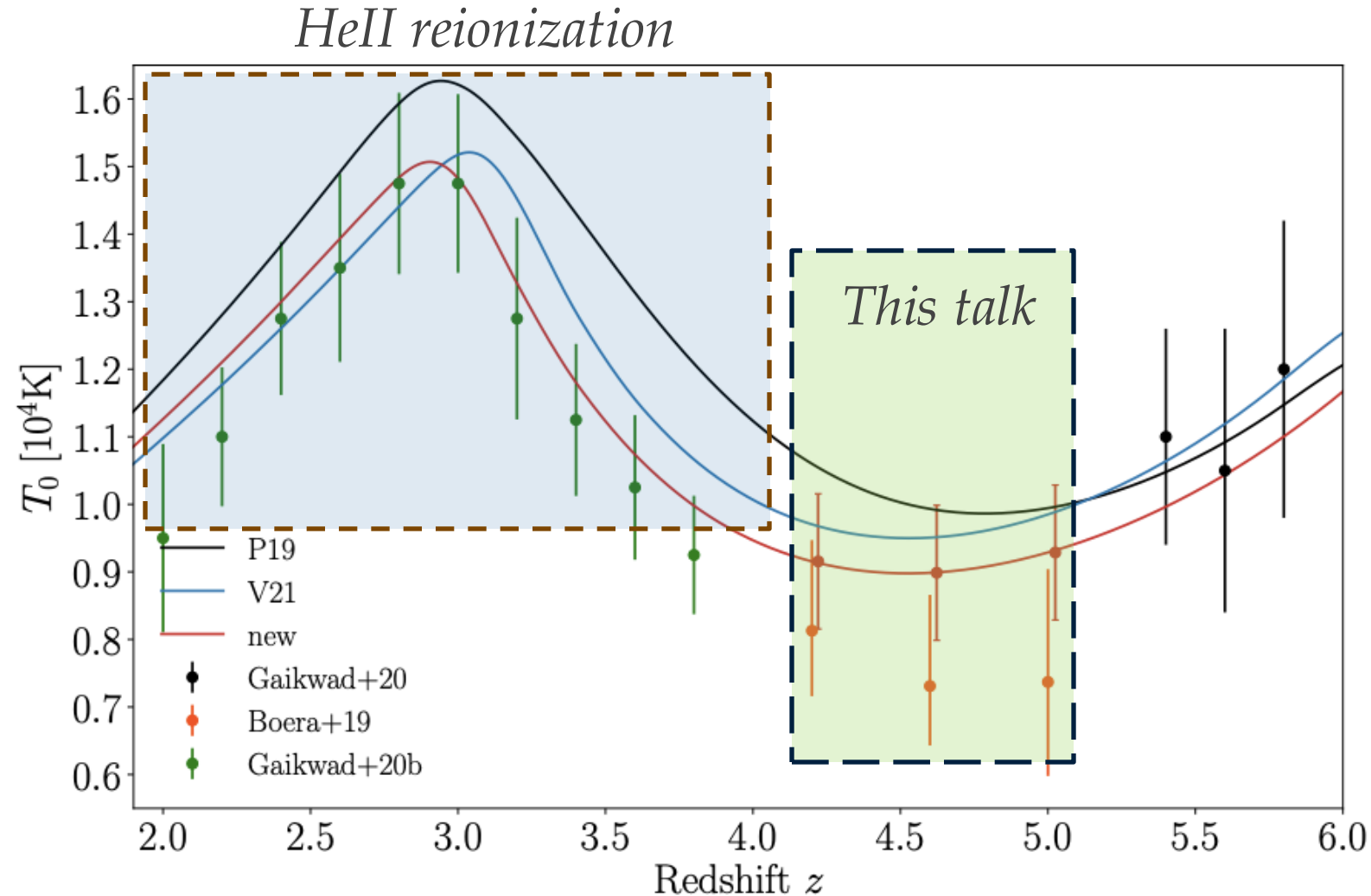


Different physical scales (on top of instrumental resolution) affect the power spectrum cutoff:

- thermal: instantaneous temperature at that redshift;
- filtering scale: depends on all the past thermal history – related to Jeans scale;
- WDM cutoffs are basically redshift independent

