

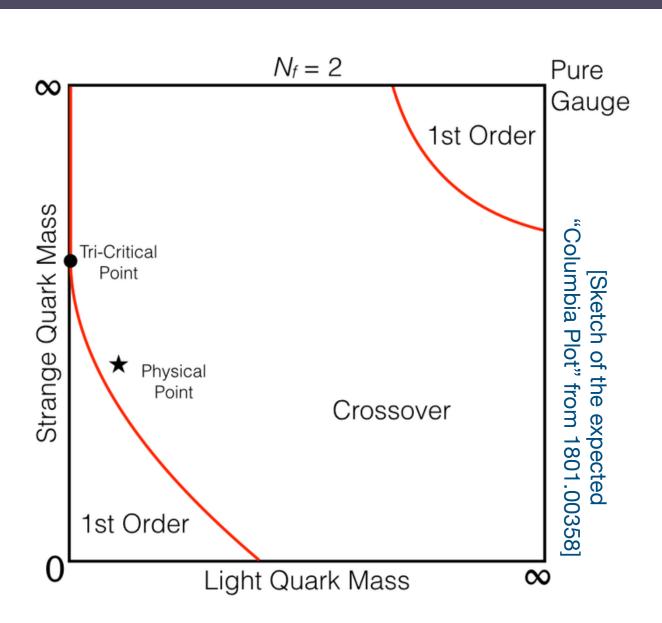
Gravitational waves from pion dark matter?

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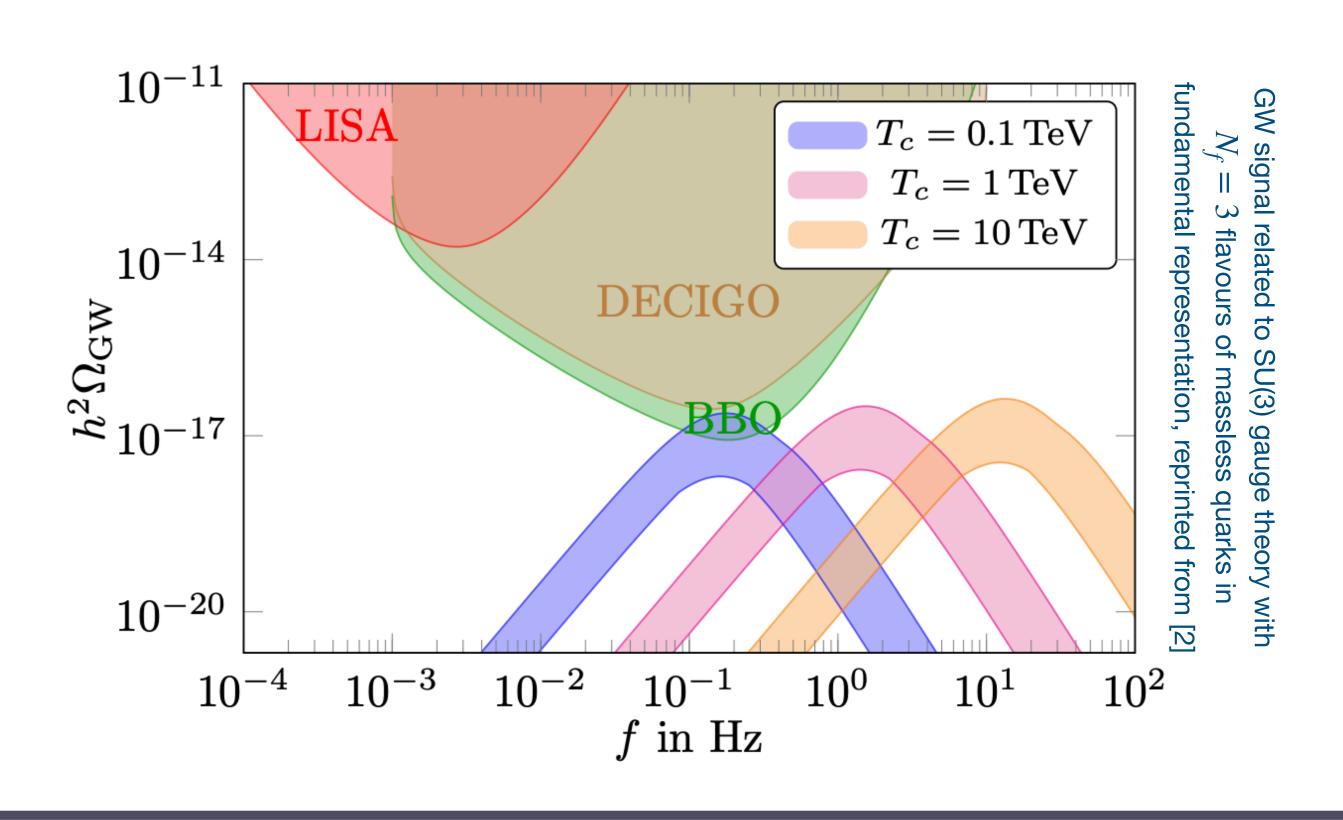
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Chiral phase transition