The IGM thermal state

Matteo Viel

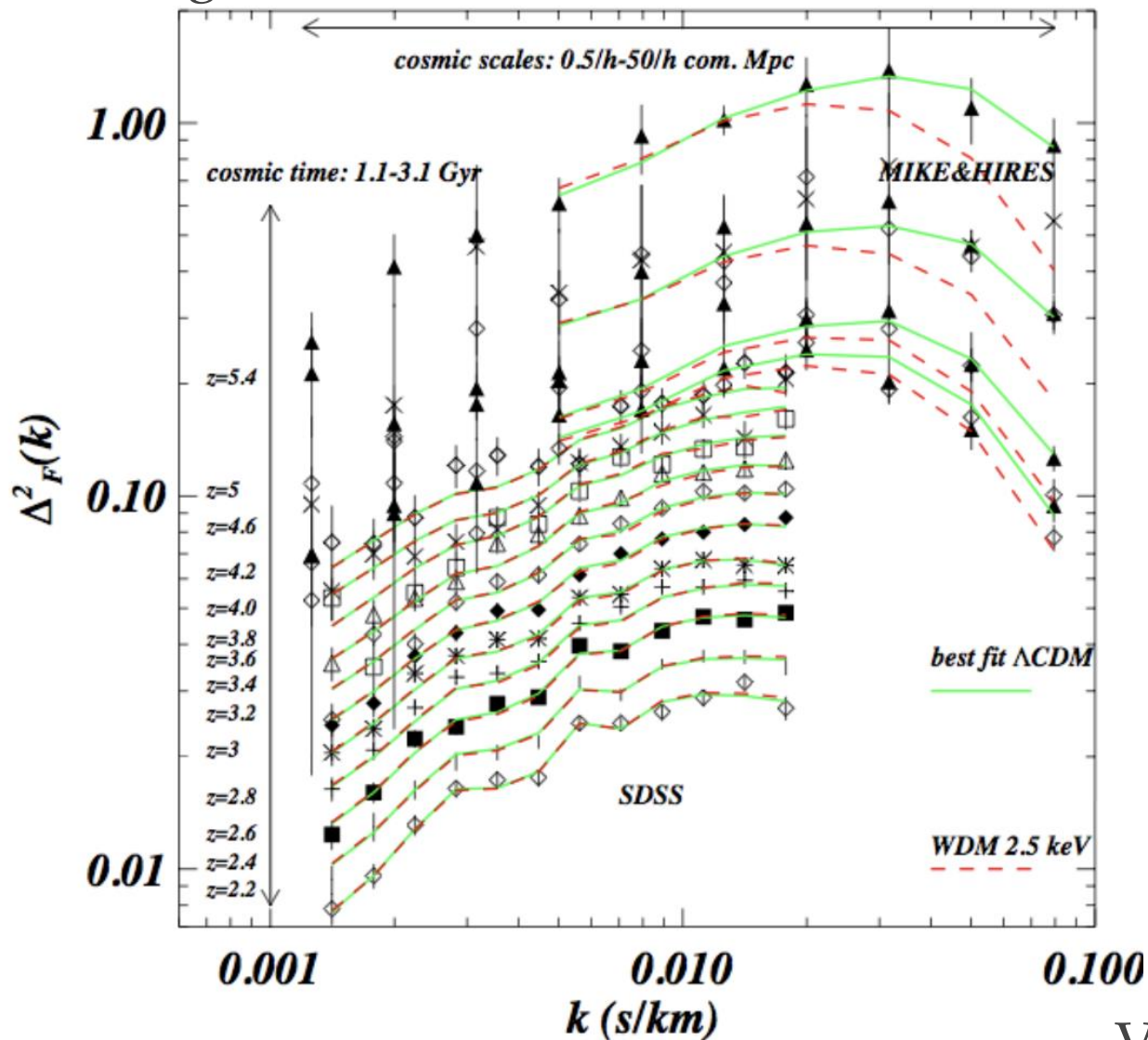


- Constraints obtained with a variety of data and methods
- Sensitive to lines rather than the lines' clustering

Large scales

Small scales

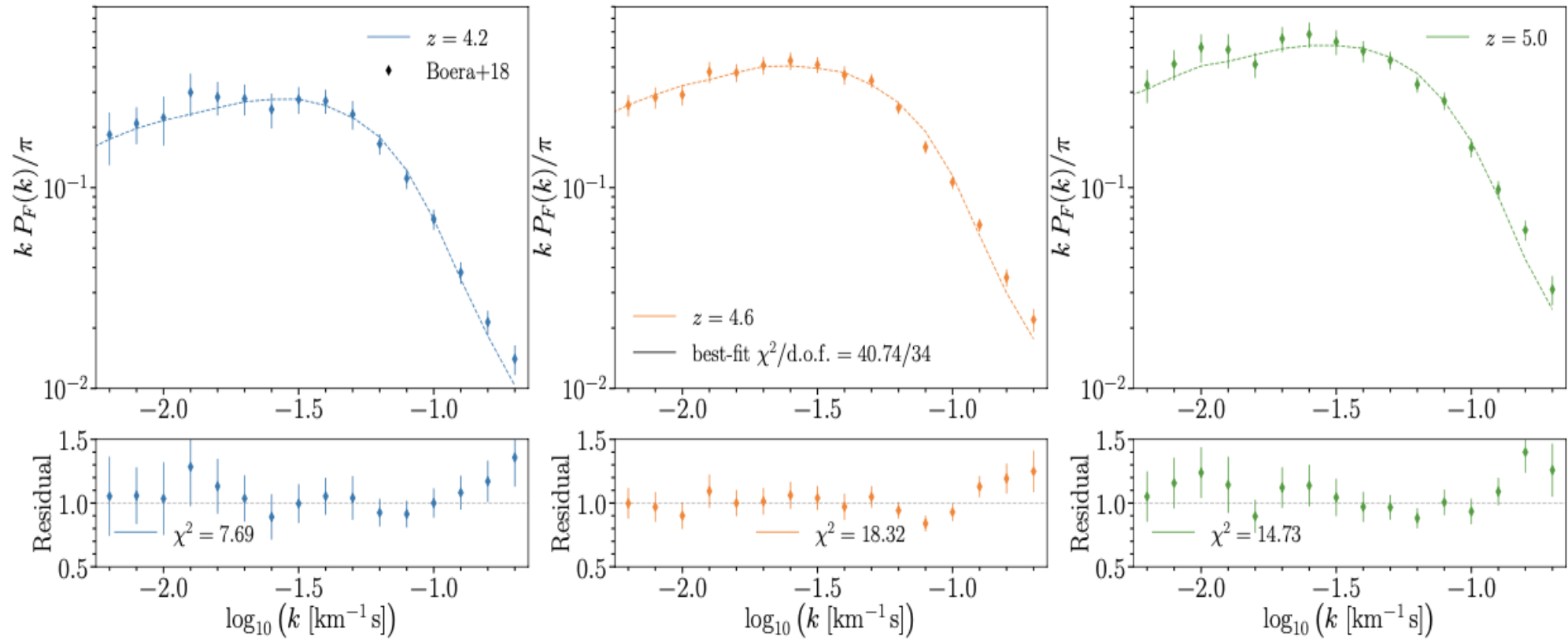
1D flux power



➤ Test of structure formation for a Λ CDM Universe in a **unique “pre-galactic” environment**

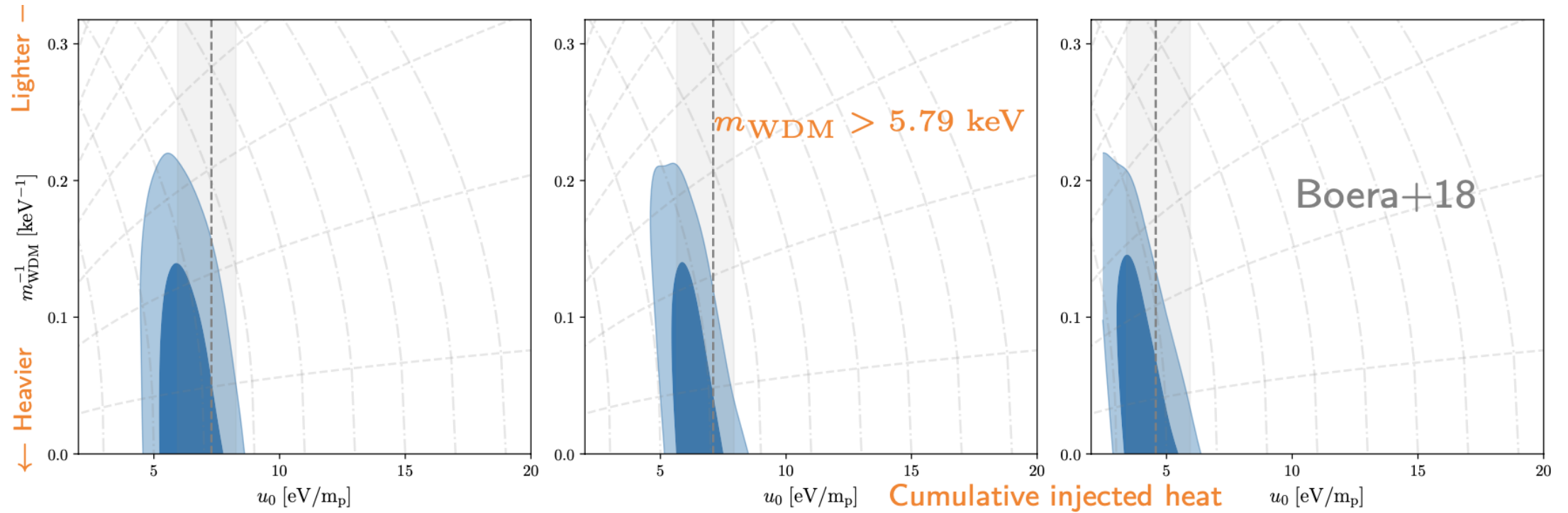
➤ $m_{\text{WDM}} > 3.3 \text{ keV}$ (2σ C.L.)

Note: 10 yrs later only a factor 2 more high z QSOs



Thermal WDM

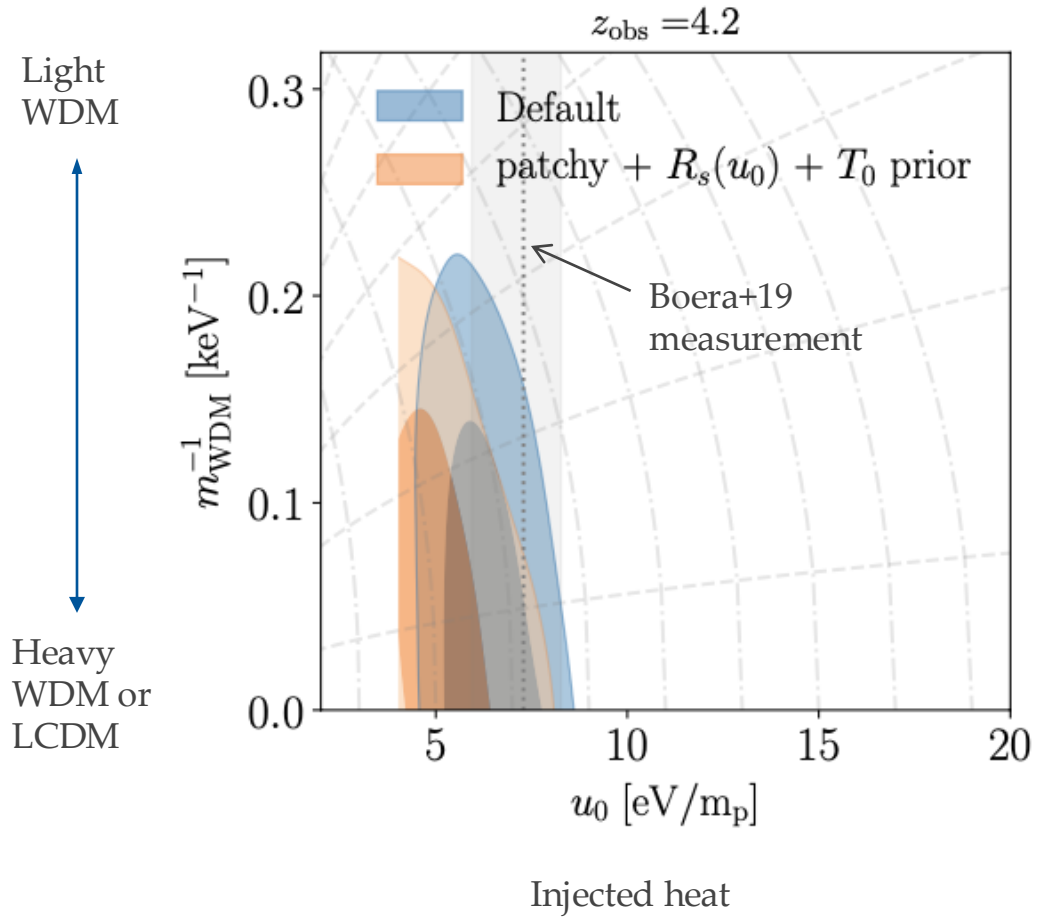
Matteo Viel



Irsic+23

Thermal WDM

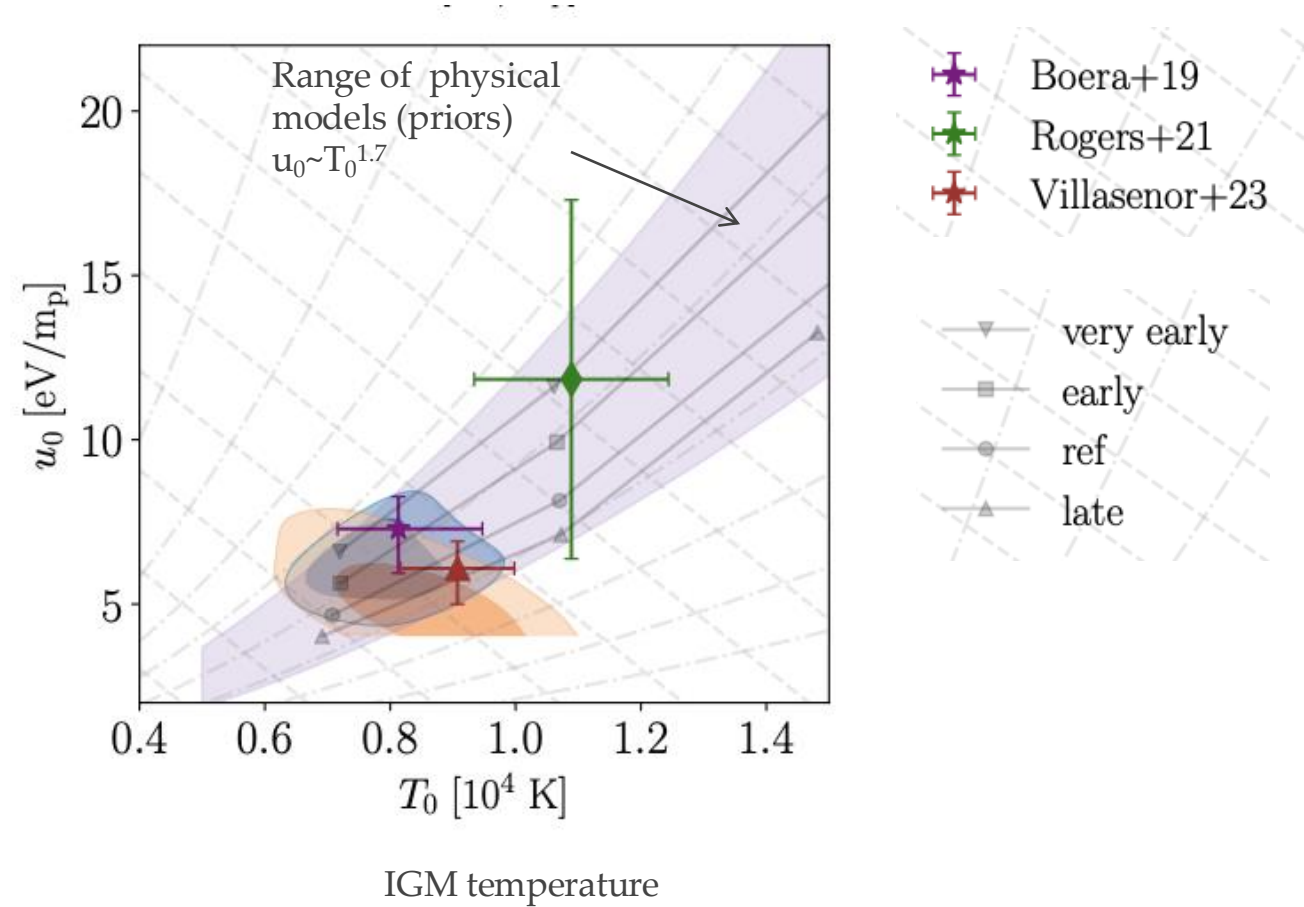
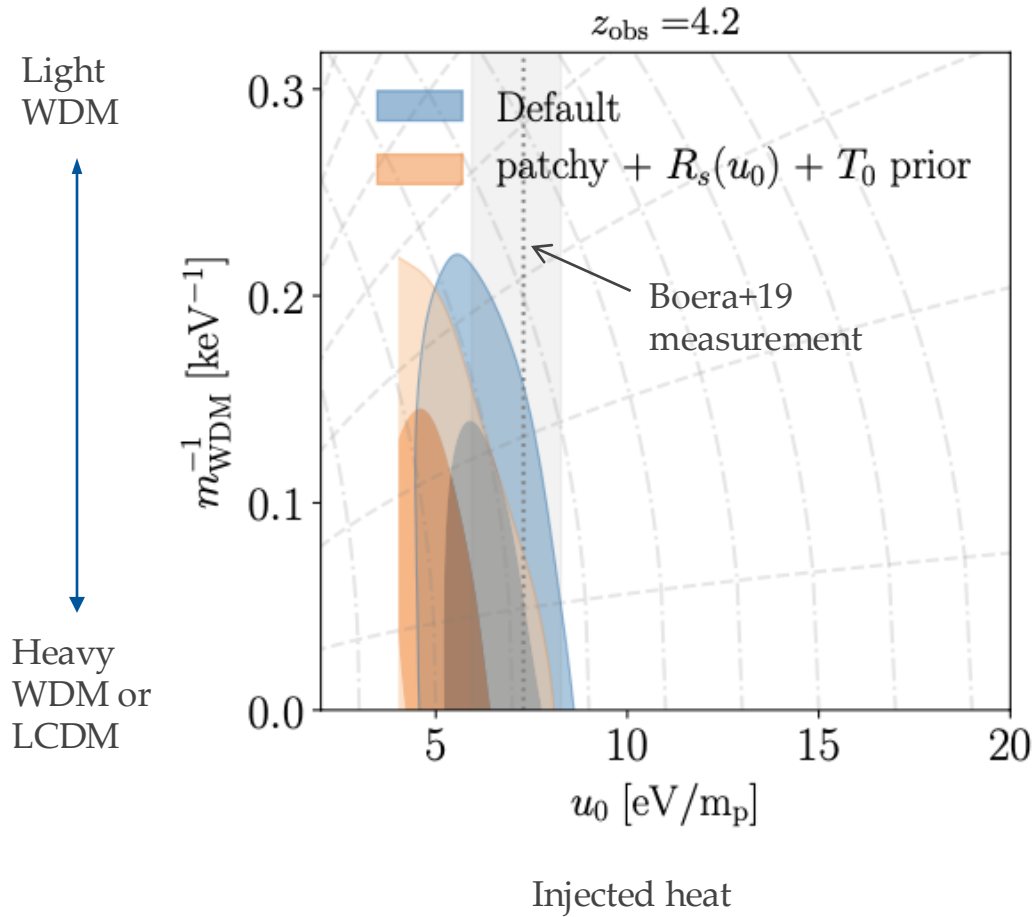
Matteo Viel