- In the Standard Model QCD, chiral phase transition is a smooth crossover. However, order of the chiral phase transition depends on the quark masses!
- First order phase transition in beyond-Standard-Model QCD-like theories?
- Lattice calculations for SU(3) gauge theory with $N_f=3$ quark flavours were not able to find the critical pion mass up to now, for $N_f=4$ Ref. [1] suggests that the phase transition is of first order for $m_\pi/T_c\sim m_\pi/f_\pi\lesssim 3$
- Gravitational wave signal from a chiral phase transition with massless quarks studied in [2] using PNJL model, sound waves assumed to be the main source, dark and visible sectors assumed to be thermalised



Gravitational waves from dark sectors

• For a (relatively weak) phase transition happening in a dark sector, gravitational wave signal suppressed if energy fraction in dark sector small [3]

$$\Omega_{\rm GW} h^2 \propto \left(\kappa_{\rm sw}[\alpha_{\rm (tot)}] \frac{\alpha_{\rm tot}}{\alpha_{\rm tot} + 1}\right)^2 \lesssim \kappa_{\rm sw}[\alpha_{\rm (tot)}]^2 \left(\frac{\rho_D}{\rho_{\rm tot}}\right)^2$$

where

$$\alpha = \frac{V_0}{\rho_D - V_0}$$
, $\alpha_{\text{tot}} = \frac{V_0}{\rho_{\text{tot}} - V_0}$

and V_0 is the vacuum energy released in the phase transition

• The efficiency coefficient $\kappa_{\rm sw}$ determines what fraction of the vacuum energy is transferred to the fluid motion, e.g., $\kappa_{\rm sw} \propto \alpha_{\rm (tot)}$ for relativistic bubble walls. It depends on α if the dark sector is decoupled or on $\alpha_{\rm tot}$ if the dark sector forms single fluid with the Standard Model plasma.