Irsic, MV +23

Thermal WDM - II

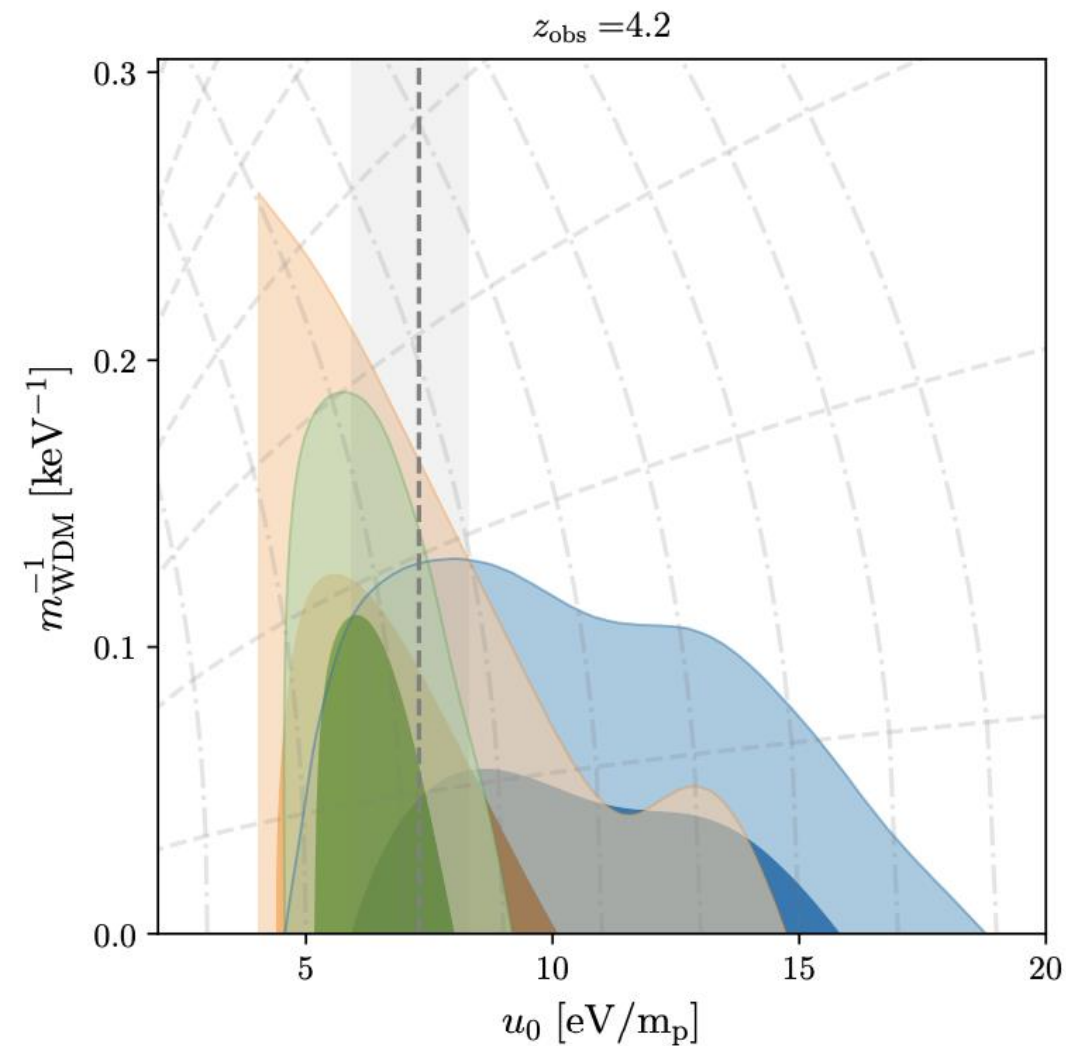
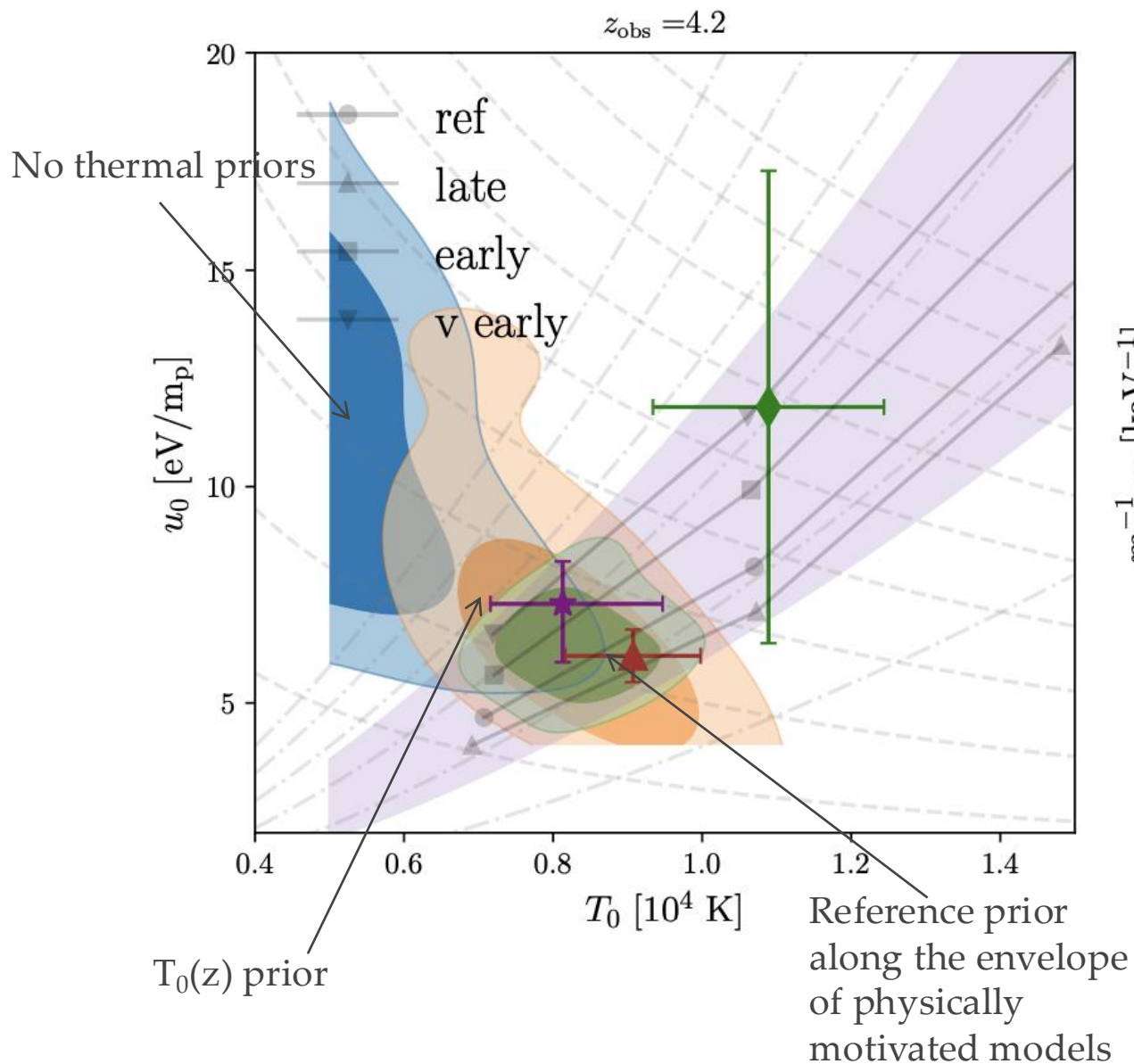
Matteo Viel



Irsic, MV +23

Thermal WDM – the effect of thermal priors

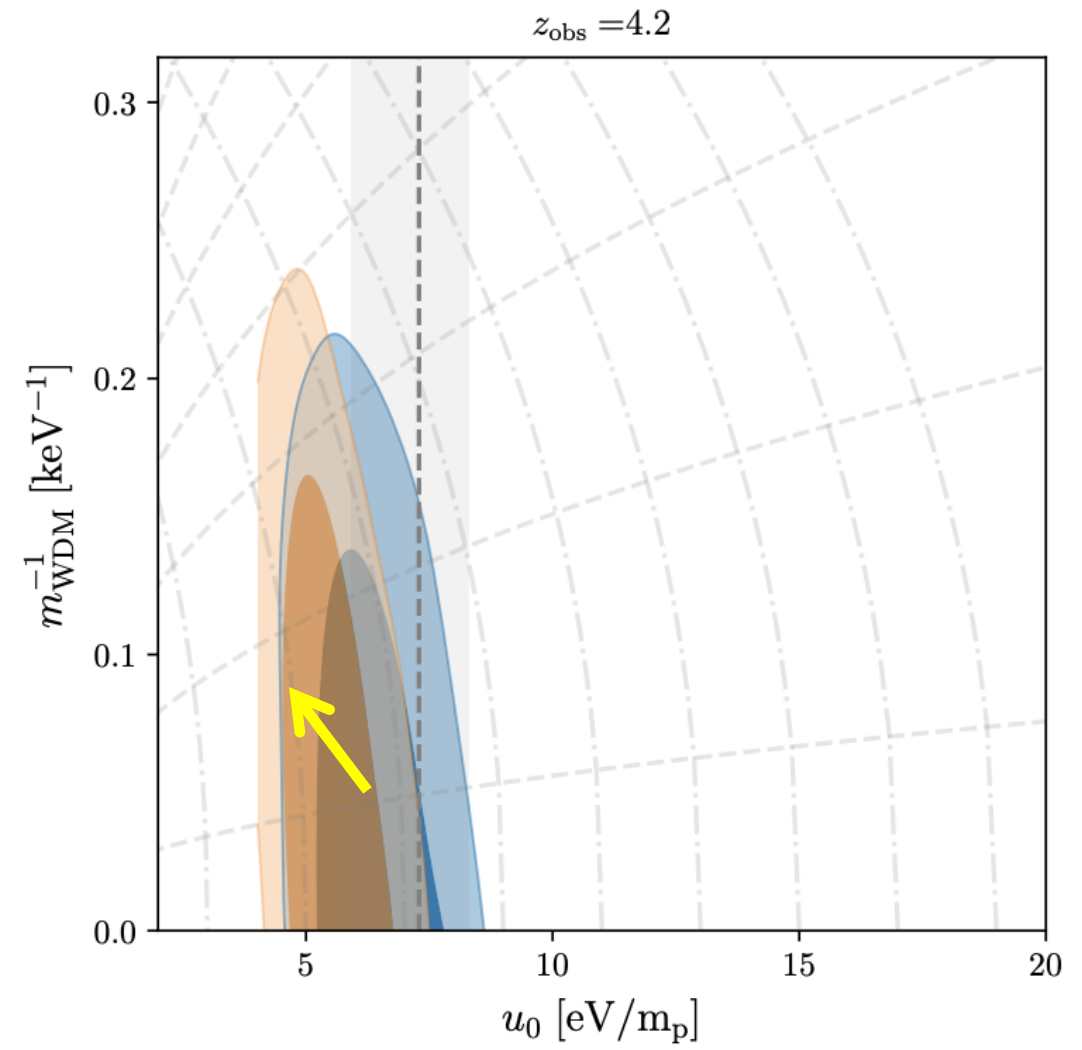
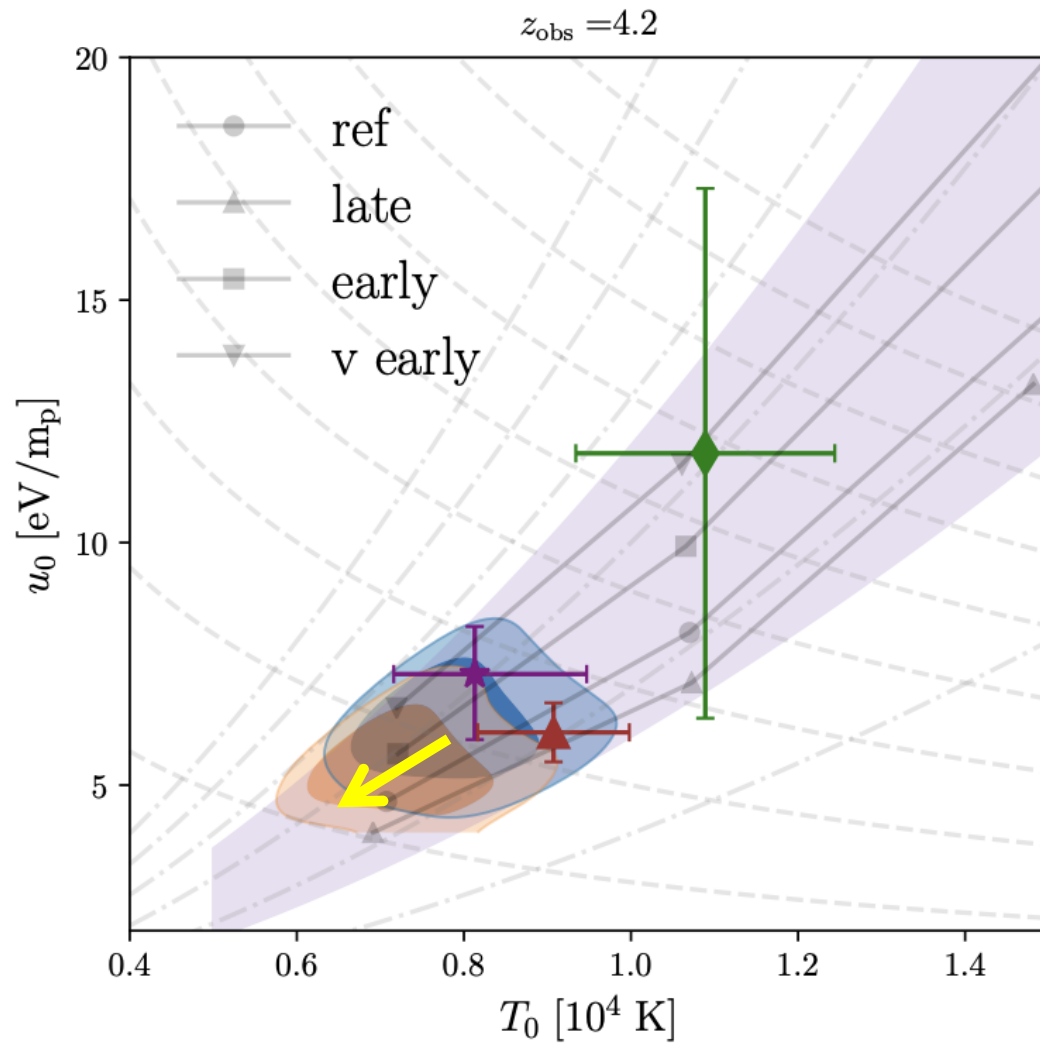
Matteo Viel



Irsic+23

Thermal WDM – inclusion of patchy correction

Matteo Viel



Irsic+23

$$\nabla_\mu \nabla^\mu \phi = m^2 \phi, \quad G_{\mu\nu} = 8\pi G T_{\mu\nu},$$

KG and Einstein equations

$$T_{\mu\nu}^\phi = g_{\mu\nu} \left(-\frac{1}{2} \partial_\rho \phi \partial^\rho \phi - \frac{1}{2} m^2 \phi^2 \right) + \partial_\mu \phi \partial_\nu \phi.$$

Energy momentum tensor
for the scalar field

$$ds^2 = -(1 + 2\Phi)dt^2 + a(t)^2(1 - 2\Phi)d\mathbf{x}^2.$$

Metric

$$\phi = \frac{1}{\sqrt{2m}} (\varphi e^{-imt} + \varphi^* e^{imt})$$

Oscillating field

$$i \left(\dot{\varphi} + \frac{3}{2} H \varphi \right) = -\frac{\partial^2 \varphi}{2a^2 m} + m \Phi \varphi,$$

Dropping higher order and averaging
over one oscillating period:
Schrodinger type eq.

$$\rho_\phi \equiv m \varphi \varphi^*, \quad v_i \equiv \frac{\partial_i \{\arg(\varphi)\}}{am} = -\frac{i}{2am} \left(\frac{\partial_i \varphi}{\varphi} - \frac{\partial_i \varphi^*}{\varphi^*} \right)$$

Defining density and velocities
of the fluid

$$\dot{v}_i + H v_i + \frac{v_j \partial_j v_i}{a} = -\frac{\partial_i \Phi}{a} + \frac{1}{2a^3 m^2} \partial_i \left(\frac{\partial^2 \sqrt{\rho_\phi}}{\sqrt{\rho_\phi}} \right)$$