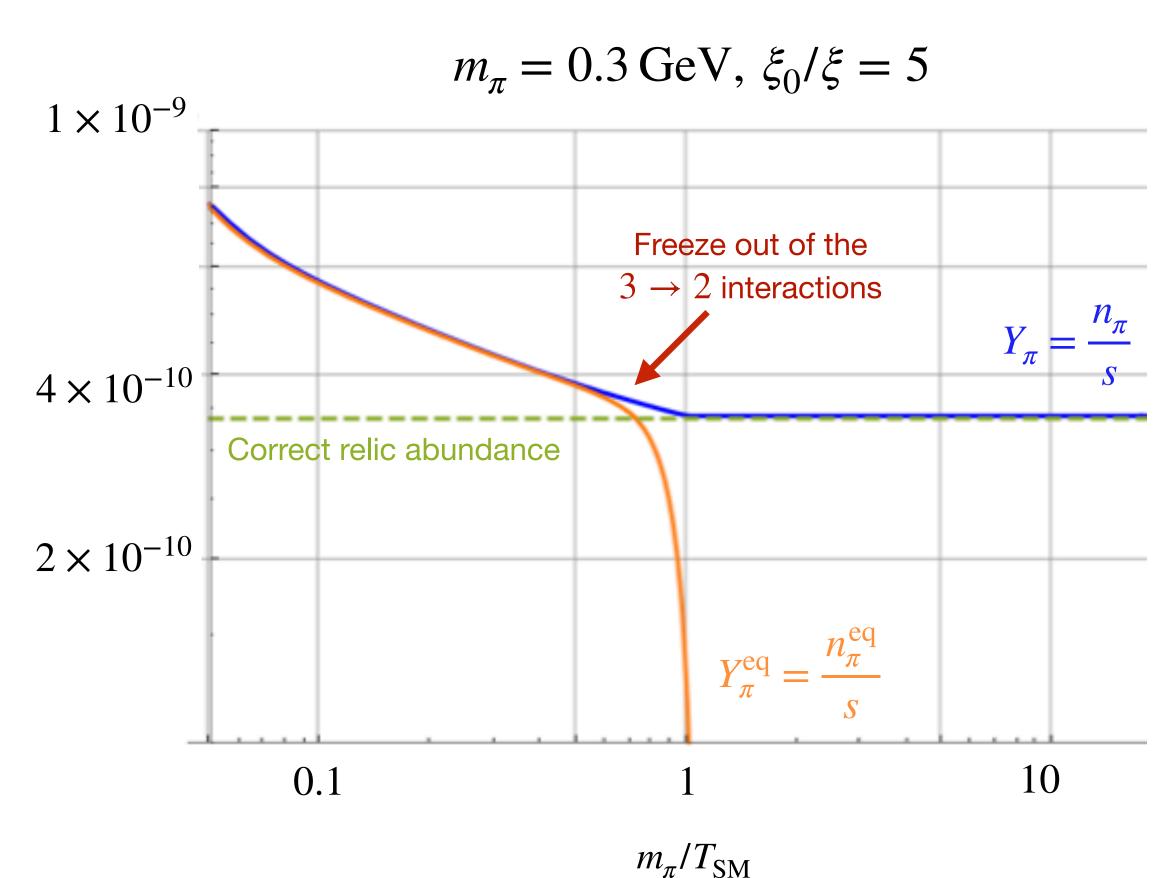
NB: For correct prediction of the GW signal, it should be checked if the dark and visible sectors form a single fluid or not!