Euler eq. NOTE the pressure term

$$\dot{\rho}_\phi + 3H \rho_\phi + \frac{\partial_i (\rho_\phi v_i)}{a} = 0.$$

Continuity

$$\delta_m = F\delta_\phi + (1 - F)\delta_c.$$

$$\ddot{\delta}_{\phi k} + 2H\dot{\delta}_{\phi k} + \frac{c_s^2 k^2}{a^2}\delta_{\phi k} - \frac{3}{2}H^2\delta_{mk} = 0,$$

$$\ddot{\delta}_{ck} + 2H\dot{\delta}_{ck} - \frac{3}{2}H^2\delta_{mk} = 0.$$

$$c_s^2 \equiv \frac{k^2}{4a^2m^2}, \quad \frac{k_J}{a} = \sqrt{Hm},$$

Linear perturbation theory
in CDM+scalar field model

$$\frac{k_{\text{Jeq}}}{a_0} = \frac{a_{\text{eq}}}{a_0} \sqrt{H_{\text{eq}} m} \approx 7 \text{ Mpc}^{-1} \left(\frac{m}{10^{-22} \text{ eV}} \right)^{1/2}$$

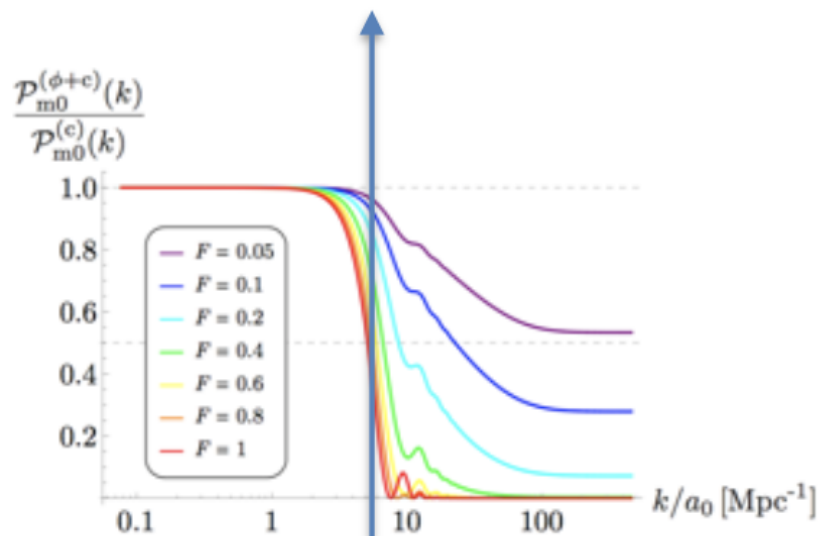
Sound speed of scalar DM and Jeans
scale definition

At $k < k_J$ no pressure

At $k > k_J$ pressure and oscillations
no growth

Comoving Jeans $k_J \sim a^{1/4}$ in MD

Important quantity is k_J at equival.

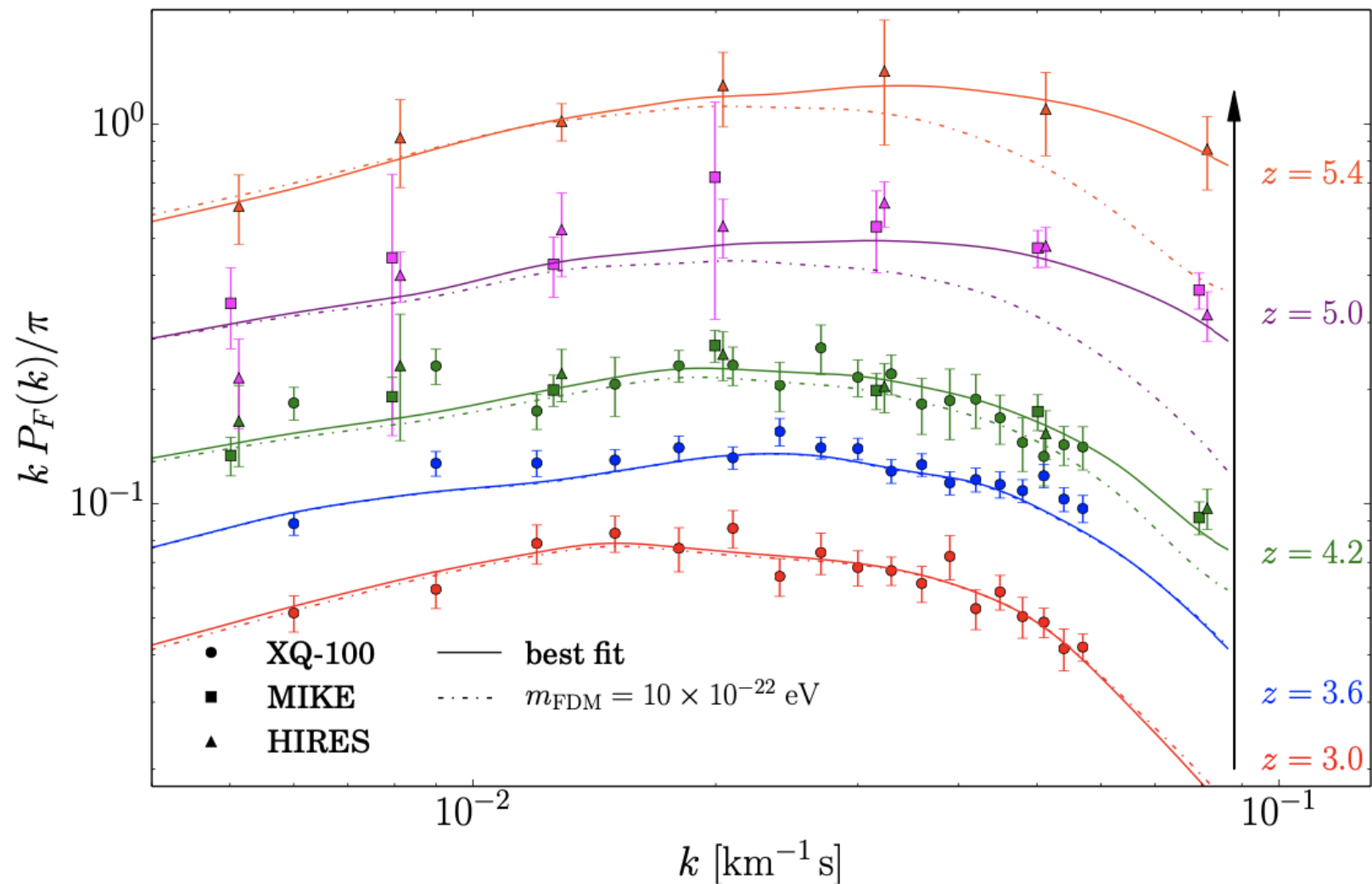


Plateau is set by FDM fraction

Cutoff scale set by FDM mass

Scalar Dark Matter - III

Matteo Viel



Irsic, Viel+ 2022 PRL

- Dark Photon Dark Matter: simple extension of the SM of particle physics

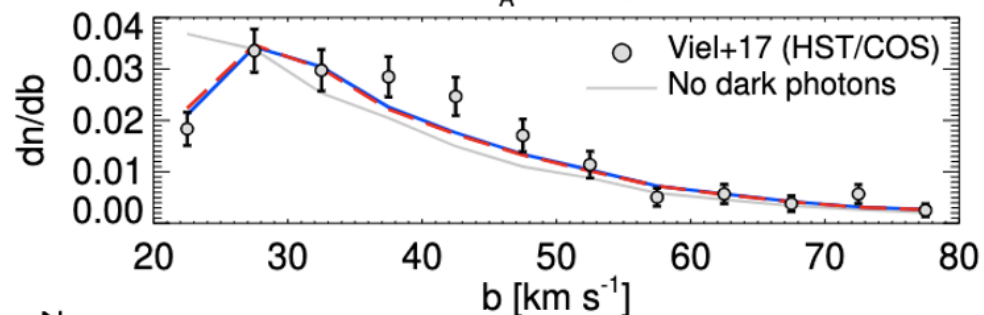
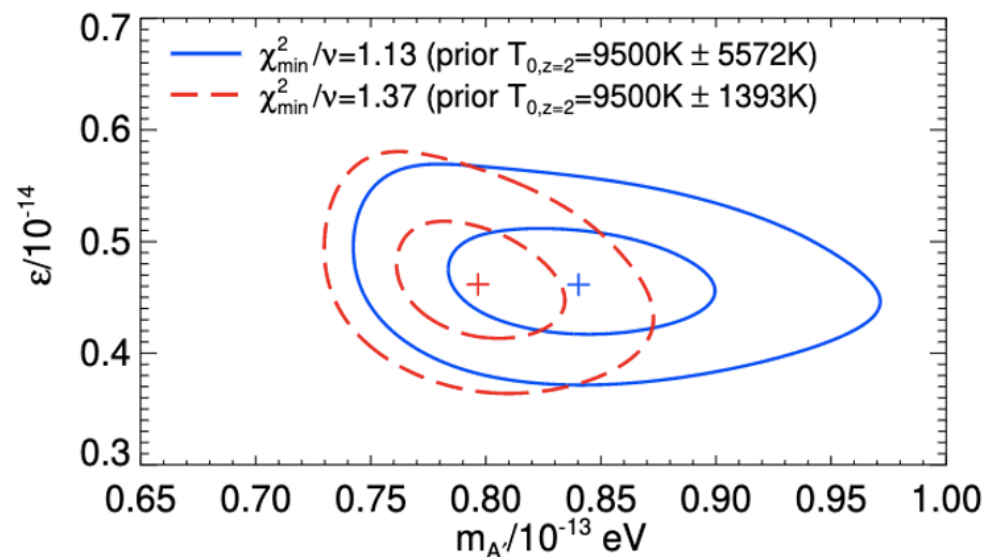
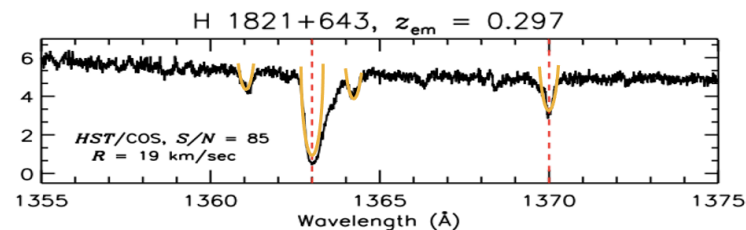
$$\mathcal{L}_{\gamma A'} = -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}(F'_{\mu\nu})^2 - \frac{\epsilon}{2}F^{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2(A'_\mu)^2$$

- Dark photon converts into standard photon when a resonance