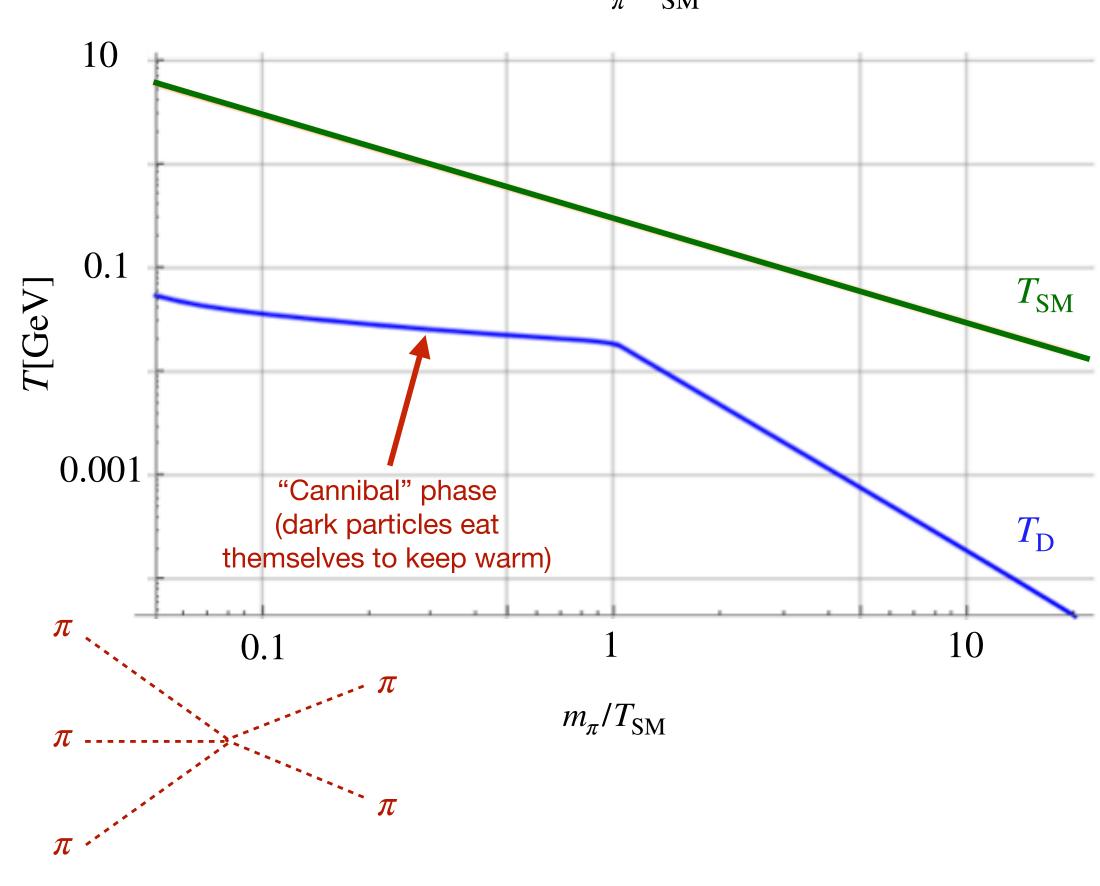
References

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- [3] M. Breitbach, J. Kopp, E. Madge, T. Opferkuch, and P. Schwaller, "Dark, Cold, and Noisy: Constraining Secluded Hidden Sectors with Gravitational Waves", JCAP **07** (2019) 007, arXiv:1811.11175 [hep-ph].
- [4] Y. Hochberg, E. Kuflik, H. Murayama, T. Volansky, and J. G. Wacker, "Model for Thermal Relic Dark Matter of Strongly Interacting Massive Particles", Phys. Rev. Lett. 115 no. 2, (2015) 021301, arXiv:1411.3727 [hep-ph].
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Strongly interacting massive particles (SIMP)

- Dark matter abundance set by $3 \rightarrow 2$ interactions among pions (due to chiral anomaly \Leftrightarrow Wess-Zumino-Witten term in Chiral Perturbation Theory) [4]
- If dark sector and SM temperatures equal: $T_D = T_{\rm SM}$, large m_π/f_π needed for correct relic abundance \Rightarrow no first order phase transition. On the other hand, low m_π/f_π achievable for **secluded SIMP dark dark sector** [5]!
- Entropy ratio in the visible and dark sectors $\xi = S_{\rm SM}/S_D$, $\xi \propto (T_{\rm SM}/T_D)^3$ for relativistic particles. Convenient normalization: $\xi_0/\xi \sim (2\times 10^8/\xi)\times m_\pi/(0.3\,{\rm GeV})$. For $\xi_0/\xi\lesssim 2$, dark matter freezes out when still relativistic.
- Boltzmann equations solved to obtain the evolution of the dark matter density and the dark sector temperature. Non-relativistic dark sector with number-changing processes cools down much slower in the "secluded" case!





• Unfortunately, $T_D=T_c\sim f_\pi$ achieved when the dark pions are still relativistic $\Rightarrow T_D\ll T_{\rm SM}$ at the phase transition, $\rho_D/\rho_{\rm tot}\propto (T_D/T_{\rm SM})^4\Rightarrow$ very suppressed gravitational wave signal \odot