$$E_{A' \rightarrow \gamma} \sim 2.5 \text{ eV} \left(\frac{\epsilon_{-14}}{0.5} \right)^2 \left(\frac{3}{1+z_{\text{res}}} \right)^{3/2} \left(\frac{m_{-13}}{0.8} \right)$$

The IGM as a thermometer - II

Matteo Viel

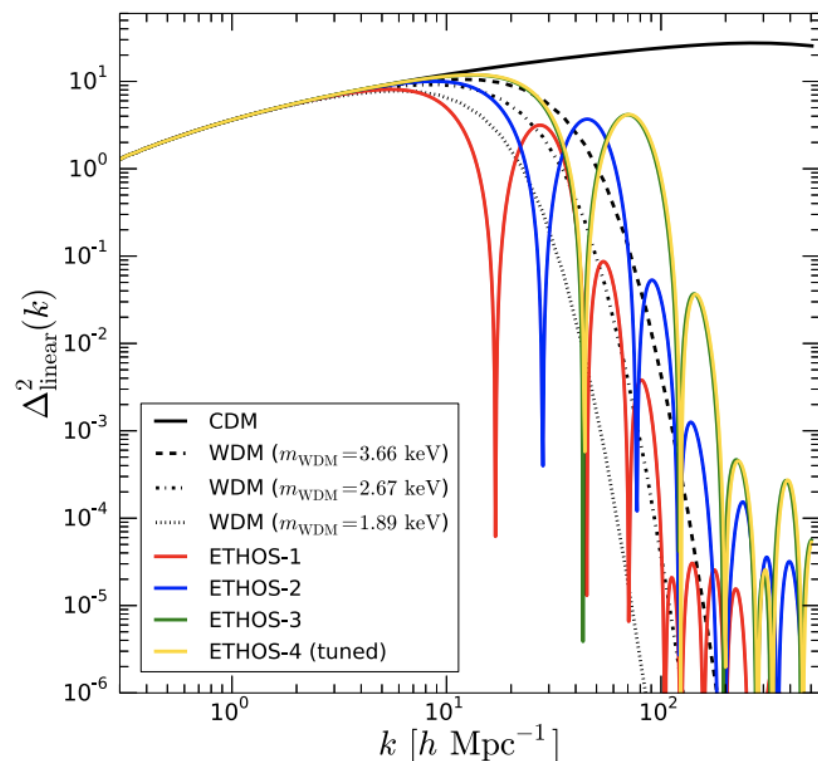


- Effect is small but can be used to place constraints on extra-heating
- At $z=0.1$ COS/HST lines are broader than

Baryon-DM or Dark radiation-DM interactions

Matteo Viel

Vogelsberger+16



- Dark Acoustic Oscillations are impacted by: 1) non-linearities; 2) projection in 1D power; 3) non-linear density-flux transformation